

Modeling Self-Heating Effects in 28 nm Technology Node Fully-Depleted SOI Devices

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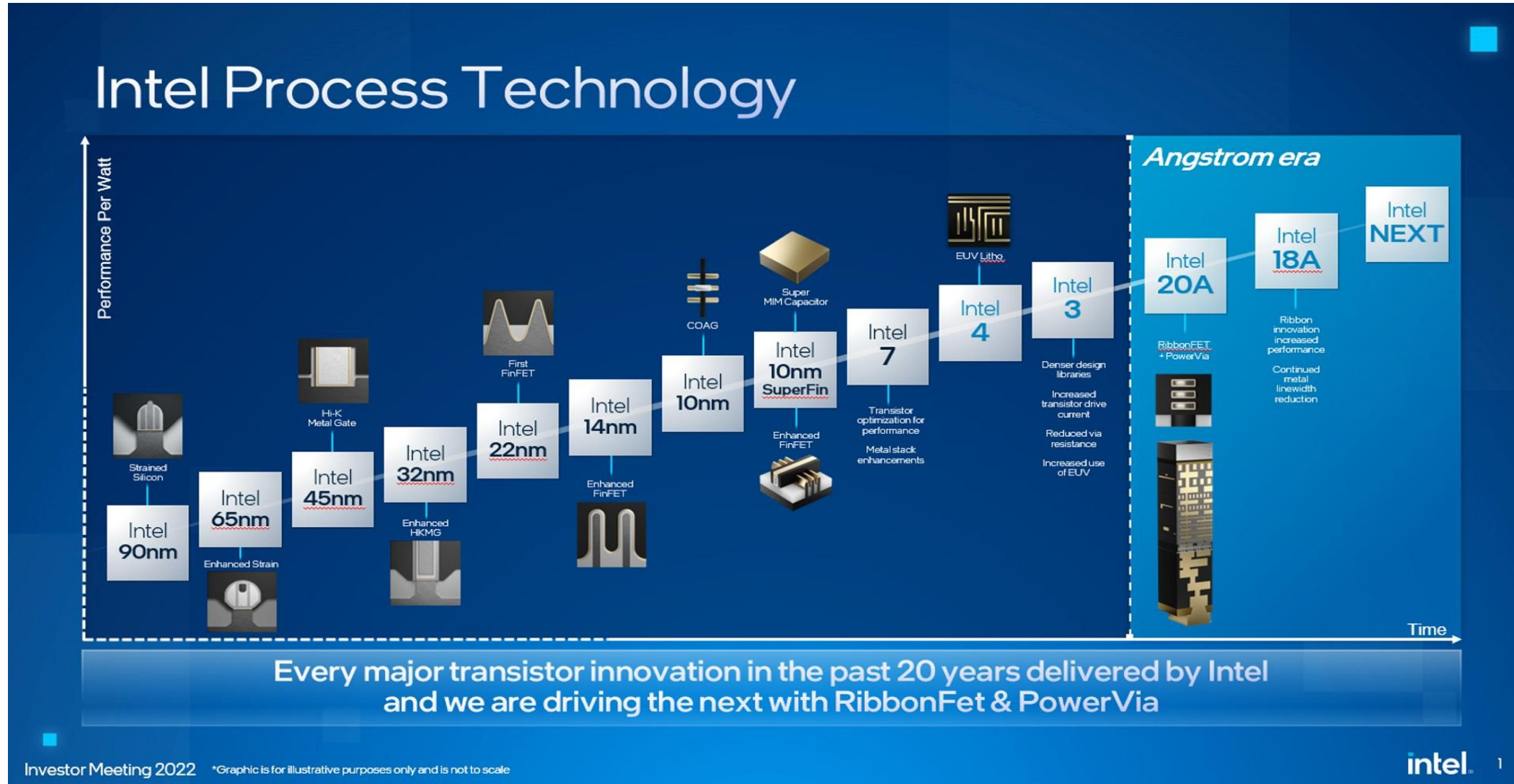
**Centro Universitario FEI, Sao Bernardo do Campo, Brazil

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Talk Outline

- Technology Trends
- Challenges to TCAD: Multiscale Nature of Self-Heating
- Approaches to Modeling Self-Heating
- ASU Approach in More Details
- Modeling Self-Heating: Thermal Conductivity
- 28 nm Technology Node FD SOI Device
- Importance of Self-Heating
- Conclusions

Technology Trends:



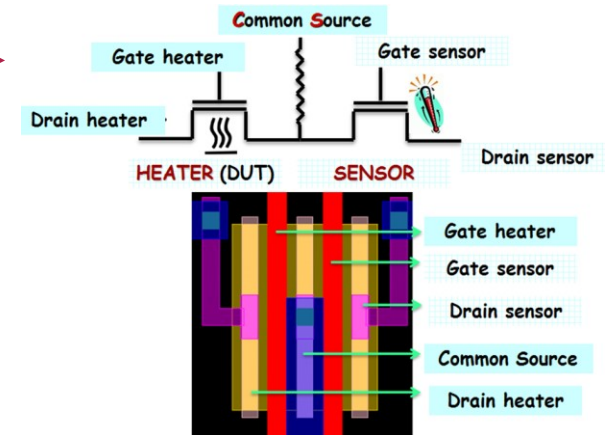
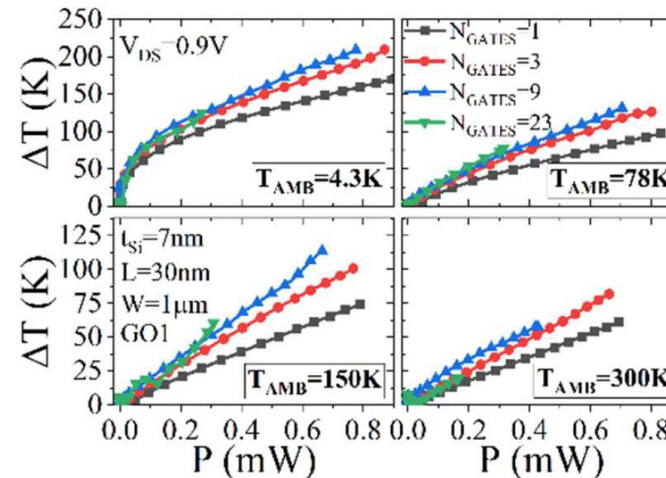
Effects:

- Random Dopant/Unintentional Dopant Fluctuations

- **Self-Heating**

Experimental Characterization:

- Heater-Sensor Approach [1]
- Gate Resistance Thermometry [2,3]



- Quantization of Charge, Tunneling and Quantum Interference

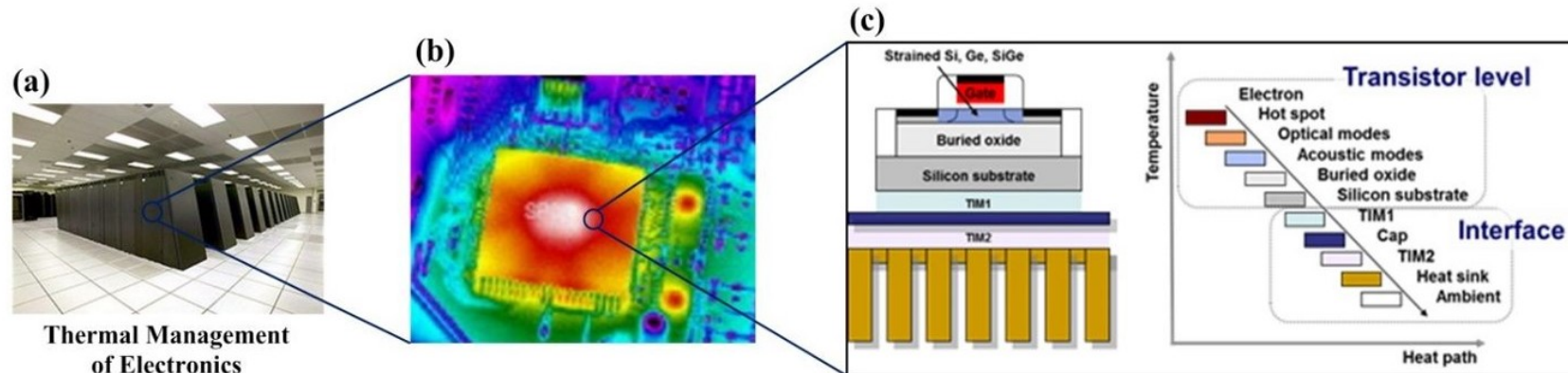
[1] E. Bury, B. Kaczer, P. J. Roussel, R. Ritzenthaler, K. Raleva, D. Vasileska, G. Groeseneken, "Experimental validation of self-heating simulations and projections for transistors in deeply scaled nodes", in Proceedings of IEEE, Reliability Physics Symposium, 2014 IEEE International, pp. XT. 8.1-XT. 8.6

[2] K. Triantopoulos et al., "Self-heating effect in FDSOI transistors down to cryogenic operation at 4.2K", IEEE Trans. Electron Devices 66, no. 8, pp. 3498-3505 (2019).

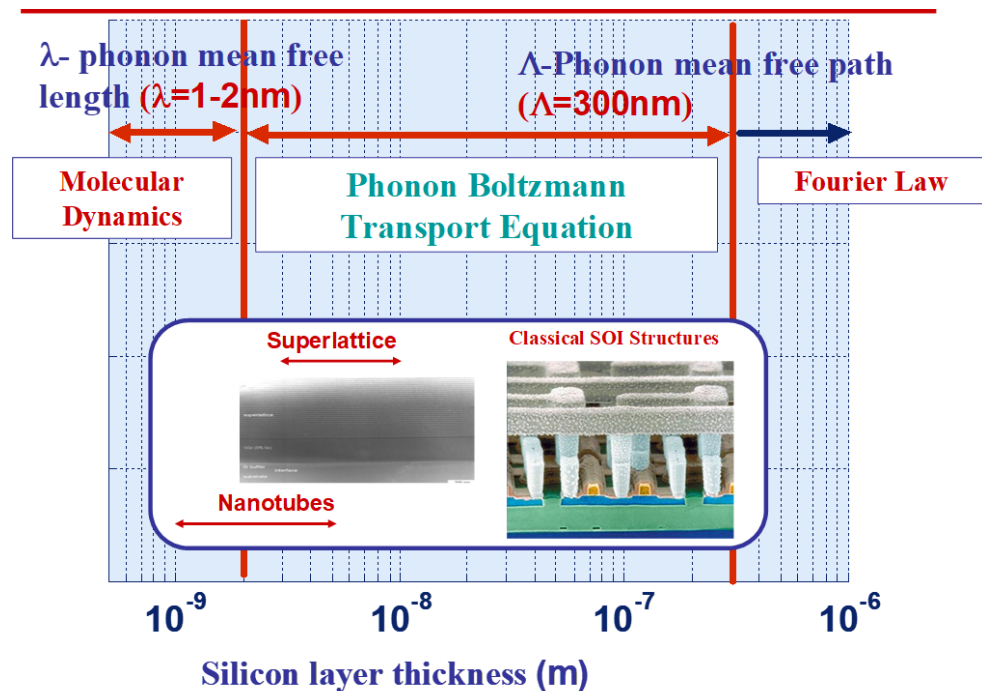
[3] M. Casse et al., "FDSOI for cryoCMOS electronics: device characterization towards compact model", IEDM (2022)

Challenge to TCAD: Multiscale Nature of Self-Heating

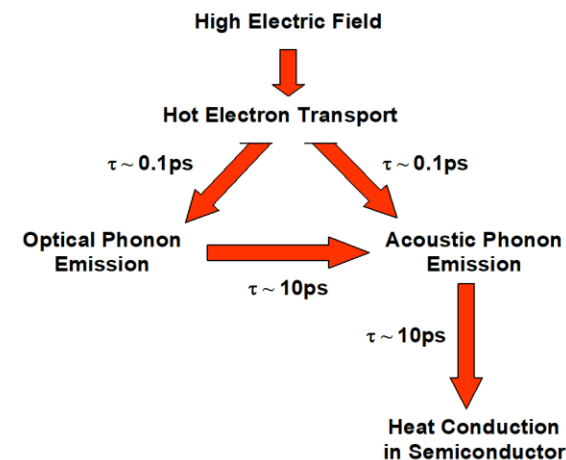
- Phonon-mediated thermal transport is inherently **multi-scale**:
 - The **wave-length of phonons** (considering phonons as waves) is typically at the **nanometer scale**;
 - The **typical size of a phonon wave energy packet** is **tens of nanometers**, while
 - The **phonon mean free path (MFP)** can be **as long as microns**.
- Multi-scale thermal transport [1]**:
 - Different heat transfer physics across **different length scales**, and
 - The physics crossing different scales is interdependent and coupled.



Approaches to Modeling Self-Heating



ASU Approach: Energy Balance Model [1,2]



$$\left(\frac{\partial}{\partial t} + v_e(\mathbf{k}) \cdot \nabla_r + \frac{e}{\hbar} \mathbf{E}(\mathbf{r}) \cdot \nabla_k \right) f = \sum_q \left\{ W_{e,q}^{k+q \rightarrow k} + W_{a,-q}^{k+q \rightarrow k} - W_{e,-q}^{k \rightarrow k+q} - W_{a,q}^{k \rightarrow k+q} \right\}$$

$$\left(\frac{\partial}{\partial t} + v_p(q) \cdot \nabla_r \right) g = \sum_k \left\{ W_{e,q}^{k+q \rightarrow k} - W_{a,q}^{k \rightarrow k+q} \right\} + \left(\frac{\partial g}{\partial t} \right)_{p-p}$$

$$C_{LO} \frac{\partial T_{LO}}{\partial t} = \frac{3nk_B}{2} \left(\frac{T_e - T_L}{\tau_{e-LO}} \right) + \frac{nm^* v_d^2}{2\tau_{e-LO}} - C_{LO} \left(\frac{T_{LO} - T_A}{\tau_{LO-A}} \right),$$

$$C_A \frac{\partial T_A}{\partial t} = \nabla \cdot (k_A \nabla T_A) + C_{LO} \left(\frac{T_{LO} - T_A}{\tau_{LO-A}} \right) + \frac{3nk_B}{2} \left(\frac{T_e - T_L}{\tau_{e-L}} \right).$$

[1] J. Lai and A. Majumdar, "Concurrent thermal and electrical modeling of submicrometer silicon devices", J. Appl. Phys., Vol. 79, 7353 (1996).

[2] K. Raleva, D. Vasilevska, S. M. Goodnick and M. Nedjalkov, "Modeling Thermal Effects in Nanodevices", IEEE Transactions on Electron Devices, vol. 55, issue 6, pp. 1306-1316, June 2008.

Approaches to Modeling Self-Heating, Cont'd

Heat Conduction Equation:

$$C_{\rho} \frac{\partial T_L}{\partial t} = \nabla \cdot (\kappa \nabla T_L) + \dot{q}_V$$

- $\dot{q}_V = \frac{1}{t} \frac{d}{dV} \sum (\hbar\omega_{em} - \hbar\omega_{ab})$
Net phonon emission approach [1,2]
- $\dot{q}_V = C_{Lo} \left(\frac{T_{Lo} - T_L}{\tau_{Lo-A}} \right) + \frac{3nk_B}{2} \left(\frac{T_e - T_L}{\tau_{e-L}} \right) + \frac{nm^*v_d^2}{2\tau_{e-L}}$
Lai and Majumdar [3]
- $\dot{q}_V = \vec{J} \cdot \vec{E} + \text{corrections}$
Wachutka model [4]

[1] N. J. Pilgrim, "Electro-thermal Monte Carlo simulation of semiconductor devices", PhD Dissertation, University of Leeds, UK, 2003.

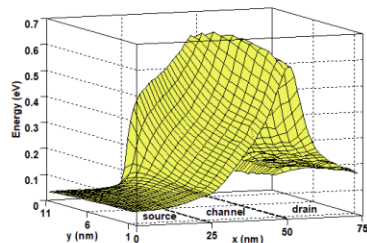
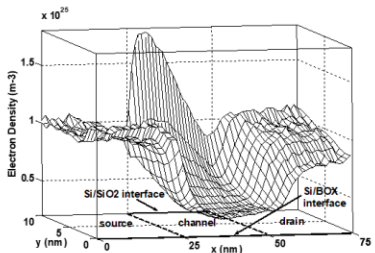
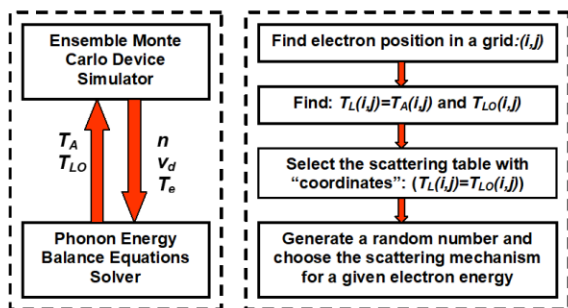
[2] E. Pop, R. W. Dutton, K. E. Goodson, "Analytic band Monte Carlo model for electron transport in Si including acoustic and optical phonon dispersion", J. Appl. Phys., Vol. 96, 4998 (2004).

[3] J. Lai and A. Majumdar, "Concurrent thermal and electrical modeling of submicrometer silicon devices", J. Appl. Phys., Vol. 79, 7353 (1996).

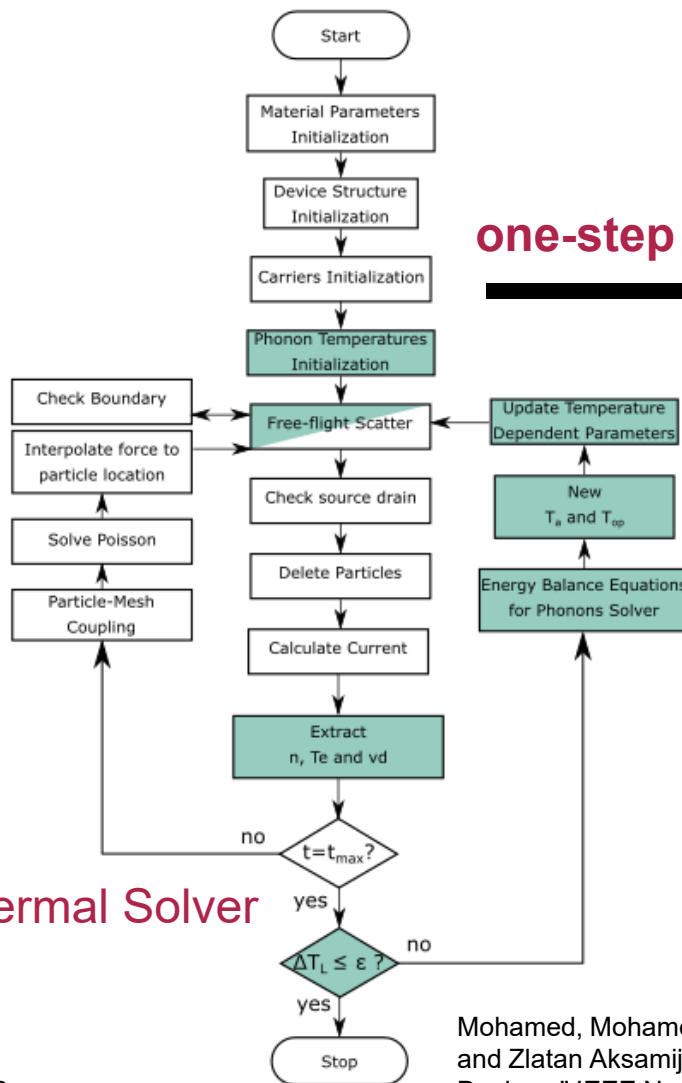
[4] Wachutka, G.K., "Rigorous Thermodynamic Treatment of Heat Generation in Semiconductor Device Modeling", *IEEE Trans., Computer-Aided Design* Vol. 9, No. 11 (1990): 1141-1149.

ASU Approach in More Details

Exchange of variables
in the electro-thermal solver

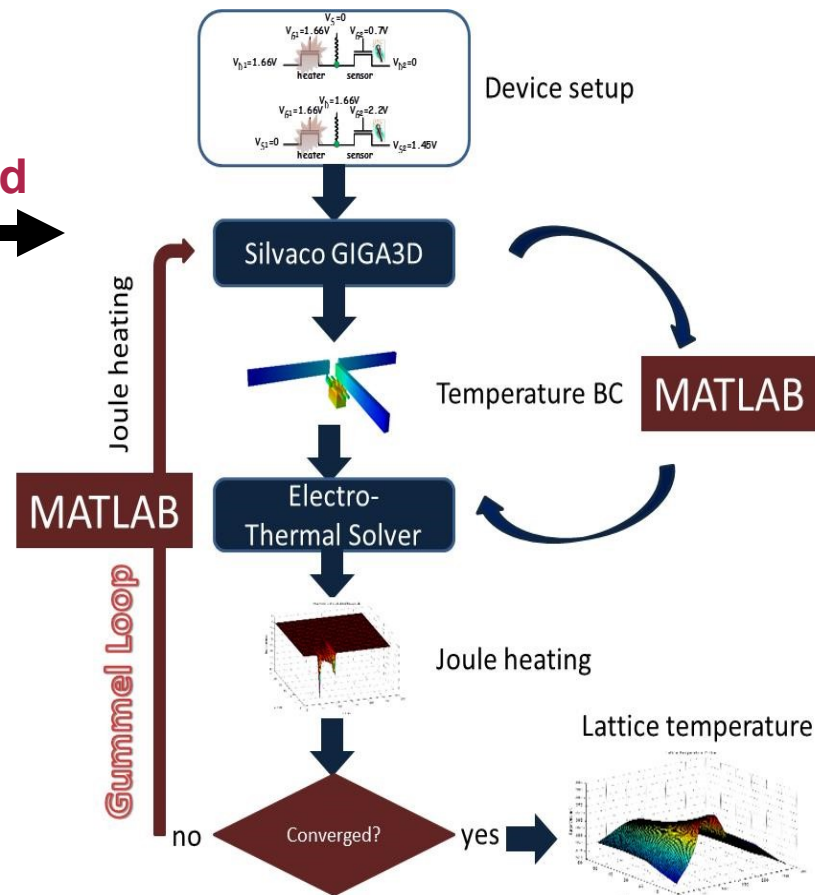


Electro-Thermal Solver



one-step forward

Multi-Scale Electro-Thermal Solver

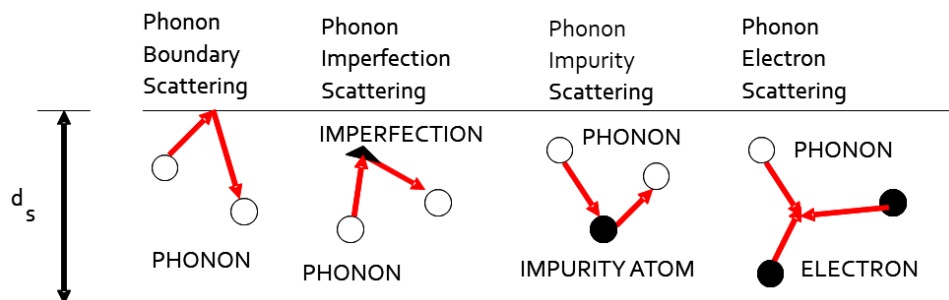


Modeling Self-Heating: Thermal Conductivity

A measure of a material's ability to transfer thermal energy by conduction.

Thermal conductivity k has two different contributions: $K = K_{phonon} + K_{electron}$

- The electronic contribution to the thermal conductivity can be calculated using the Wiedemann-Franz law.
- The phonon contribution depends upon:
 - Scattering by Lattice Imperfections
 - Defects, dislocations, **boundaries** \longrightarrow
(ELASTIC, energy and momentum are conserved)
 - Phonon-Electron Scattering
 - Phonon-Phonon Interactions**
 - NORMAL N – processes \rightarrow Energy and momentum are conserved
 - UMKLAPP U-processes \rightarrow Only energy is conserved

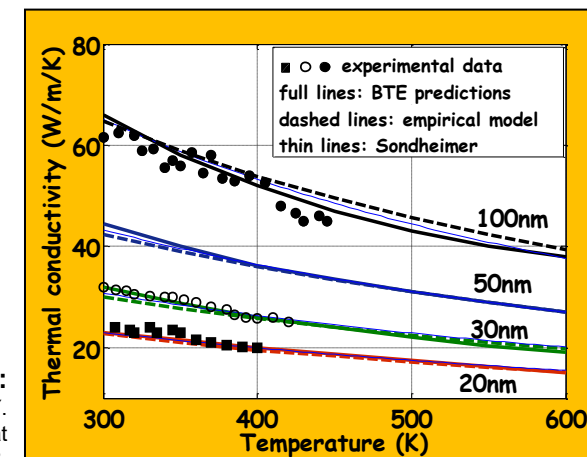


Phonon-Boundary Scattering (thin films)

$$\kappa(z) = \kappa_0(T) \int_0^{\pi/2} \sin^3 \theta \left\{ 1 - \exp\left(-\frac{a}{2\lambda(T)\cos\theta}\right) \cosh\left(\frac{a-2z}{2\lambda(T)\cos\theta}\right) \right\} d\theta$$

$$\lambda(T) = \lambda_0 (300/T)$$

$$\kappa_0(T) = \frac{135}{a + bT + cT^2} \text{ W/m/K}$$



Experimental Data:

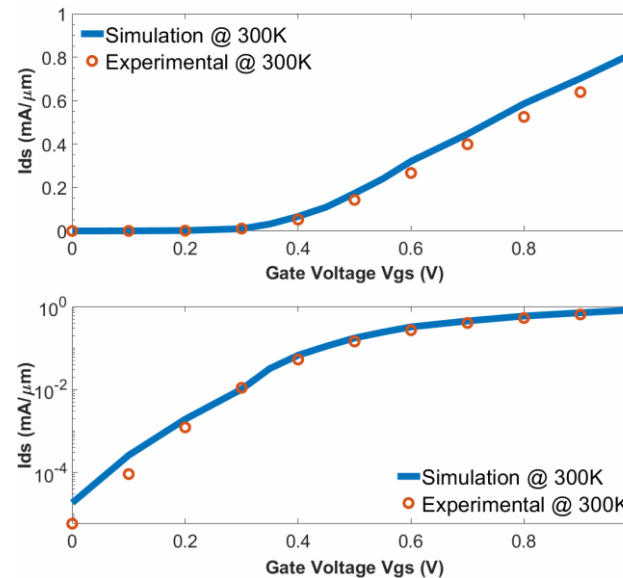
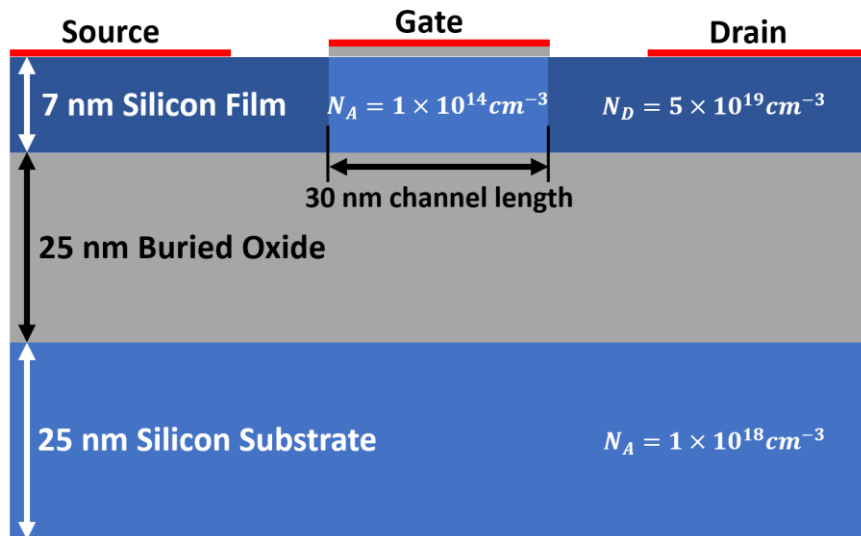
M. Asheghi, M. N. Touzelbaev, K. E. Goodson, Y. K. Leung, and S. S. Wong, ASME Journal of Heat Transfer, Vol.120, pp. 30-33, 1998.

D. Vasileska, K. Raleva and S. M. Goodnick, "Electrothermal Studies of FD SOI Devices That Utilize a New Theoretical Model for the Temperature and Thickness Dependence of the Thermal Conductivity", IEEE Transactions on Electron Devices, Vol. 57, pp. 726 – 728 (2010).

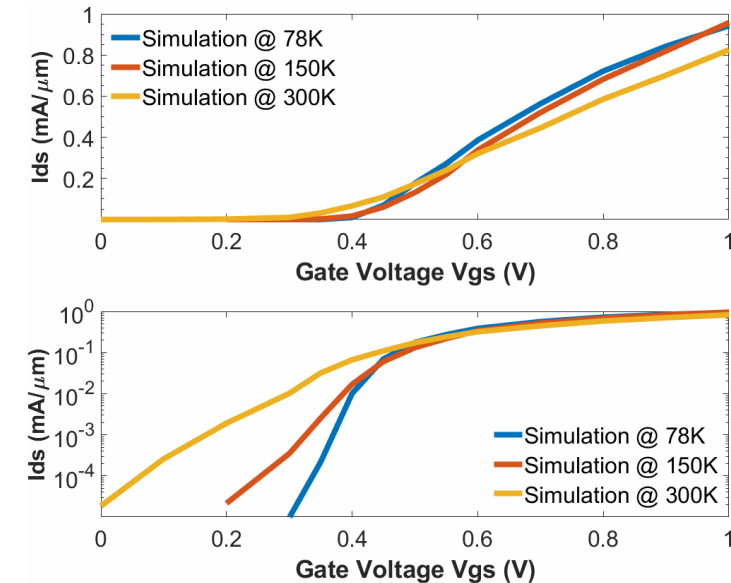
28 nm Technology Node FD SOI Device

Quantum Features:

- Superposition of states → Phase coherence → Low-temperature → No scattering
- Entanglement
- Quantum Tunneling



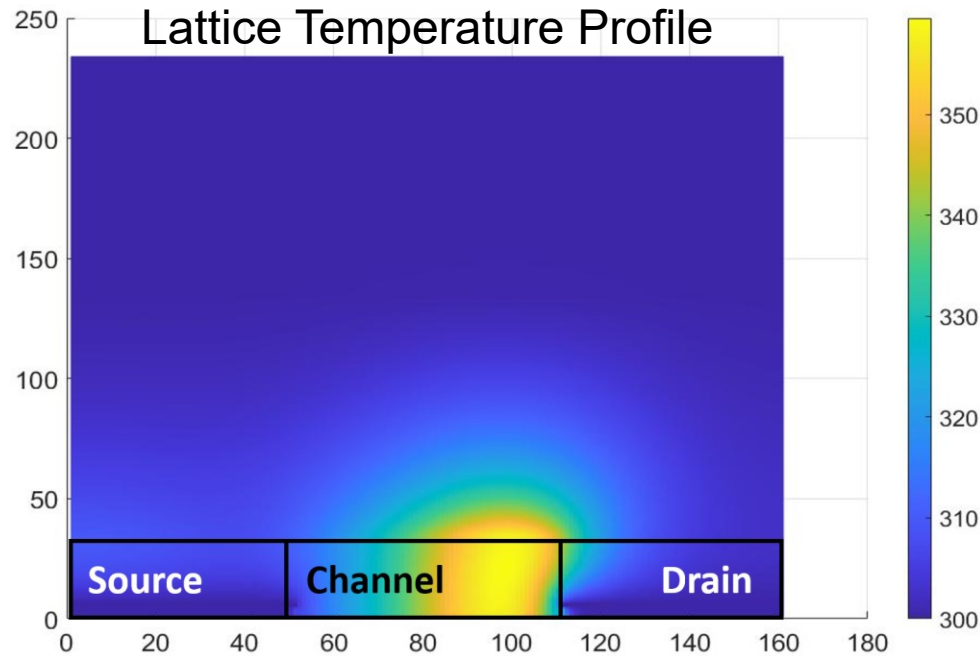
Simulated transfer characteristics compared with available experimental data from Ref. [1] at $T=300\text{K}$. The applied drain voltage is $V_{ds}=0.9\text{V}$.



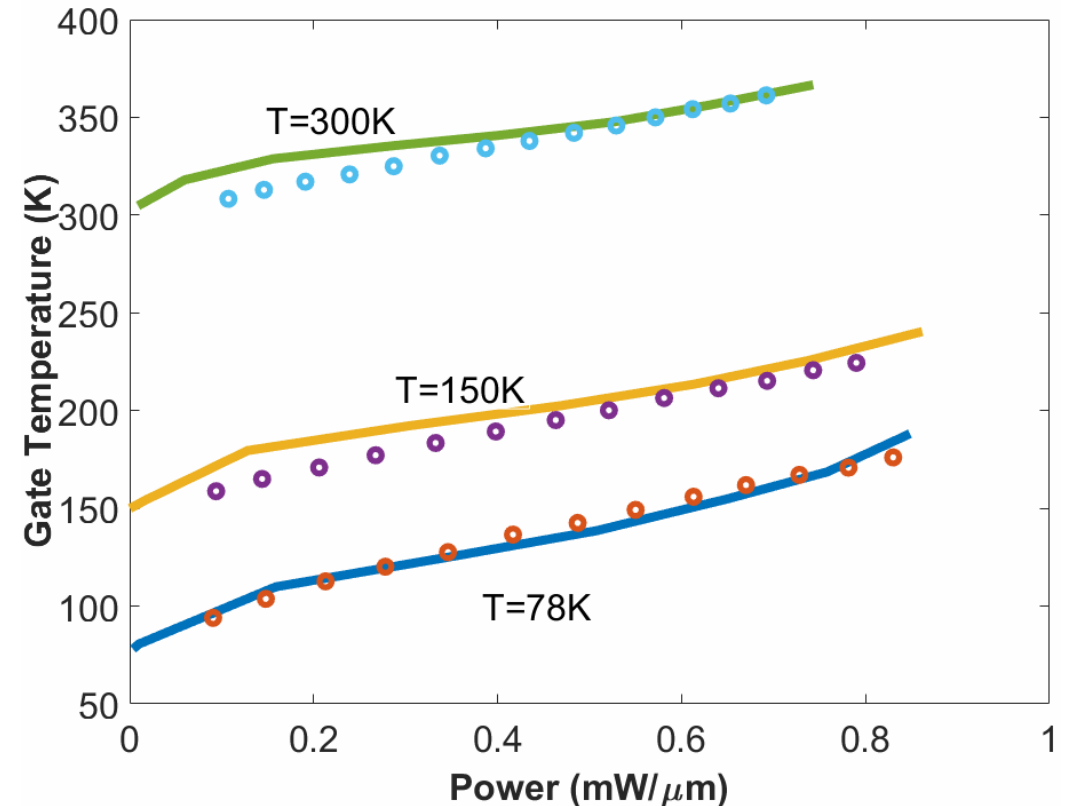
Simulated transfer characteristics at 78K, 150K and 300K. The applied drain voltage is $V_{ds}=0.9\text{V}$.

[1] M. Casse et al., "FDSOI for cryoCMOS electronics: device characterization towards compact model", IEDM (2022)

Importance of Self-Heating



- Lattice temperature profile at ambient temperature $T=300\text{K}$ for $V_{gs}=0.6\text{V}$ and $V_{ds}=0.9\text{V}$.
- We used the following boundary conditions in the simulation:
 - Fixed temperature $T=300\text{K}$ at source and drain contacts,
 - Zero heat flux at the gate contact.



Comparison of simulated (solid lines) and experimental (open circles) temperature under the gate for various input powers at $T=78\text{K}$, 150K and 300K .

Conclusions

- A 2D/3D electro-thermal device simulator has been developed at ASU to study self-heating effects in:
 - FD SOI Devices, nanowire transistors, dual gate device structures, FinFETs, CMOS inverters and simple two transistor (CS and CD) circuits.
- Low temperature simulations of FD SOI devices confirm the experimental findings that self-heating is very important at cryogenic temperatures.

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Thank You!