

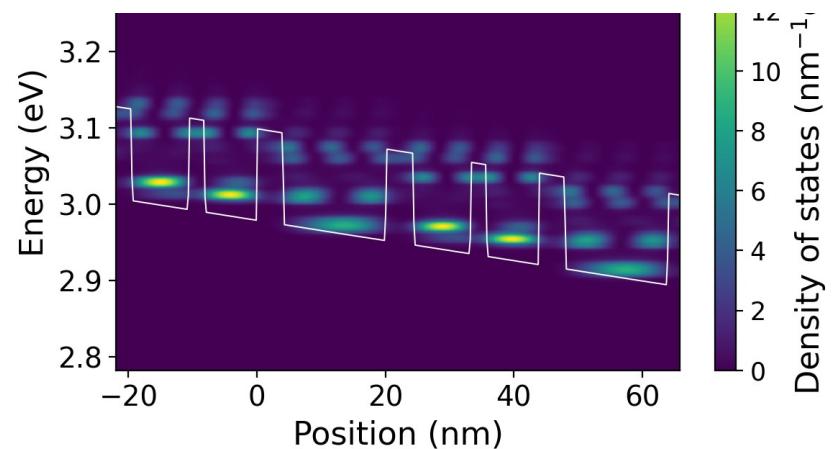
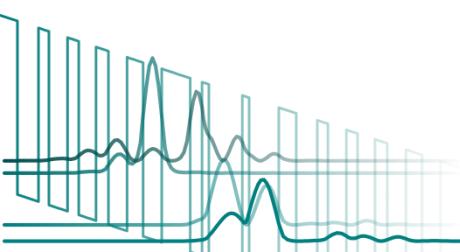
NEGF Basis in Multiband Models for Broken-Gap Sb-Based Tunneling Devices

Takuma Sato^{1,2}

Stefan Birner^{1,2}, Christian Jirauschek², Thomas Grange³

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2. School of Computation, Information and Technology, Technical University of Munich, Germany
3. nextnano Lab SAS, Corenc, France

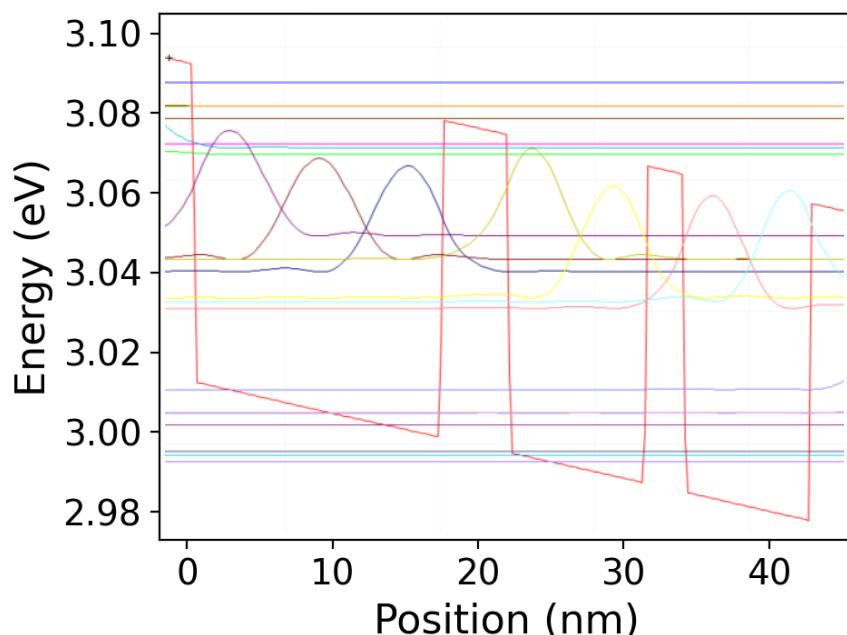
- Schrodinger-Poisson-drift-diffusion solver ([nextnano++](#))
Quasi-equilibrium semiclassical transport



Tunneling charge transport

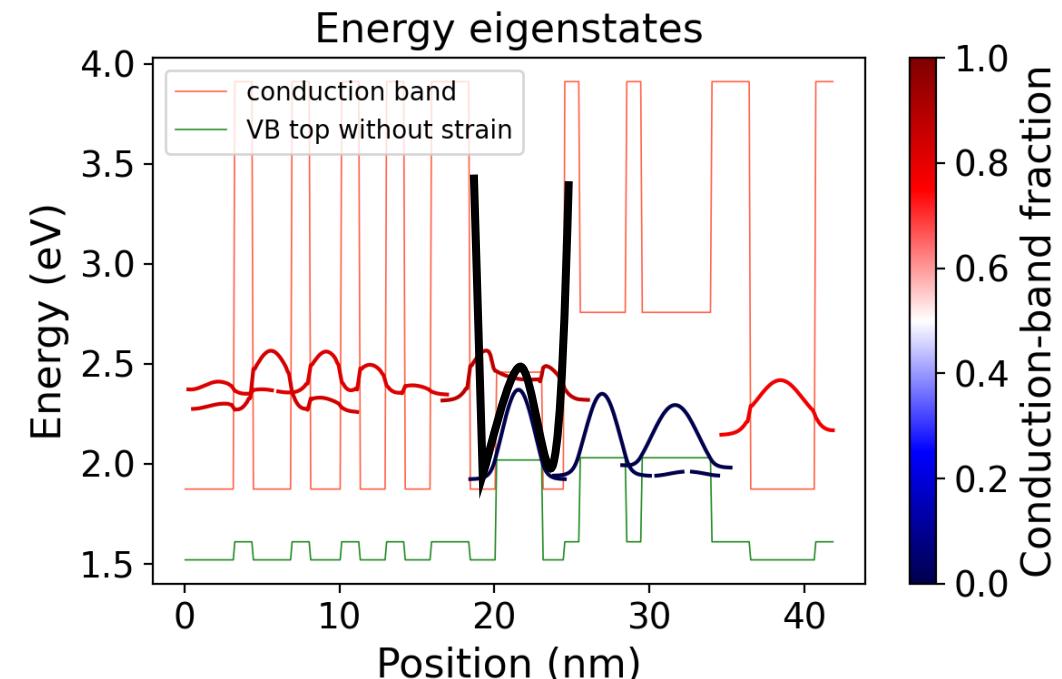
III-Sb broken-gap tunneling devices – photodetector, lasers

Intersubband e.g. QCLs



nextnano.NEGF

Interband e.g. ICLs

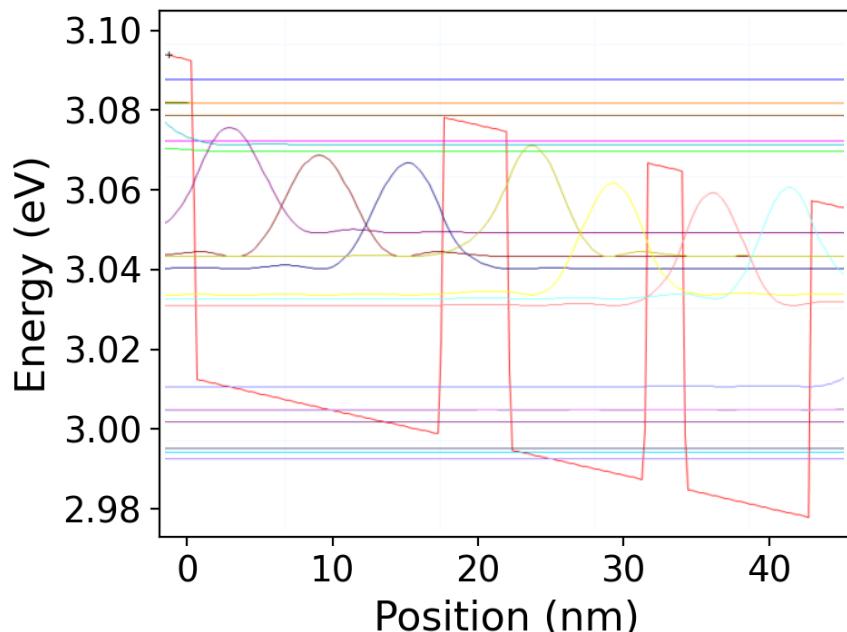


nextnano.NEGF x (this project)

Tunneling charge transport

III-Sb broken-gap tunneling devices – photodetector, lasers

Intersubband e.g. QCLs



Interband e.g. ICLs

We will simulate **non-equilibrium charge distribution** in steady states

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nextnano.NEGF x (this project)

Tunneling charge transport

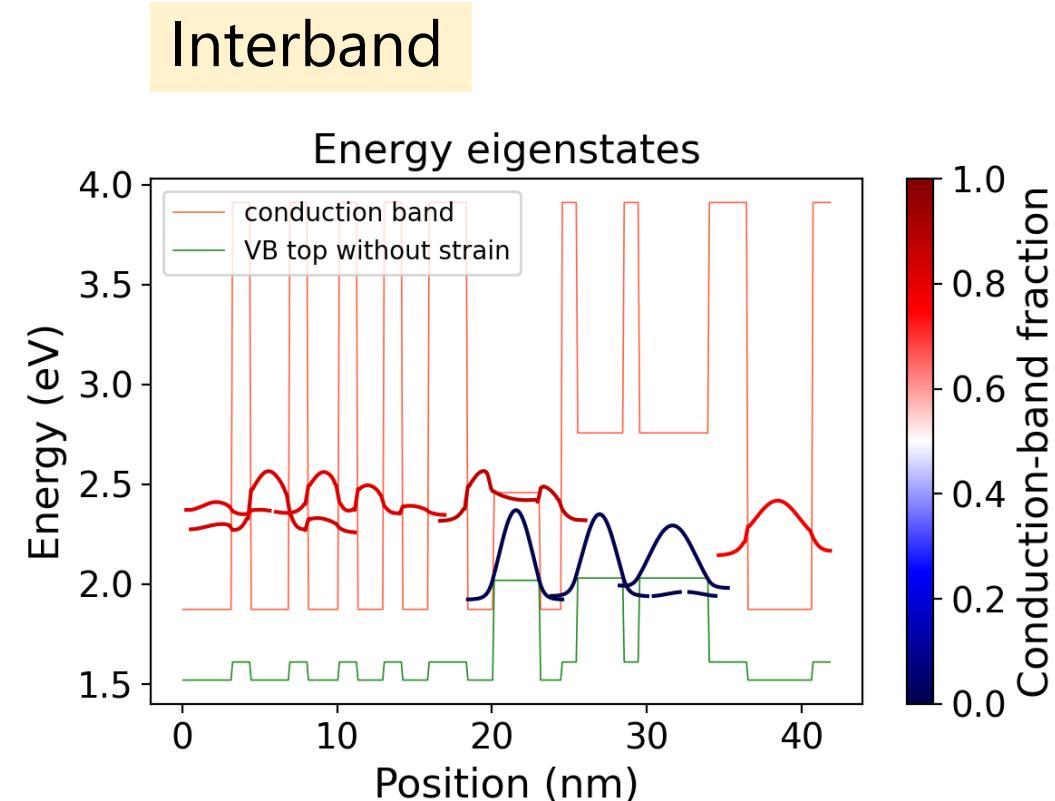
Accuracy

Reduce
problem size

$$H_0 + H_{scatt}$$

- NEGF formalism
rather than Schrodinger-Poisson
- 8-band $k \cdot p$ model
for accurate energy levels & interband coupling
- Consistent Hamiltonian discretization

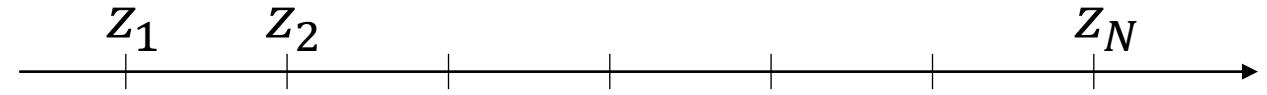
Robust against spurious solutions



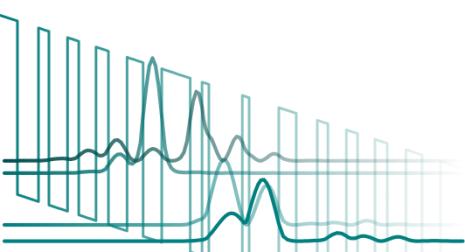
Consistent Hamiltonian discretization

8-band k.p z-dependence

$$\mathcal{H} = \mathcal{H}_{\text{band}} \otimes \mathcal{H}_{\text{spatial}}$$



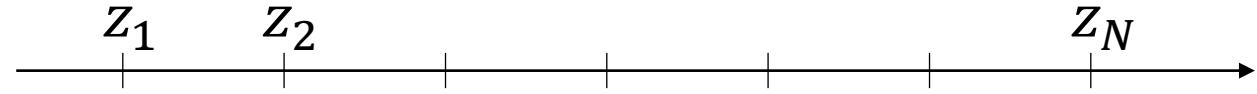
Discretization = matrix representation of the Hamiltonian in \mathcal{H}



Consistent Hamiltonian discretization

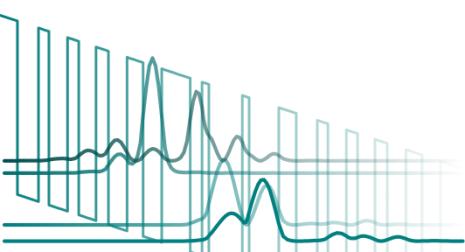
8-band k.p z-dependence

$$\mathcal{H} = \mathcal{H}_{\text{band}} \otimes \mathcal{H}_{\text{spatial}}$$



1. Heterostructure \leftrightarrow bulk
2. Avoid oscillatory solutions and peaks at the interfaces
3. Consistent discretization of the 1st- and 2nd-order derivatives
4. Hamiltonian must be Hermitian!

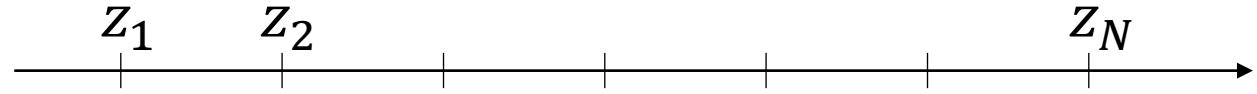
$$\psi_{n\mathbf{k}} = e^{i\mathbf{k}\mathbf{r}} u_{n\mathbf{k}}(\mathbf{r})$$



Consistent Hamiltonian discretization

8-band k.p z-dependence

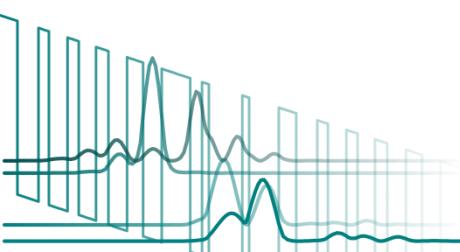
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1. Heterostructure \leftrightarrow bulk
2. Avoid oscillatory solutions and peaks at the interfaces
Foreman, PRB (1997)
3. Consistent discretization of the 1st- and 2nd-order derivatives
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$$\psi_{n\mathbf{k}} = e^{i\mathbf{k}\mathbf{r}} u_{n\mathbf{k}}(\mathbf{r})$$

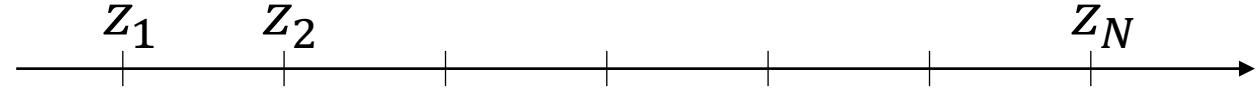
$$H = H^{(0)} + H^{(L)} \frac{d}{dz} + \frac{d}{dz} H^{(R)} + \frac{d}{dz} H^{(2)} \frac{d}{dz}$$

Consistent Hamiltonian discretization

8-band k.p z-dependence

$$\mathcal{H} = \mathcal{H}_{\text{band}} \otimes \mathcal{H}_{\text{spatial}}$$

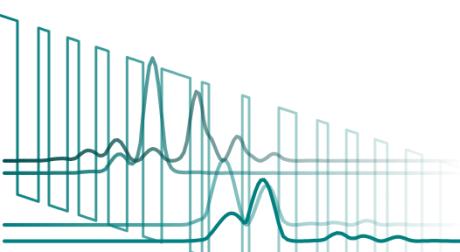


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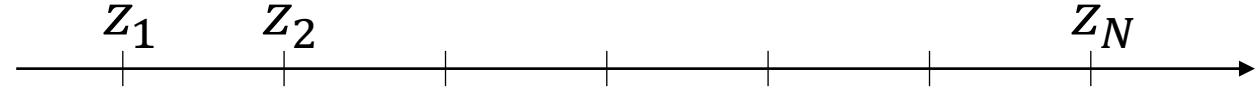
$$\hat{d}_F \gamma f(z) \Big|_i = \frac{1}{h} (\gamma_{i+1} f_{i+1} - \gamma_i f_i),$$
$$\gamma \hat{d}_F f(z) \Big|_i = \frac{1}{h} \gamma_i (f_{i+1} - f_i),$$
$$\hat{d}_B \gamma f(z) \Big|_i = \frac{1}{h} (\gamma_i f_i - \gamma_{i-1} f_{i-1}),$$
$$\gamma \hat{d}_B f(z) \Big|_i = \frac{1}{h} \gamma_i (f_i - f_{i-1}),$$



Consistent Hamiltonian discretization

8-band k.p z-dependence

$$\mathcal{H} = \mathcal{H}_{\text{band}} \otimes \mathcal{H}_{\text{spatial}}$$



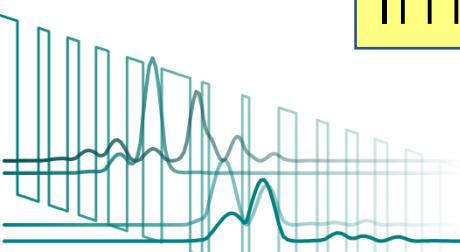
1. Heterostructure \leftrightarrow bulk
2. Avoid oscillatory solutions and peaks at the interfaces
Foreman, PRB (1997)
3. Consistent discretization of the 1st- and 2nd-order derivatives
Frensley *et al.*, arXiv (2014)
4. Hamiltonian must be Hermitian!

Implemented

$$\psi_{n\mathbf{k}} = e^{i\mathbf{k}\mathbf{r}} u_{n\mathbf{k}}(\mathbf{r})$$

$$H = H^{(0)} + H^{(L)} \frac{d}{dz} + \frac{d}{dz} H^{(R)} + \frac{d}{dz} H^{(2)} \frac{d}{dz}$$

$$\hat{d}_F \gamma f(z) \Big|_i = \frac{1}{h} (\gamma_{i+1} f_{i+1} - \gamma_i f_i),$$
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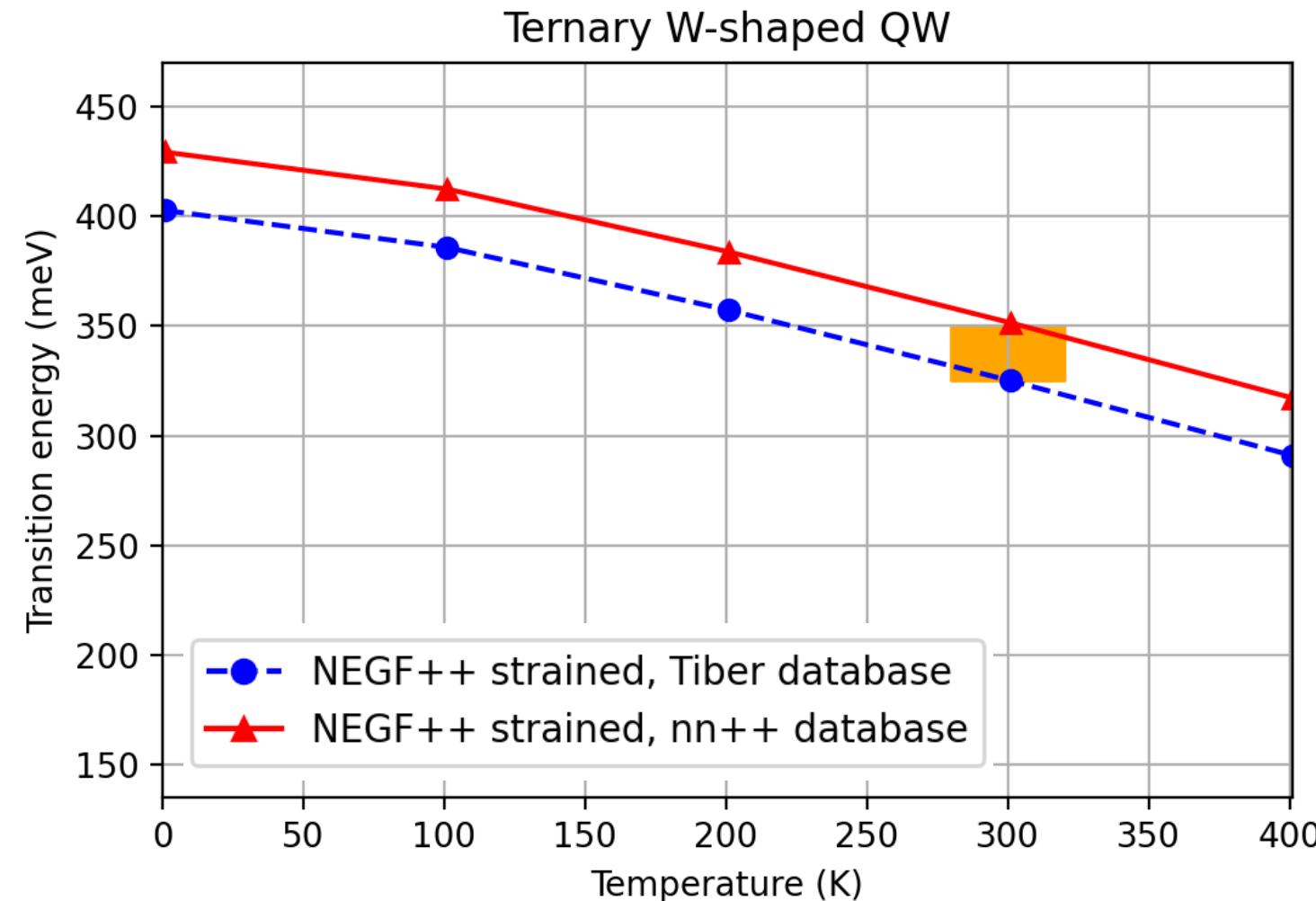
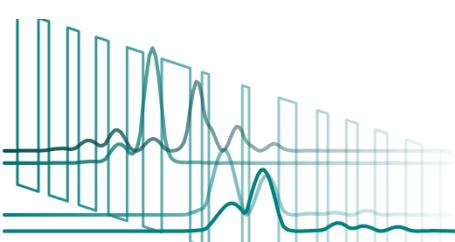
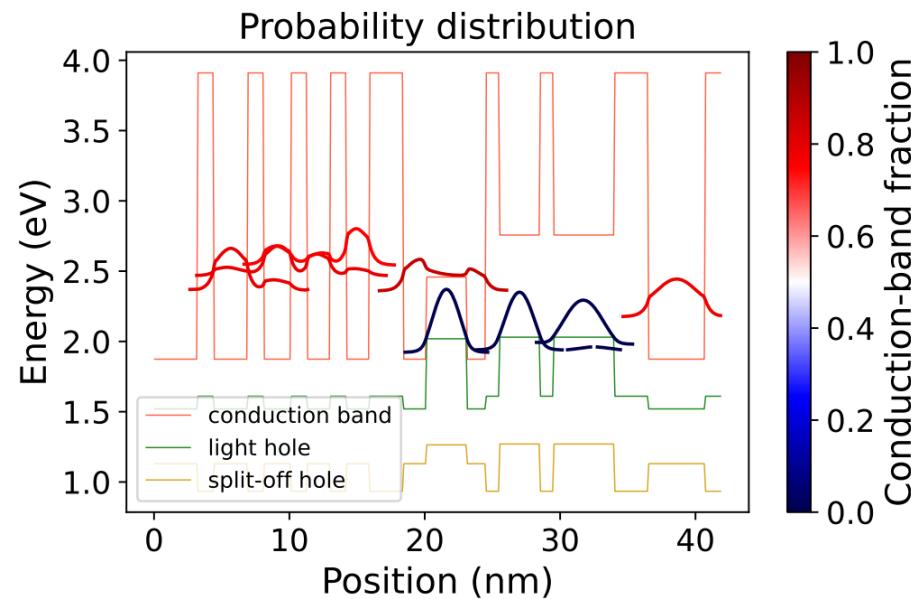


Transition energy in W-shaped QW

Measurement:

325 – 350 meV (3.6 - 3.9 μm)

Vurgaftman, et al., Nat. Comm. (2012)

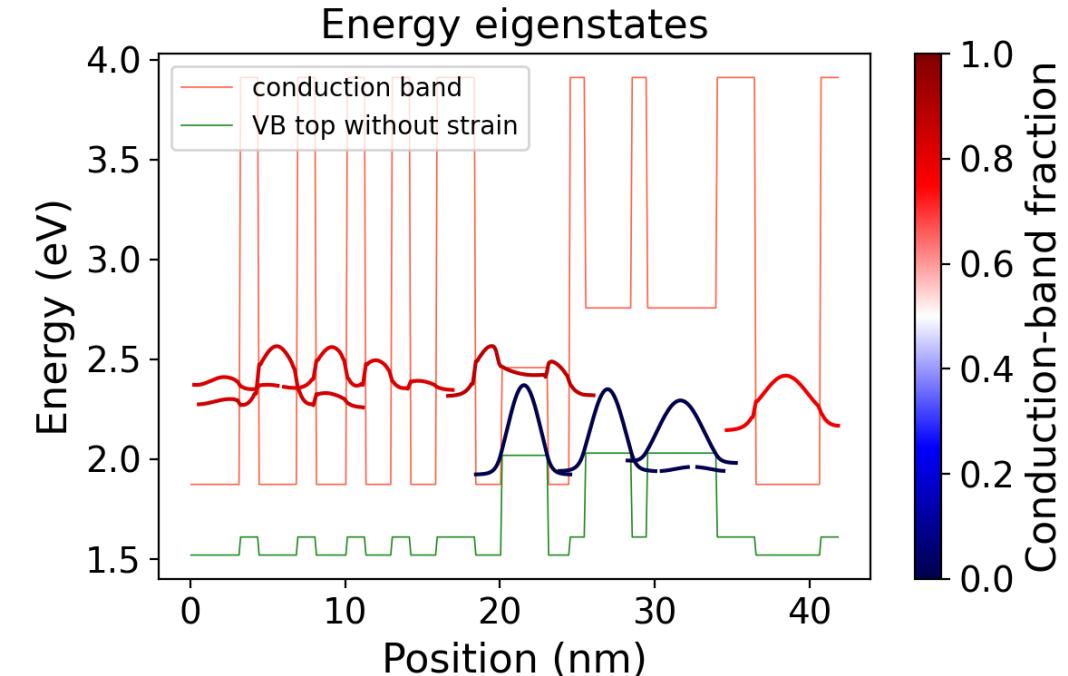


Tunneling charge transport

Accuracy

Reduce
problem size

Interband



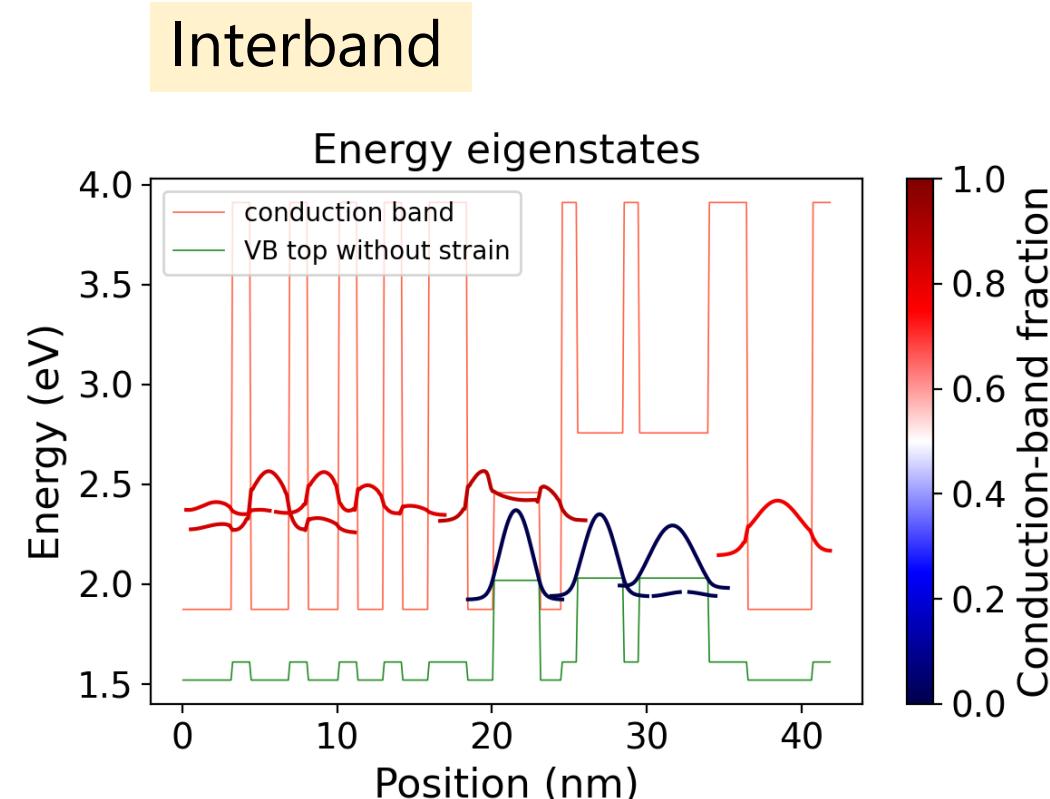
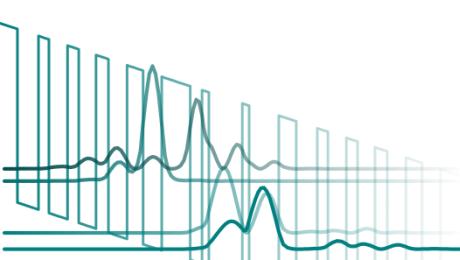
Tunneling charge transport

Accuracy

Reduce
problem size

$$G^<(\mathbf{r}_1, \mathbf{r}_2; t_2 - t_1)$$

- Solve 8-band $\mathbf{k}\cdot\mathbf{p}$ model
- Position eigenstates basis



Reduced Real Space basis

1. Reduce size by selecting relevant energy states

2. Diagonalize reduced position operator

$PH_0P = \left(\begin{array}{cc} & \diagdown \\ \diagup & \end{array} \right)$

$PzP = \left(\begin{array}{c} \text{green spring} \\ \diagdown \end{array} \right)$

$U^\dagger PH_0 P U = \left(\begin{array}{cc} & \text{blue double-headed arrow} \\ \text{blue double-headed arrow} & \end{array} \right)$

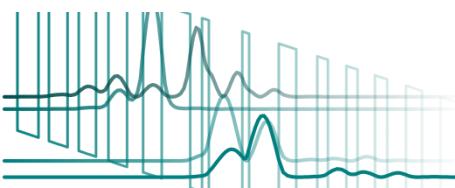
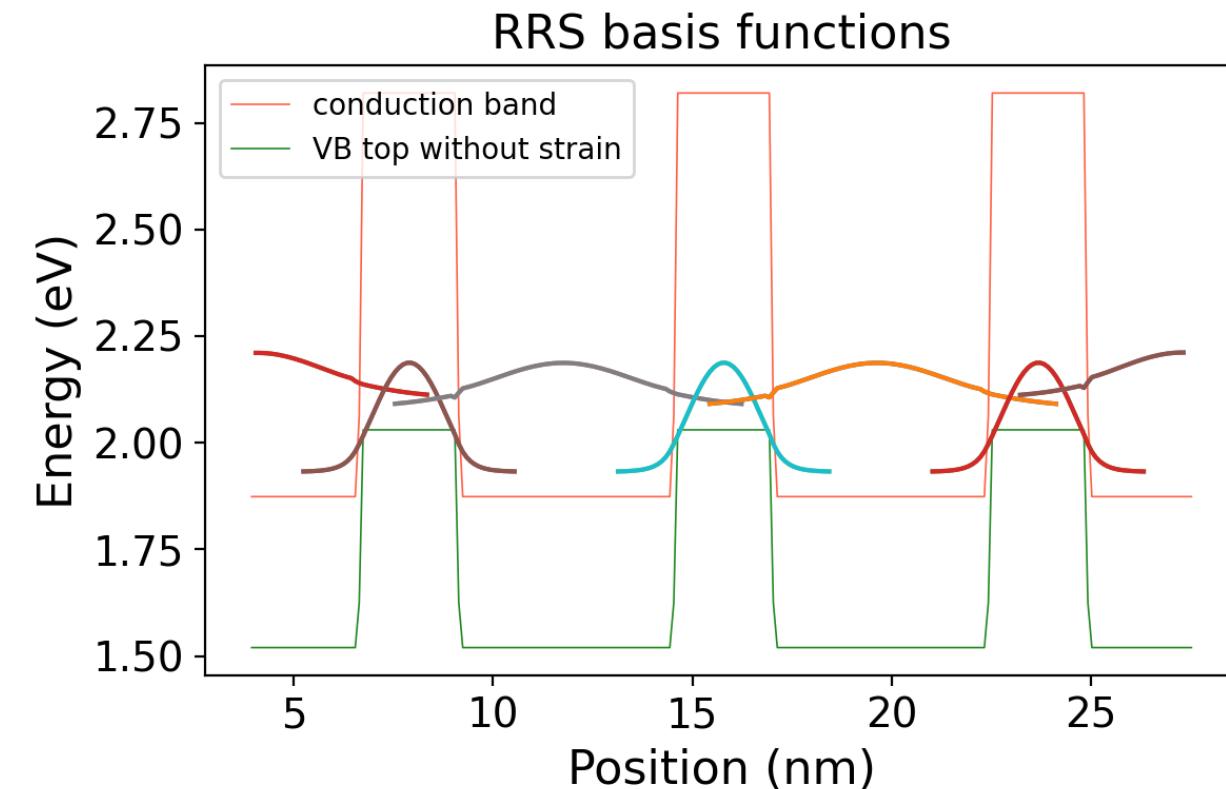
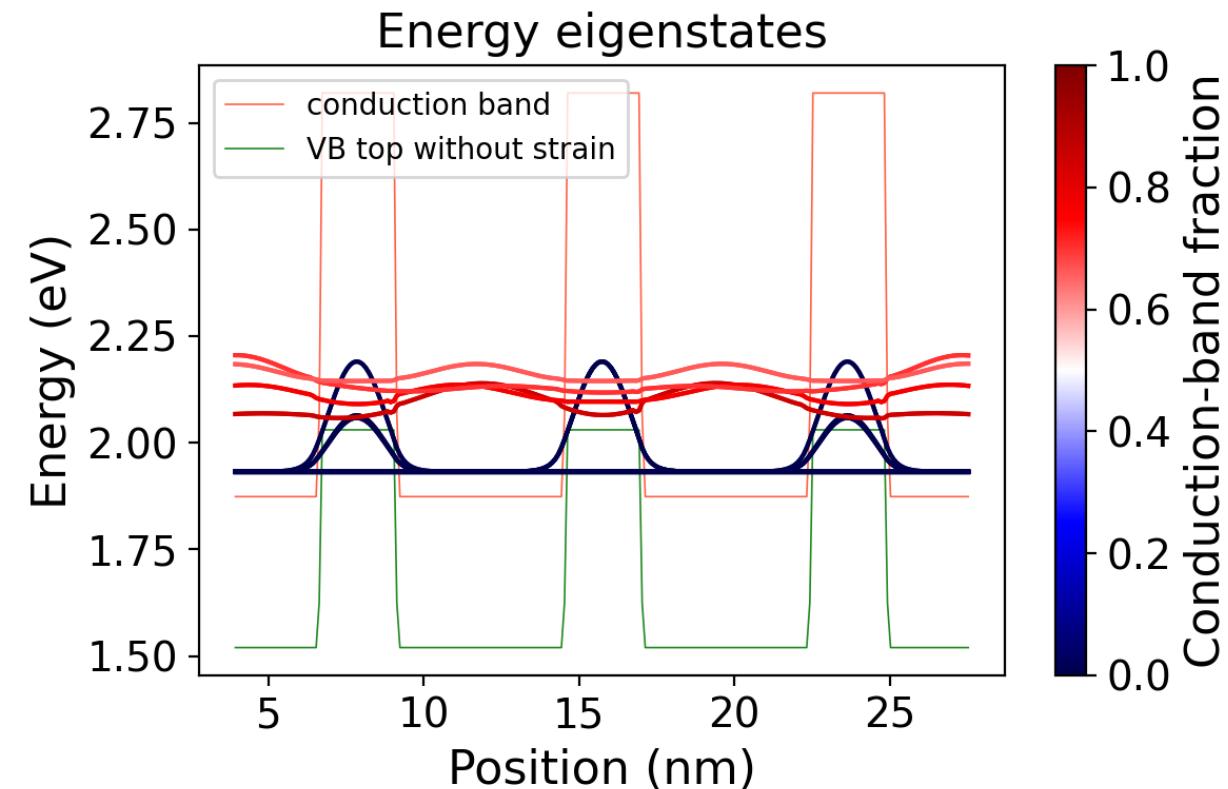
Coupling in RRS Hamiltonian

Position eigenbasis $|n\rangle$

Slide 14

Reduced Real Space basis

InAs/GaSb SLs



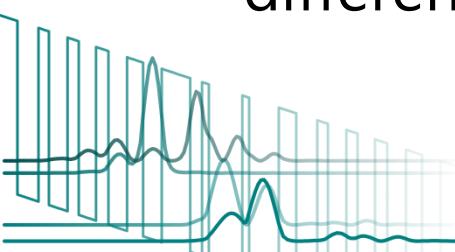
Basis set ready for NEGF

III-Sb broken-gap tunneling devices – photodetector, lasers

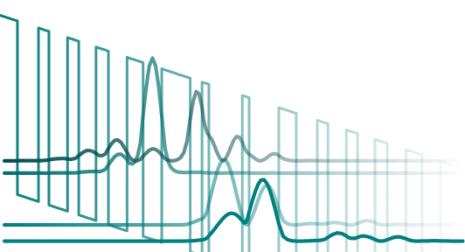
Accuracy

Reduce
problem size

- NEGF for non-equilibrium charge distribution
- 8-band $k \cdot p$ solutions to account for interband coupling & nonparabolicity
- Consistent discretization of differential operators
- Position eigenstates reflecting interband coupling

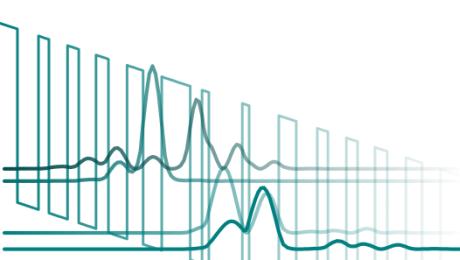
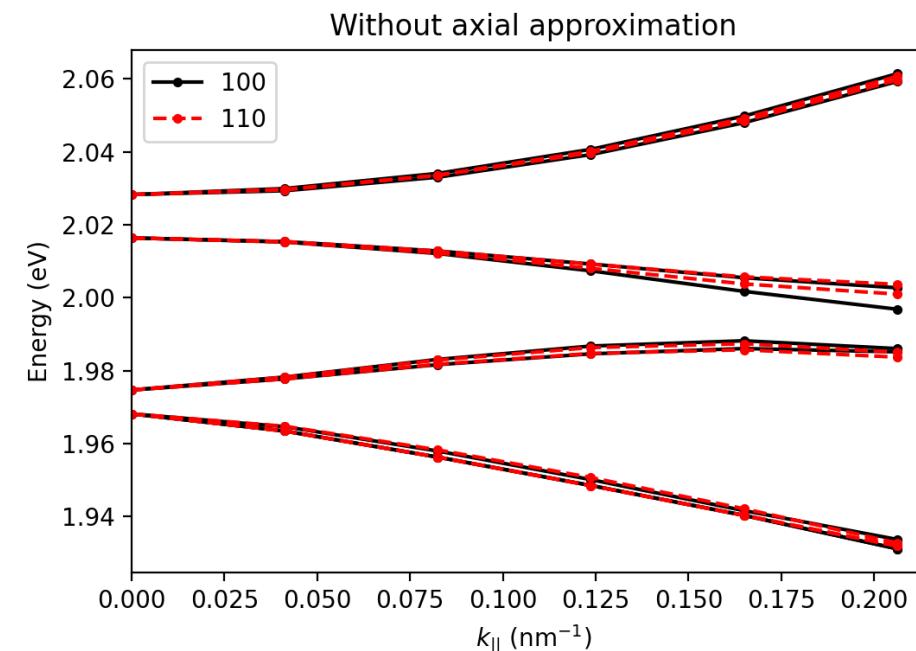
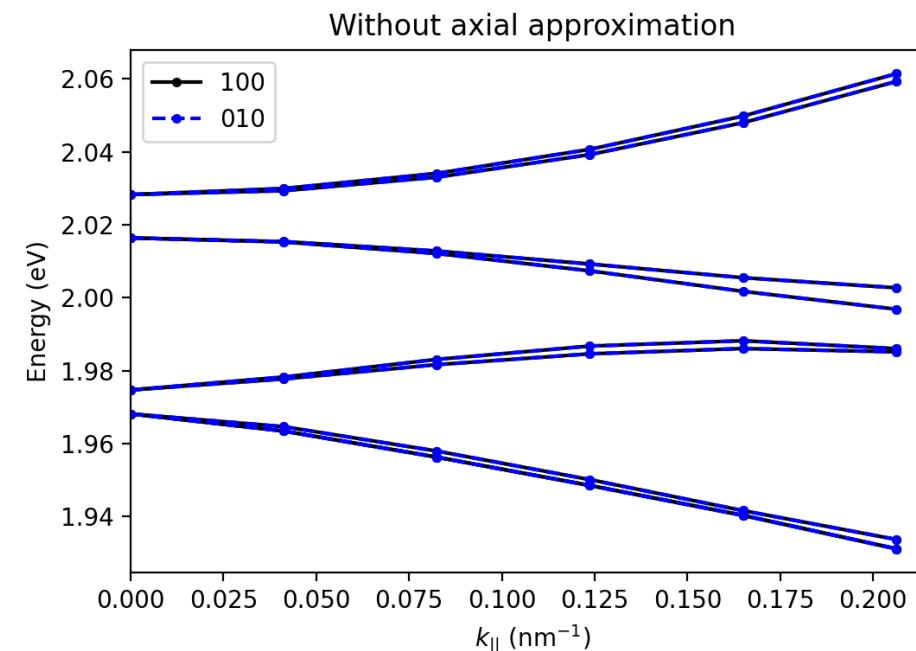


MSCA-ITN-2020 Funding Scheme from the European Union's Horizon 2020 programme under Grant agreement ID: 956548



In-plane anisotropy

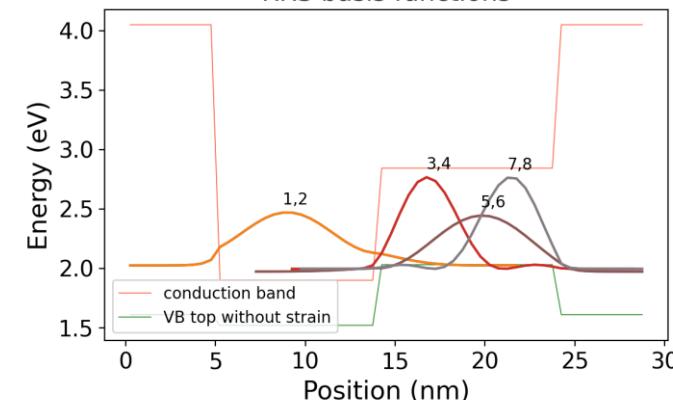
- Unlike nonparabolicity, in-plane anisotropy isn't significant.



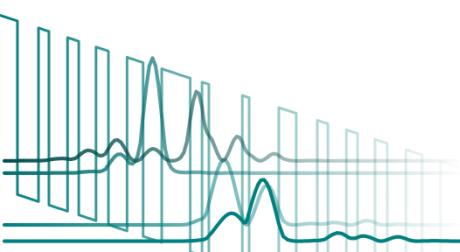
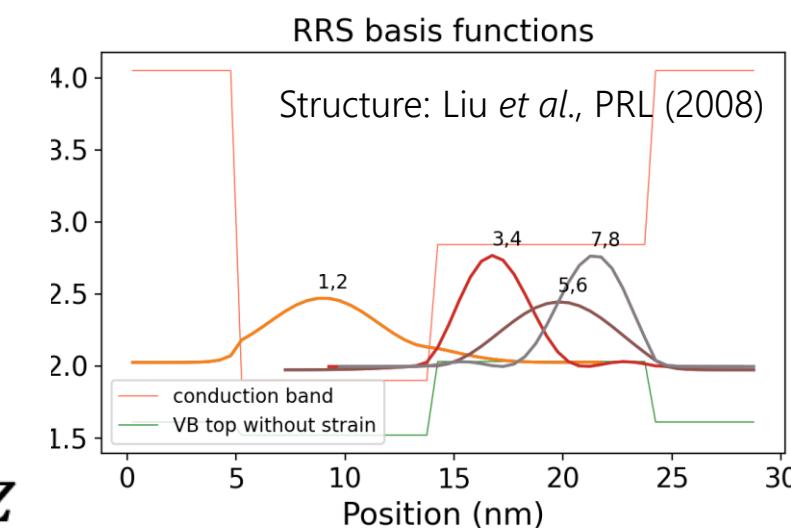
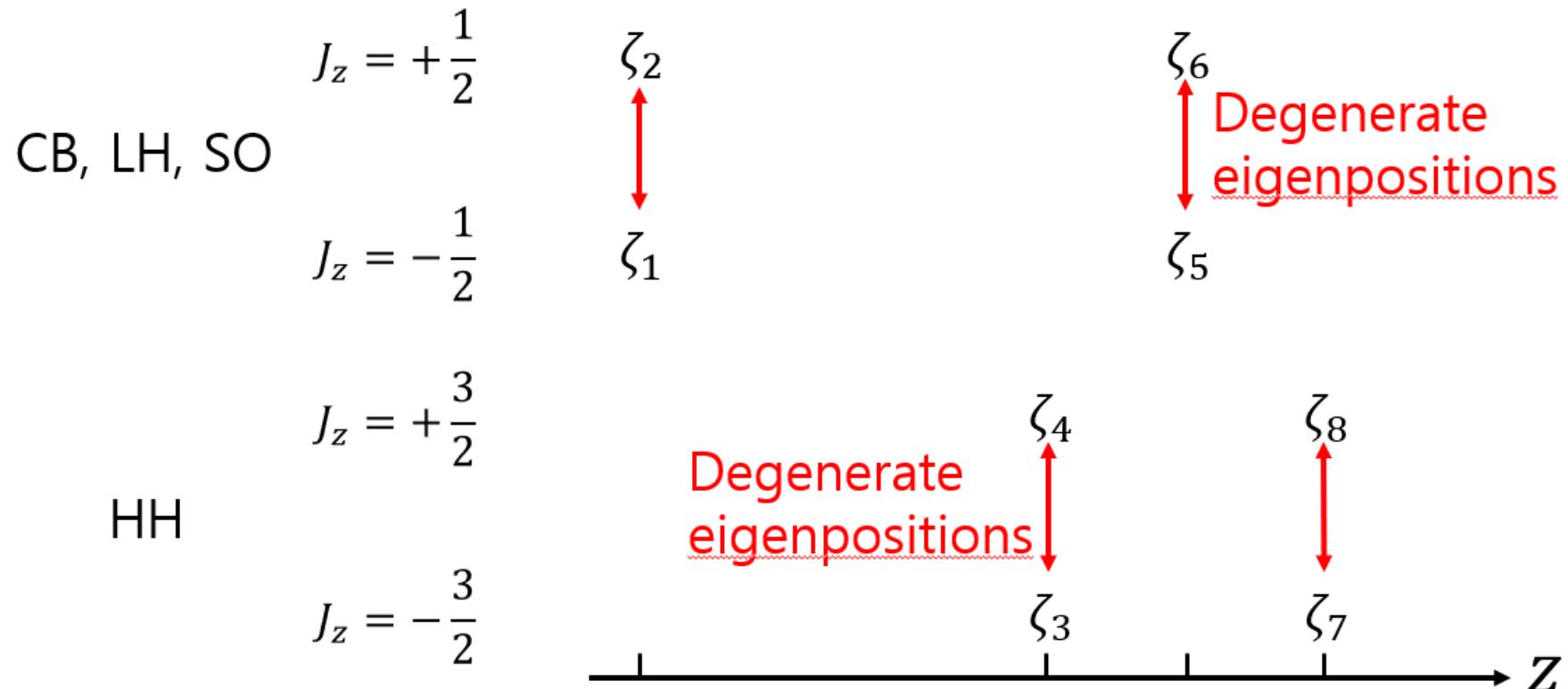
AlSb/InAs/GaSb/AlSb

Structure: Liu *et al.*, PRL (2008)

RRS basis functions



Degeneracy & coupling in Reduced Real Space



Degeneracy & coupling in Reduced Real Space

CB, LH, SO

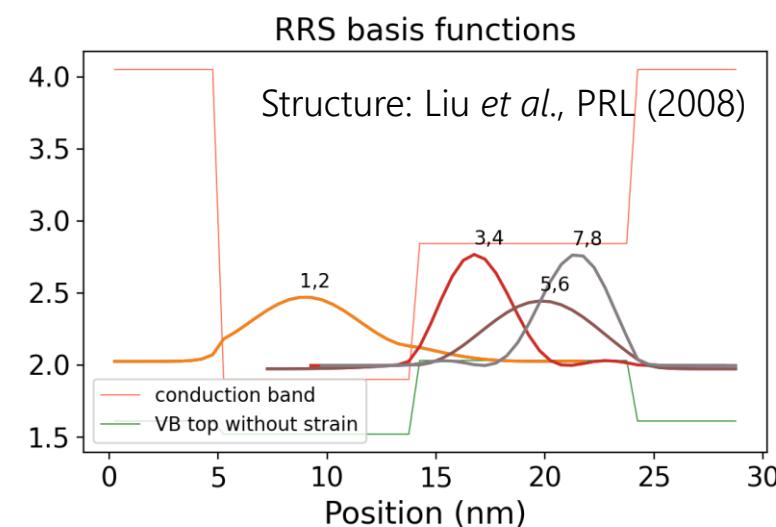
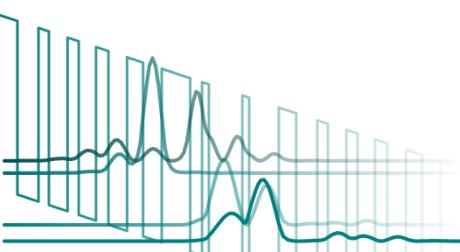
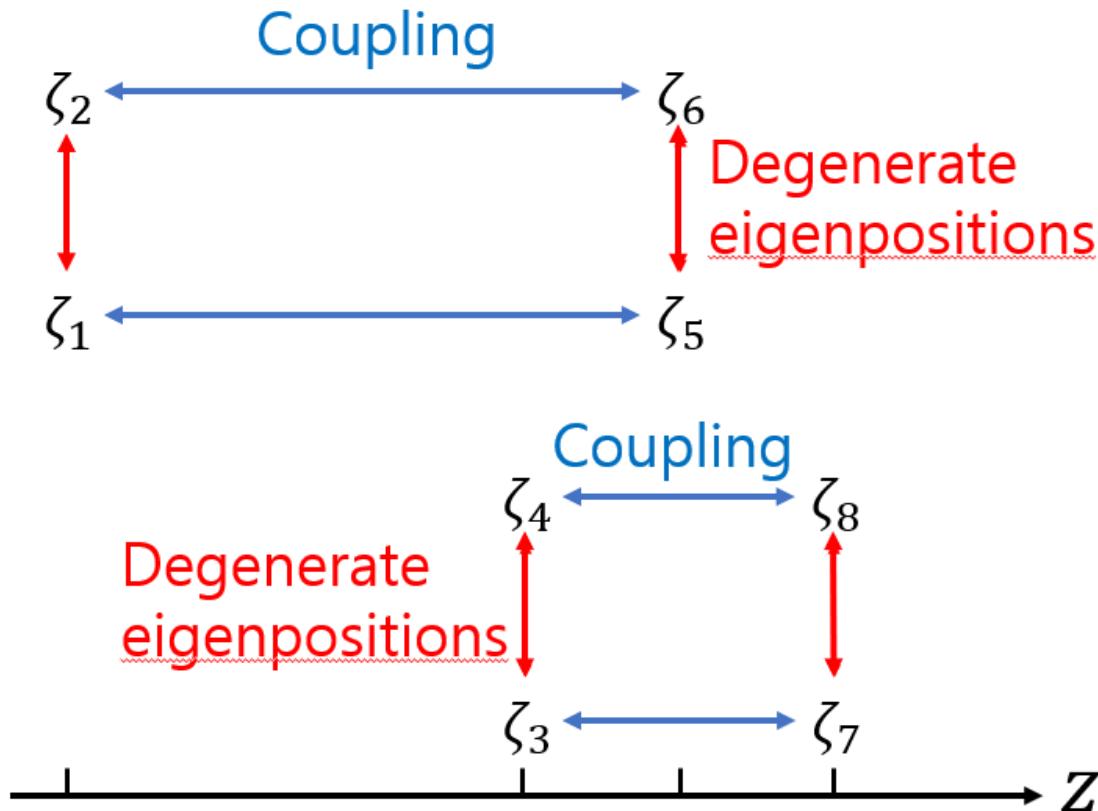
$$J_z = +\frac{1}{2}$$

$$J_z = -\frac{1}{2}$$

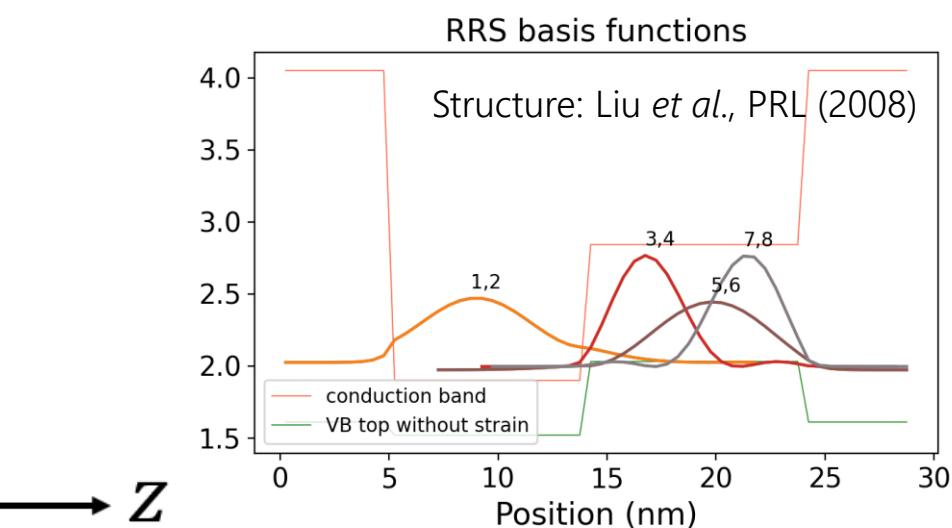
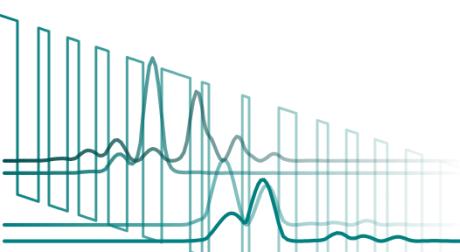
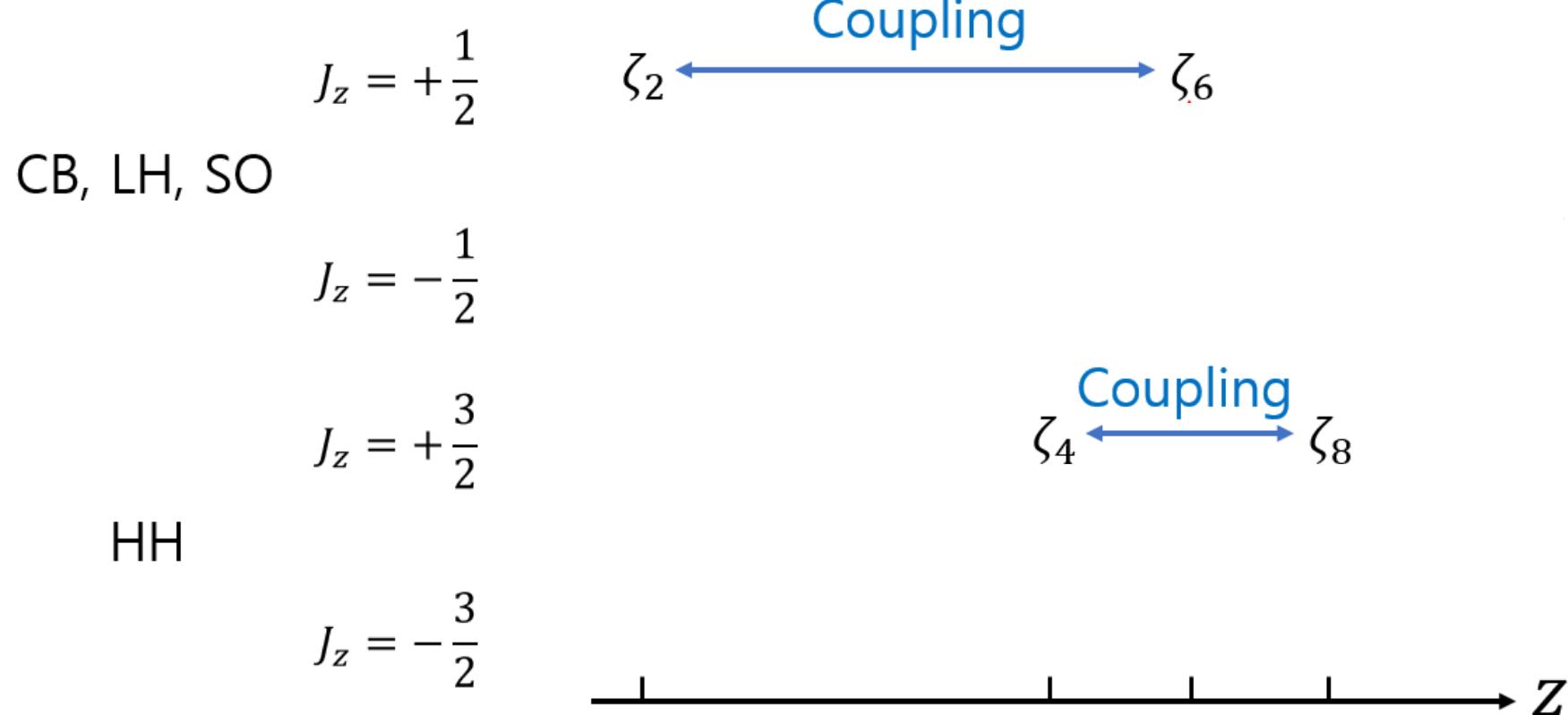
$$J_z = +\frac{3}{2}$$

HH

$$J_z = -\frac{3}{2}$$

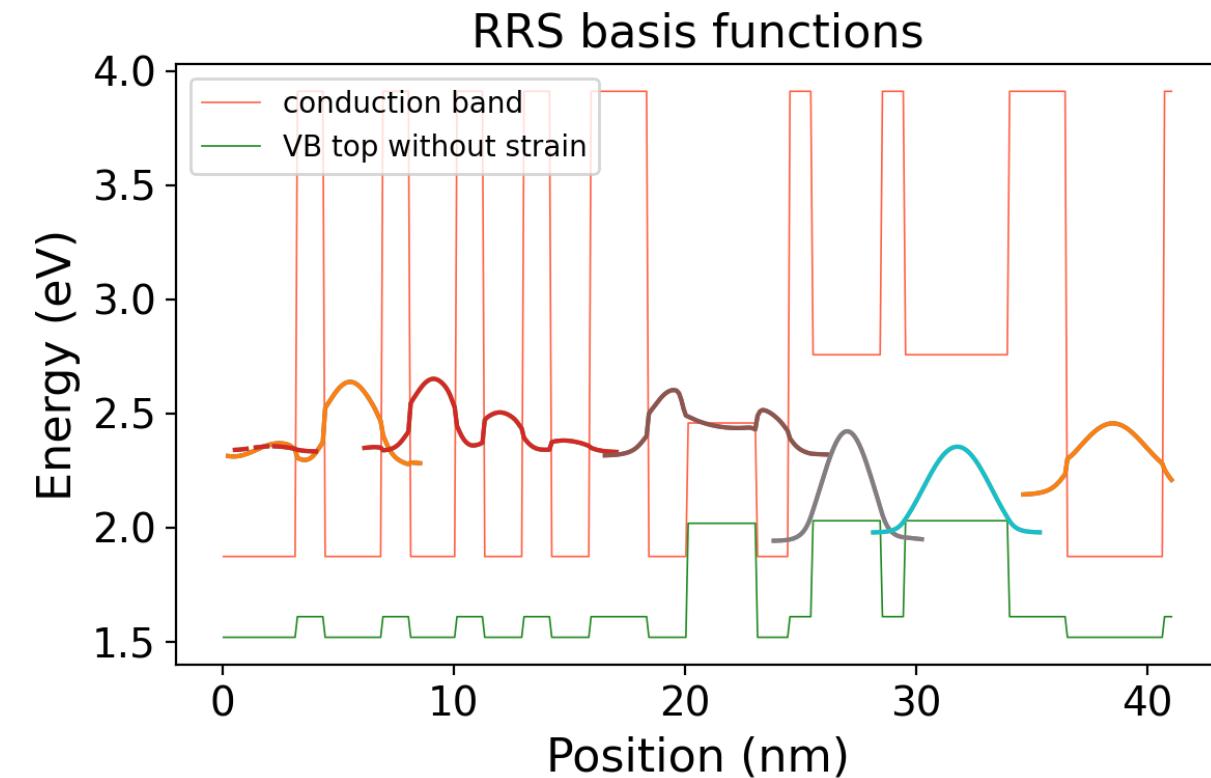
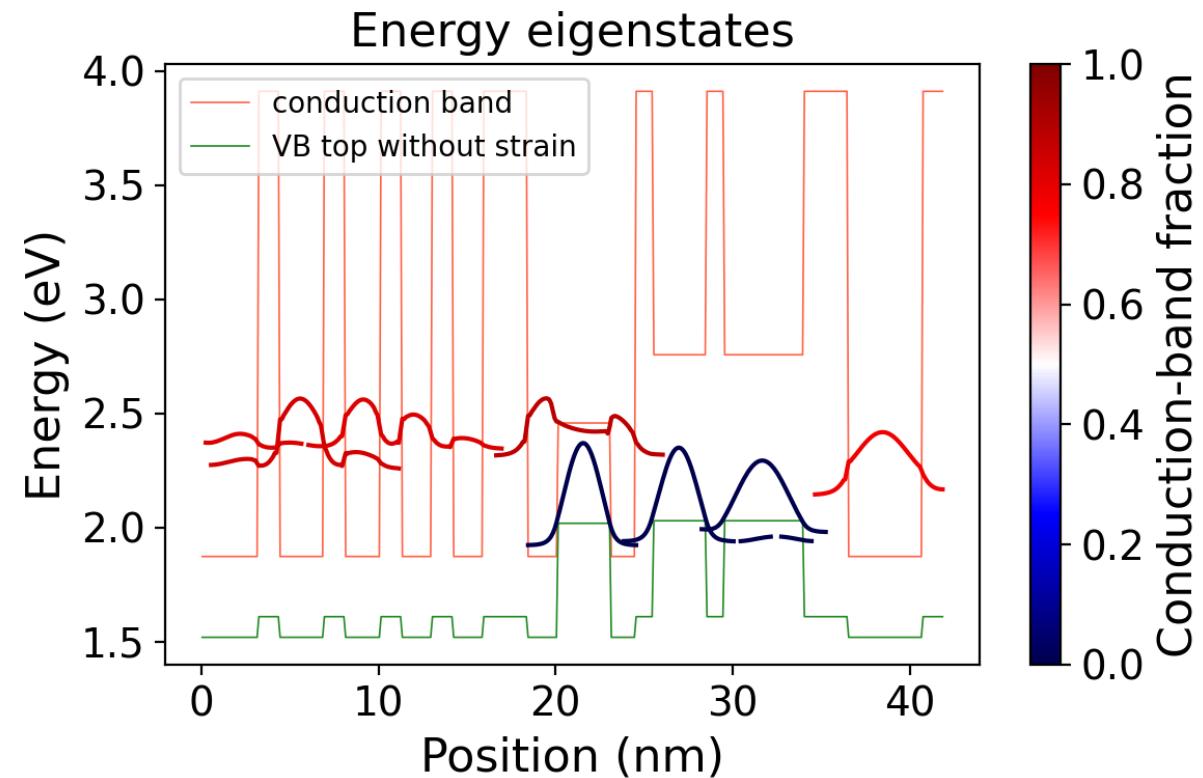


Degeneracy & coupling in Reduced Real Space

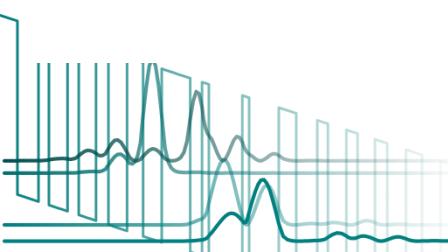


Reduced Real Space basis

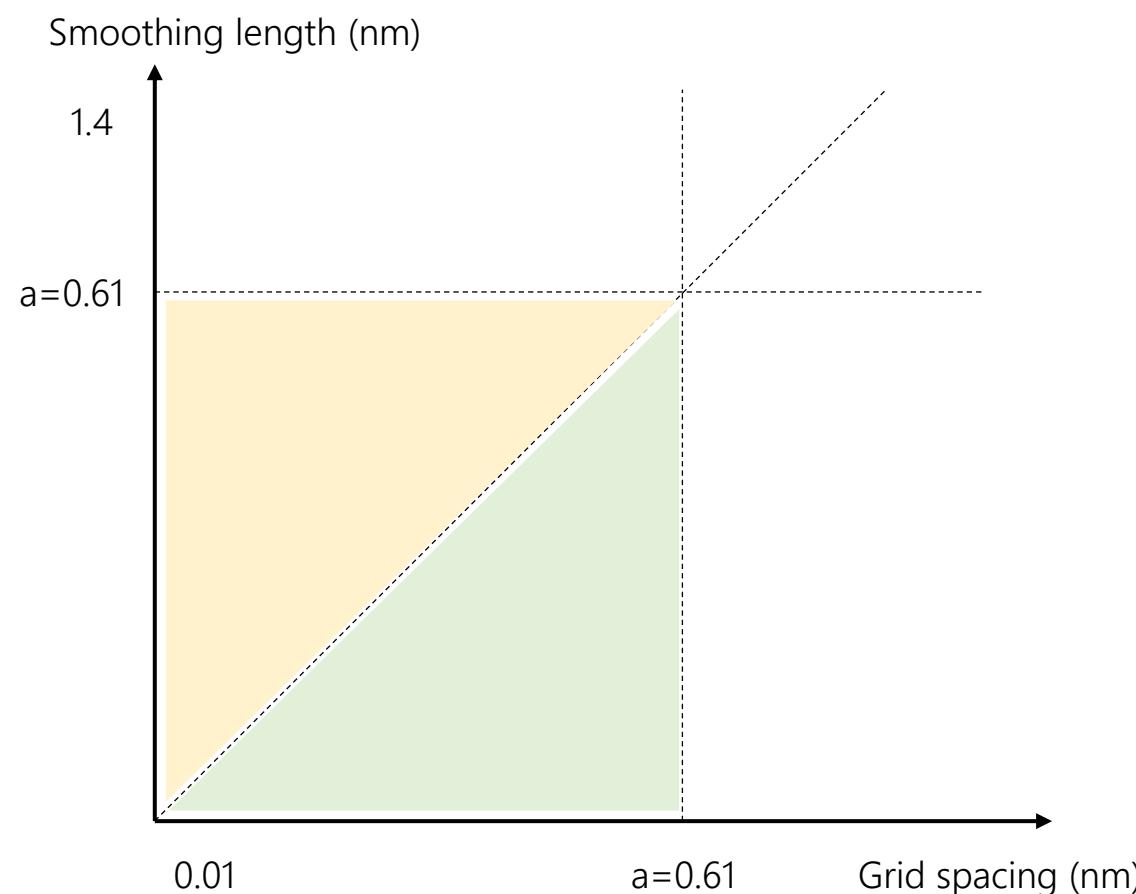
InAs/GaInSb ICLs



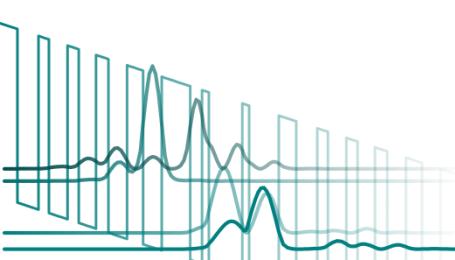
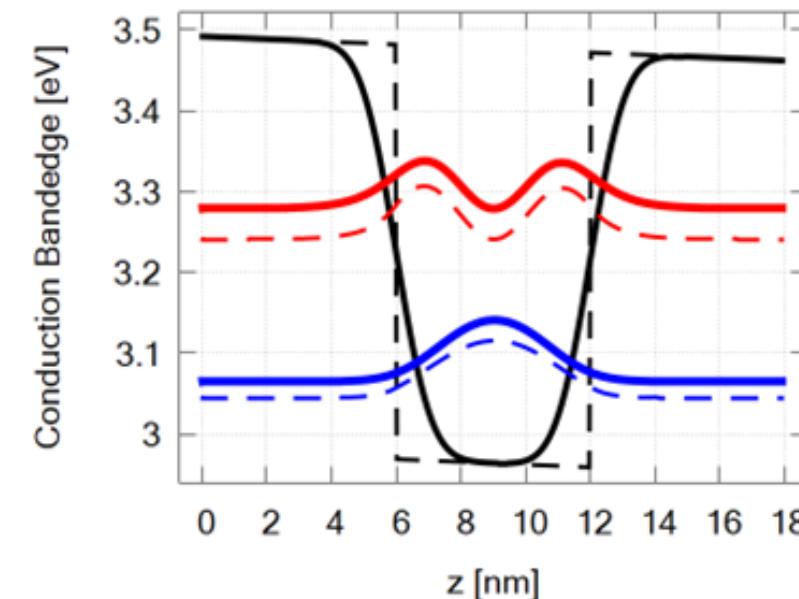
Structure: Vurgaftman *et al.*, Nat. Comm. (2012)



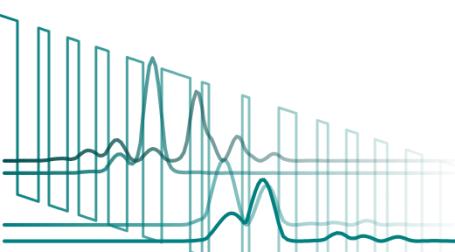
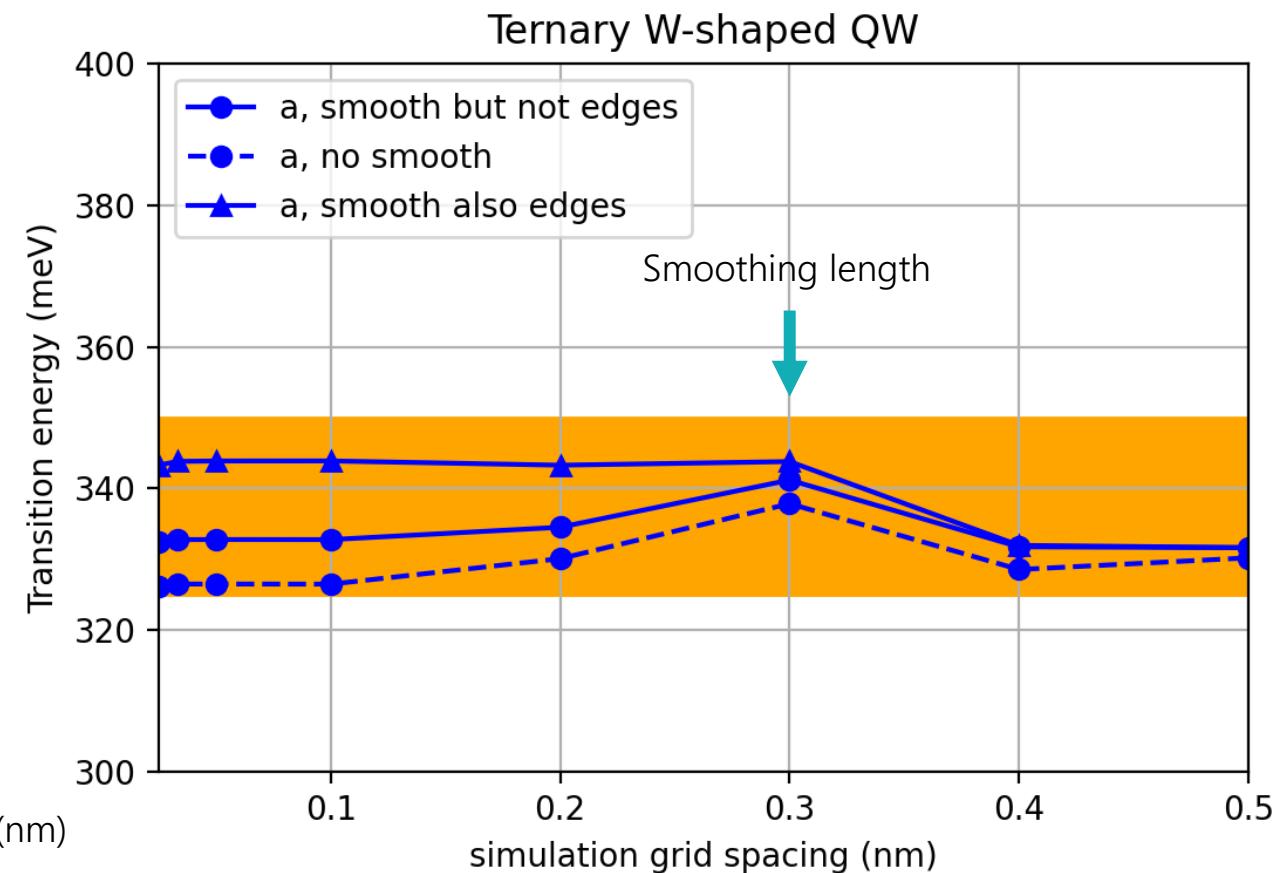
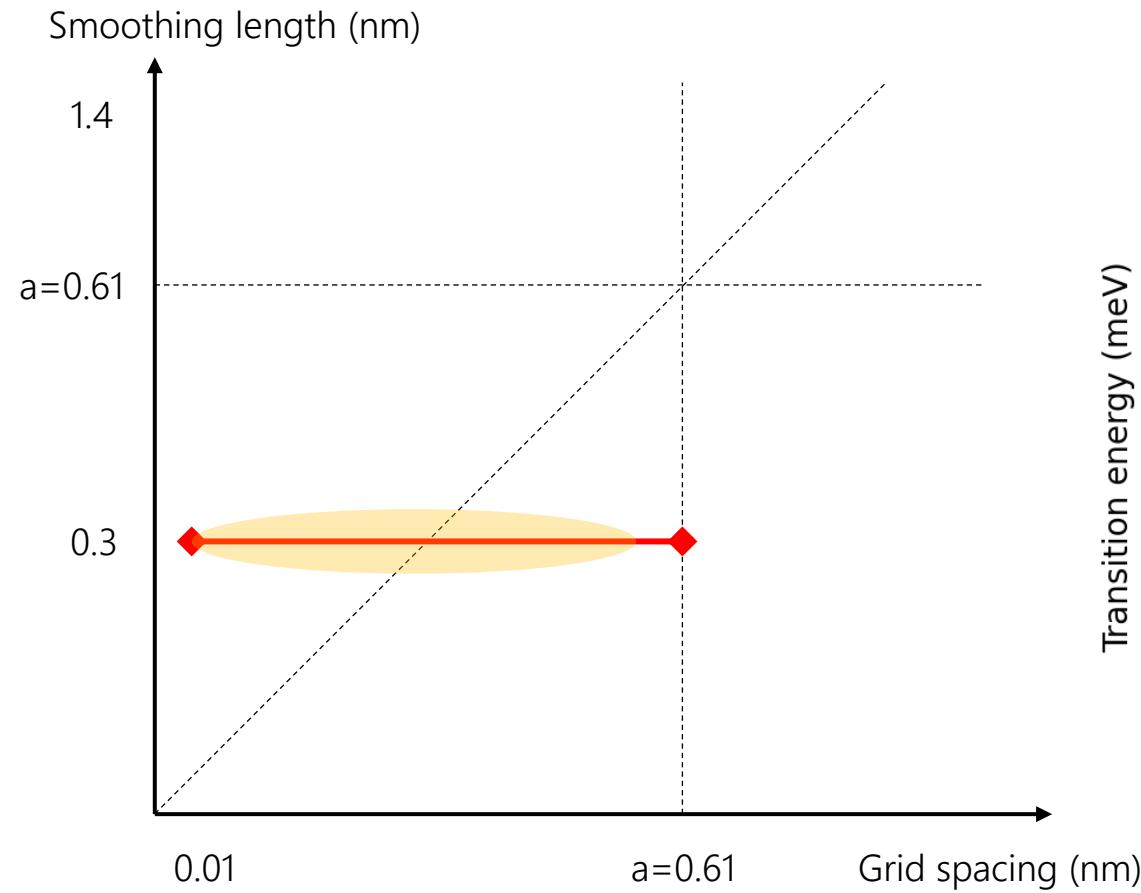
Transition energy in W-shaped QW



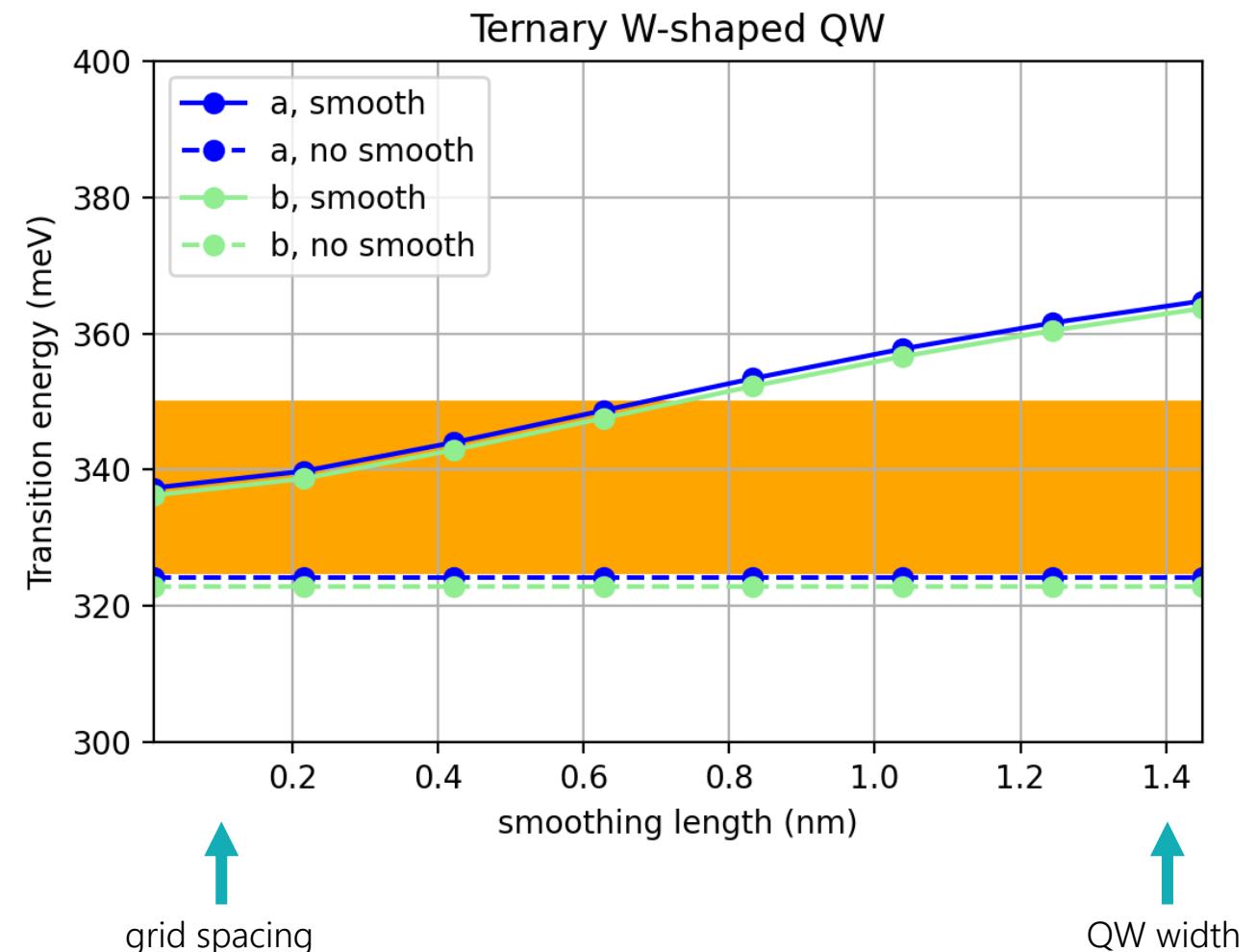
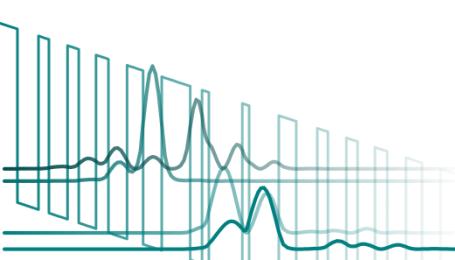
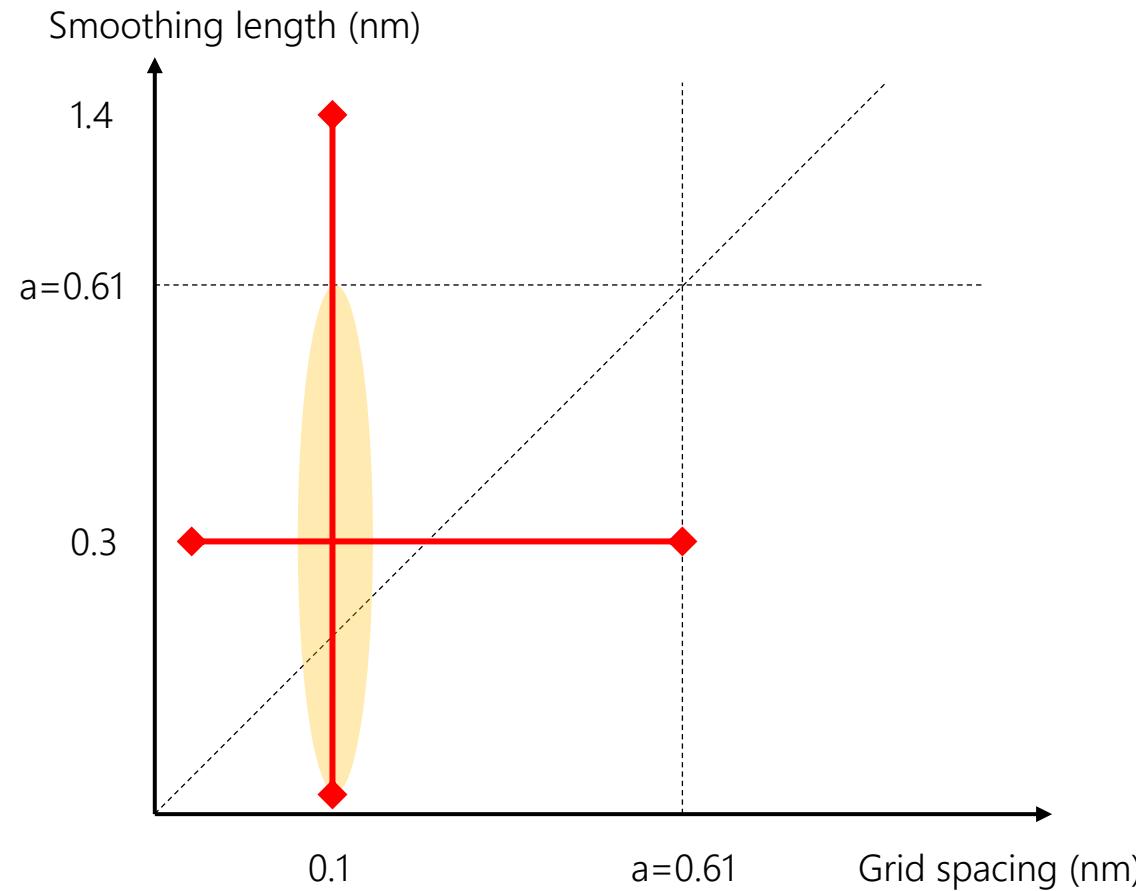
Smoothing of material parameters



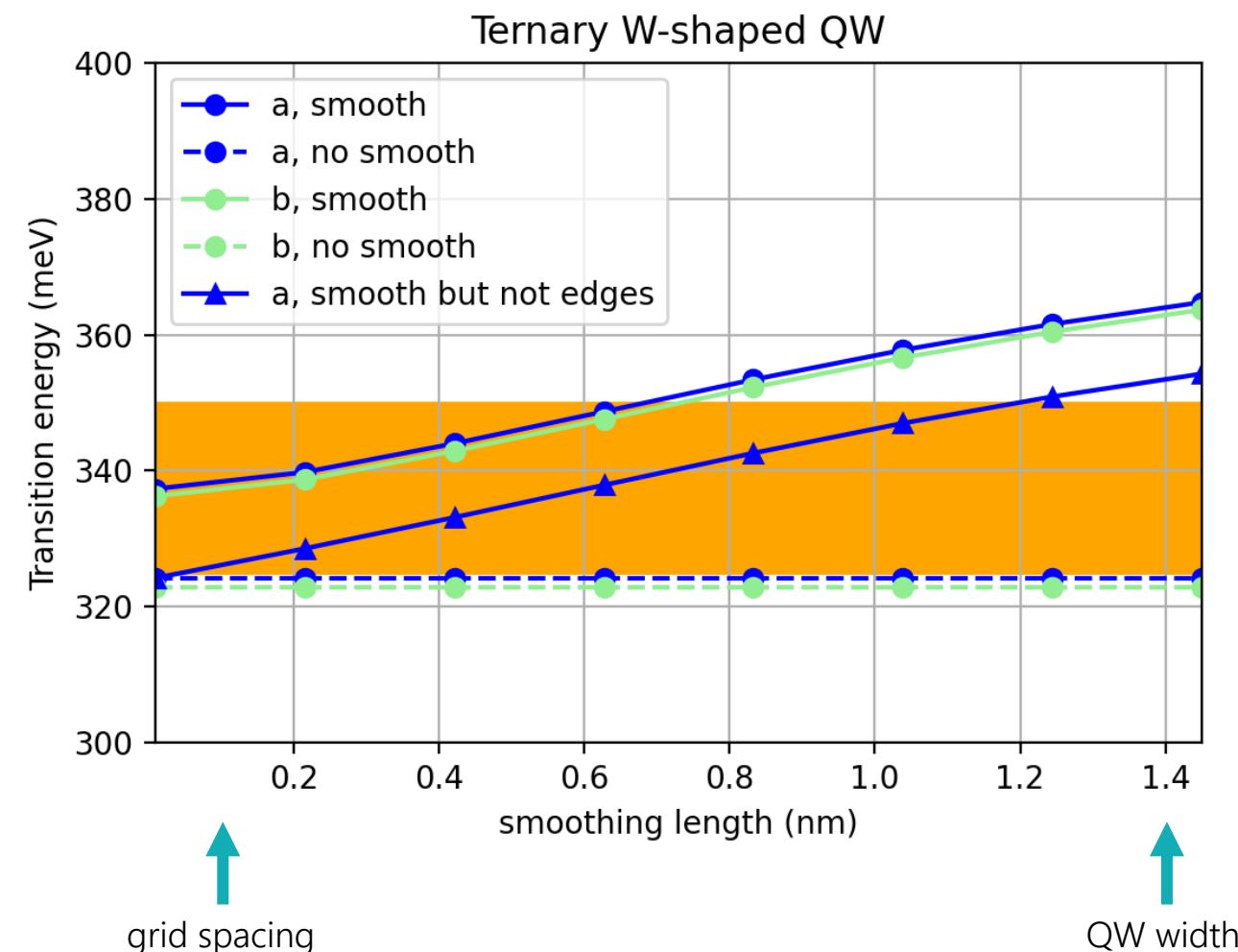
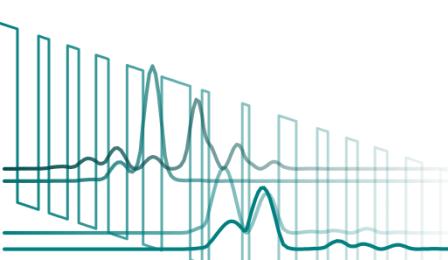
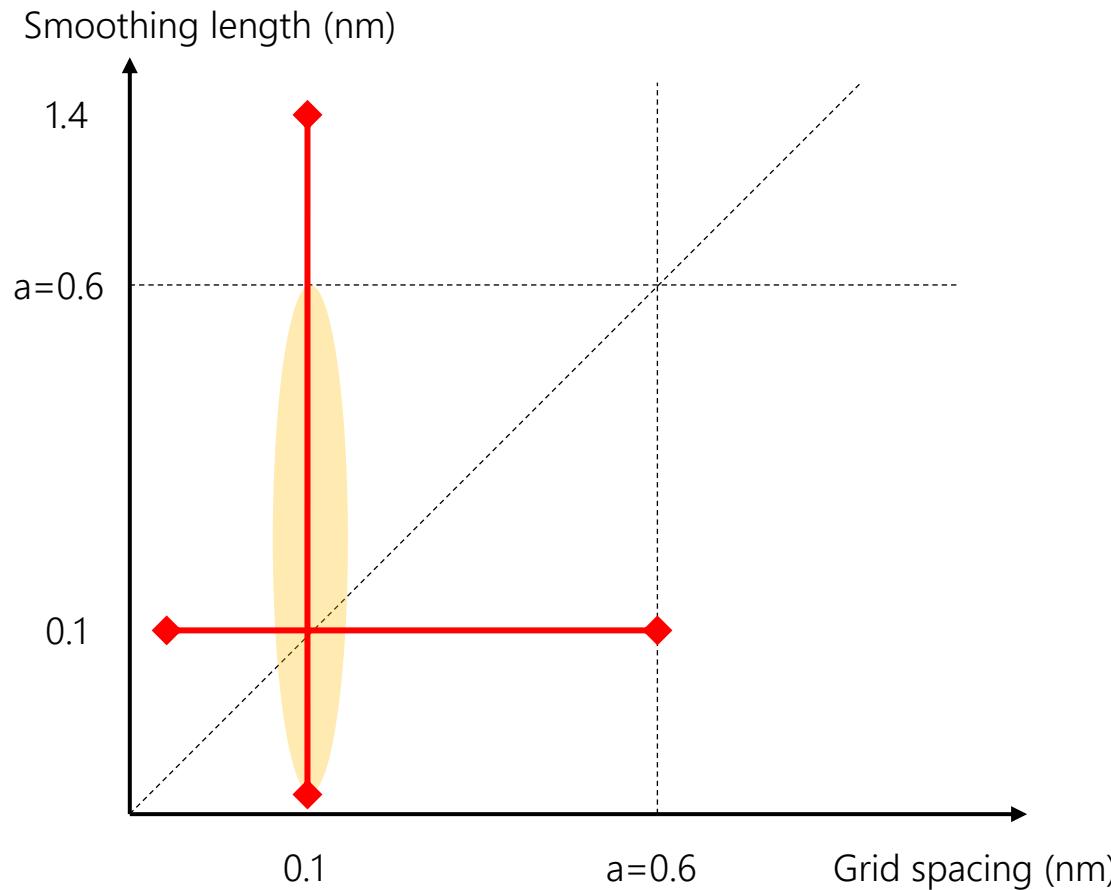
Transition energy in W-shaped QW



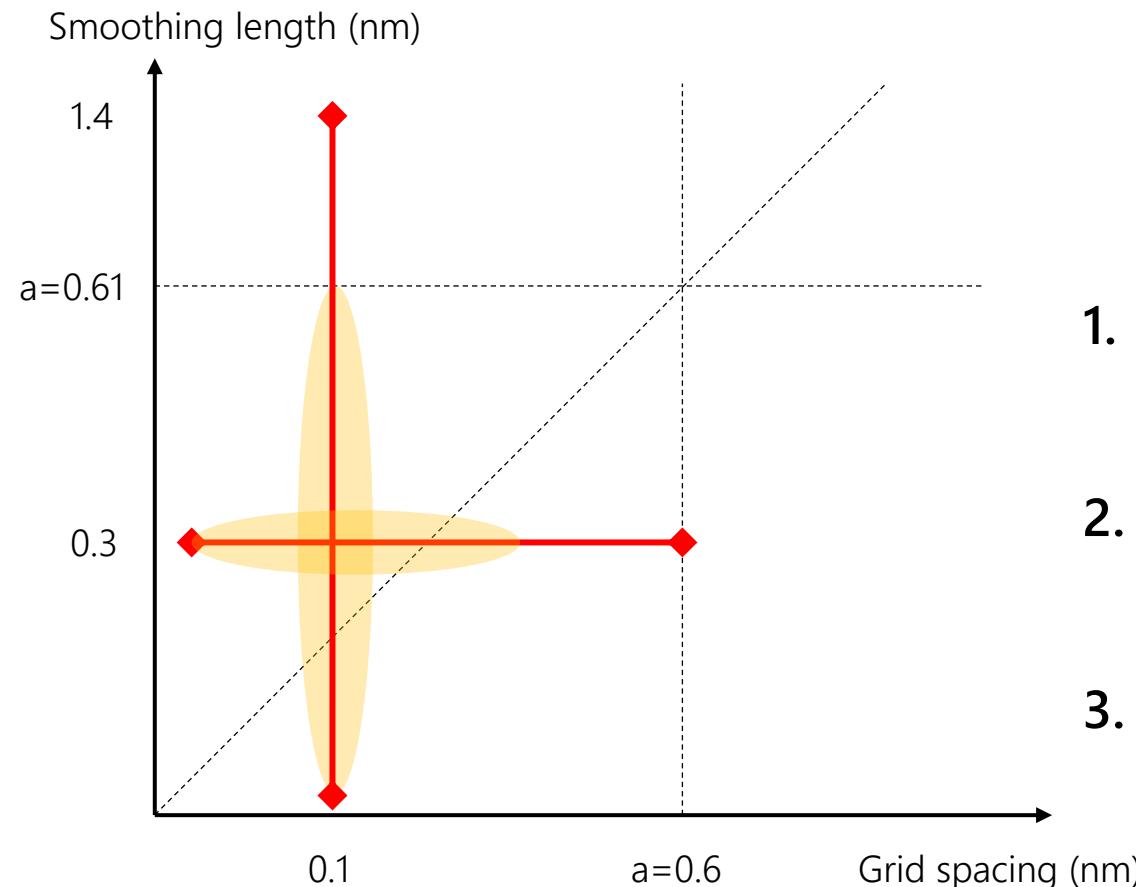
Transition energy in W-shaped QW



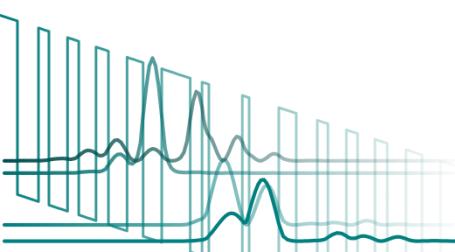
Validity of bandedge smoothing



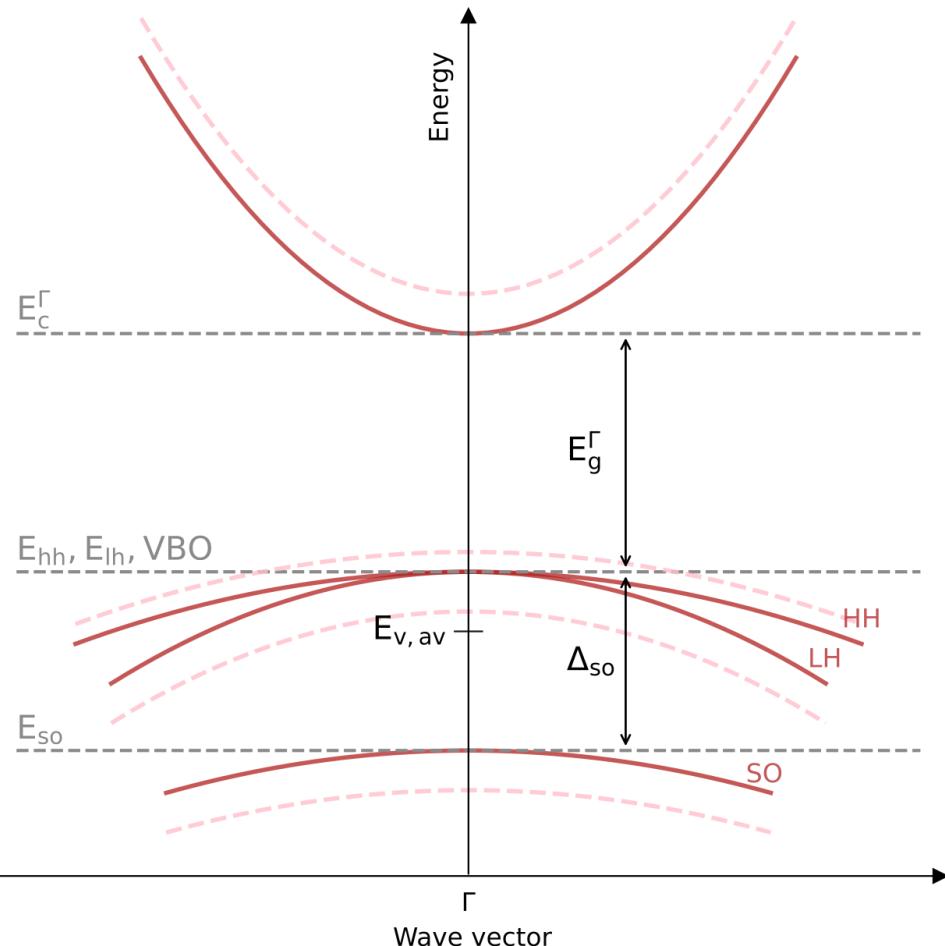
Transition energy in W-shaped QW



1. If kept < lattice const., numerical parameters do not change results significantly.
2. Choice of (a) or (b) does not affect results significantly if numerical parameters < lattice const.
3. Small smoothing seems necessary (and justifiable) to avoid jump of wave functions at the material interface.



Note: bandedge calculation



Option 1 (nn3, nn++, nnNEGFF):

$$E_c = E_{v,\text{av}} + E_g(T) + \Delta_{SO}/3$$

Option 2 (nnNEGFF):

$$E_{\text{hh},\text{lh}} = E_c - E_g(T)$$

$$E_{SO} = E_c - E_g(T) - \Delta_{SO}$$

Option 3 (nn3):

$E_{v,\text{av}}$ and E_c given in database

E_c in the database is for $T=0\text{K}$.
However, Varshni correction is applied to E_c in the code.