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Image-Force Barrier Lowering of Contact Resistance for Two-Dimensional Materials: Direct Determination and Method of Images on a Cone Manifold

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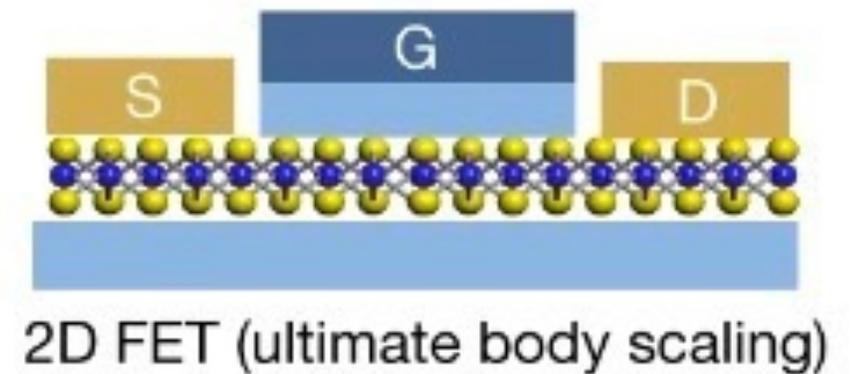
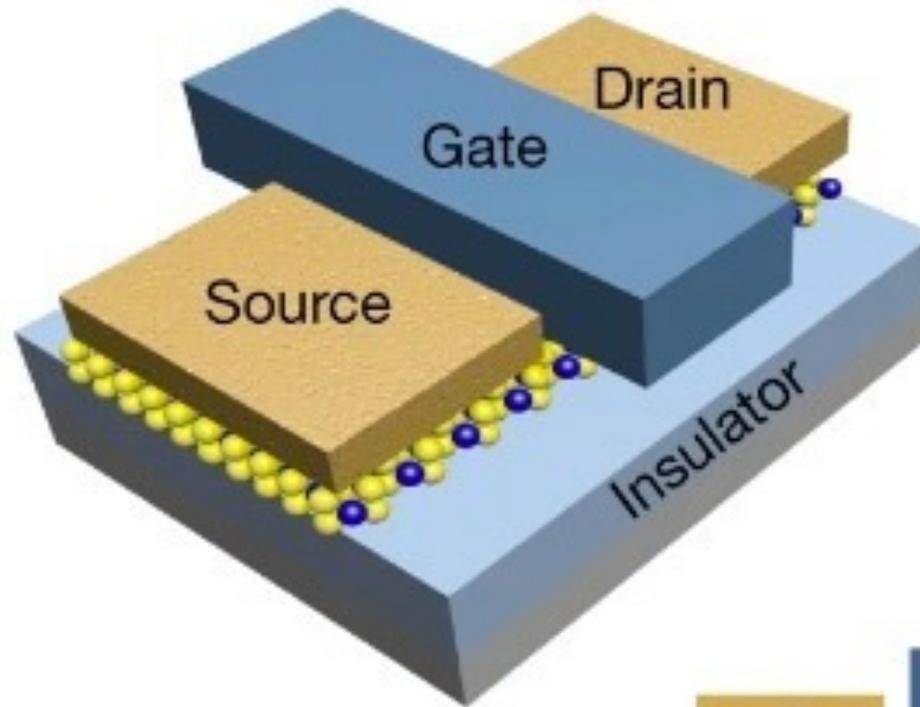
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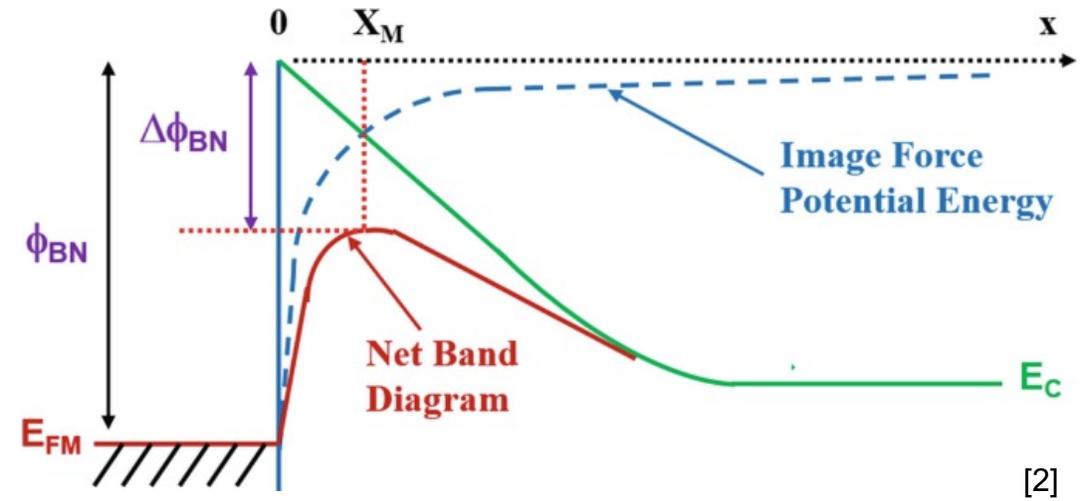
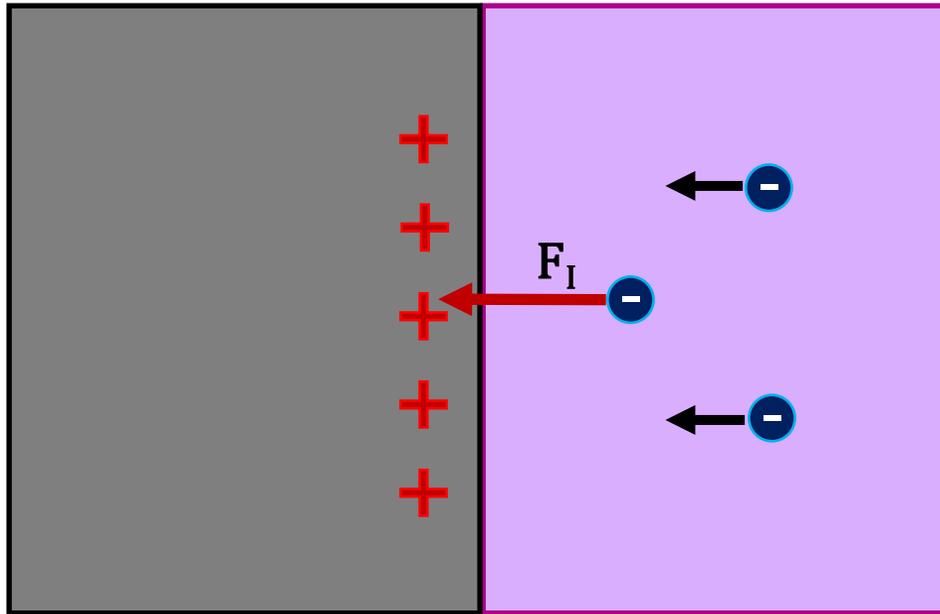
Introduction

- 2D Materials
 - TMD as channel material
- Limited by high:
 - Schottky barriers
 - Contact Resistance (R_C)



[1] Y. Liu, *et al.*, "Promises and prospects of two-dimensional transistors," *Nature*, 2021

Image Force Barrier Lowering (IFBL)



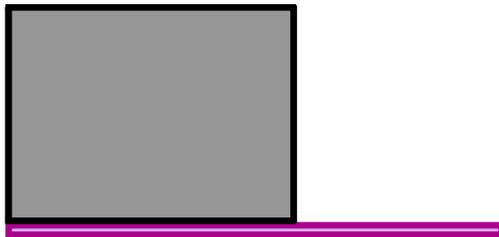
[2]

[2] Baliga, B.J., *Springer* (2019).

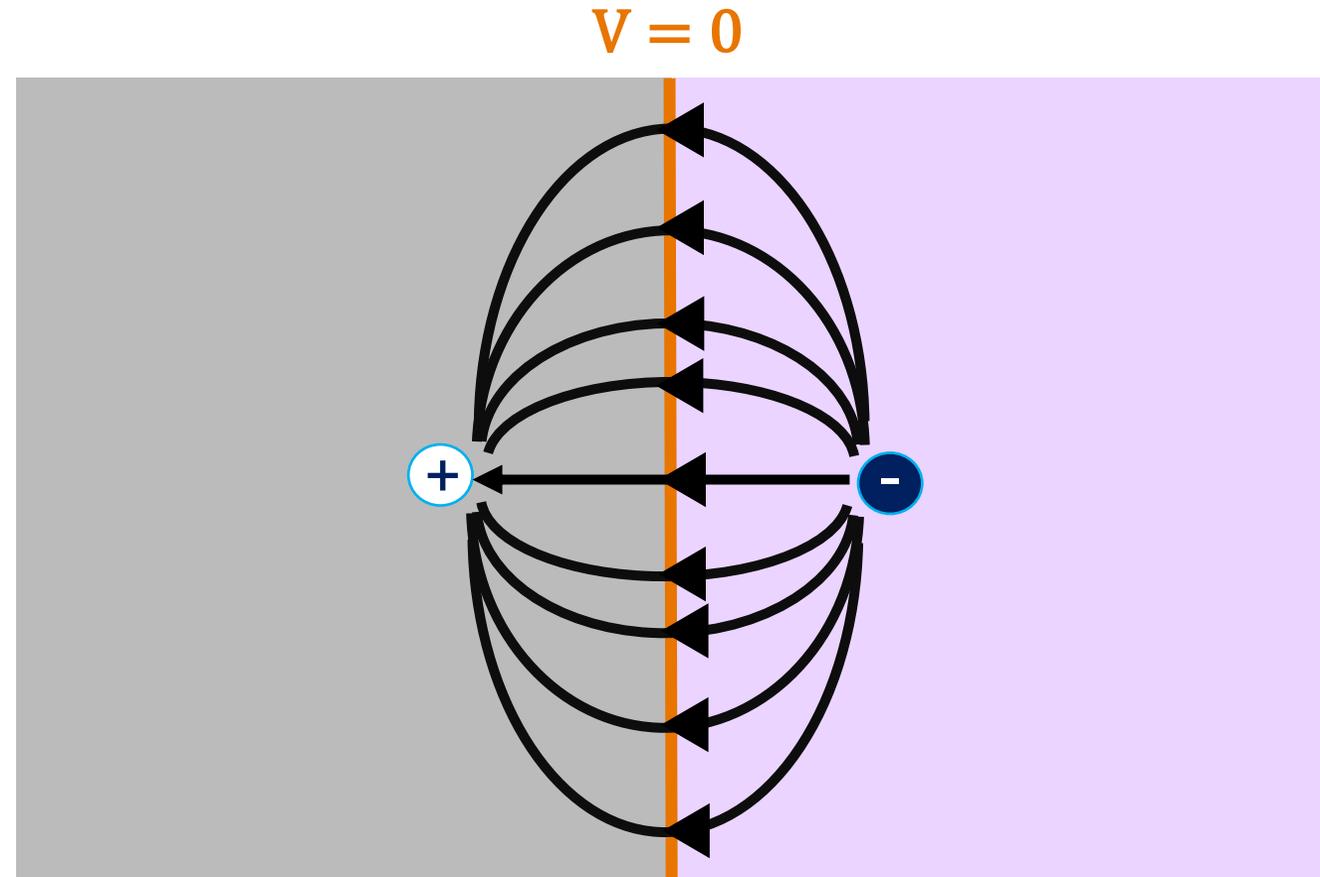
How can we solve for IFBL?

- Method of Images!

For asymmetrical geometry?

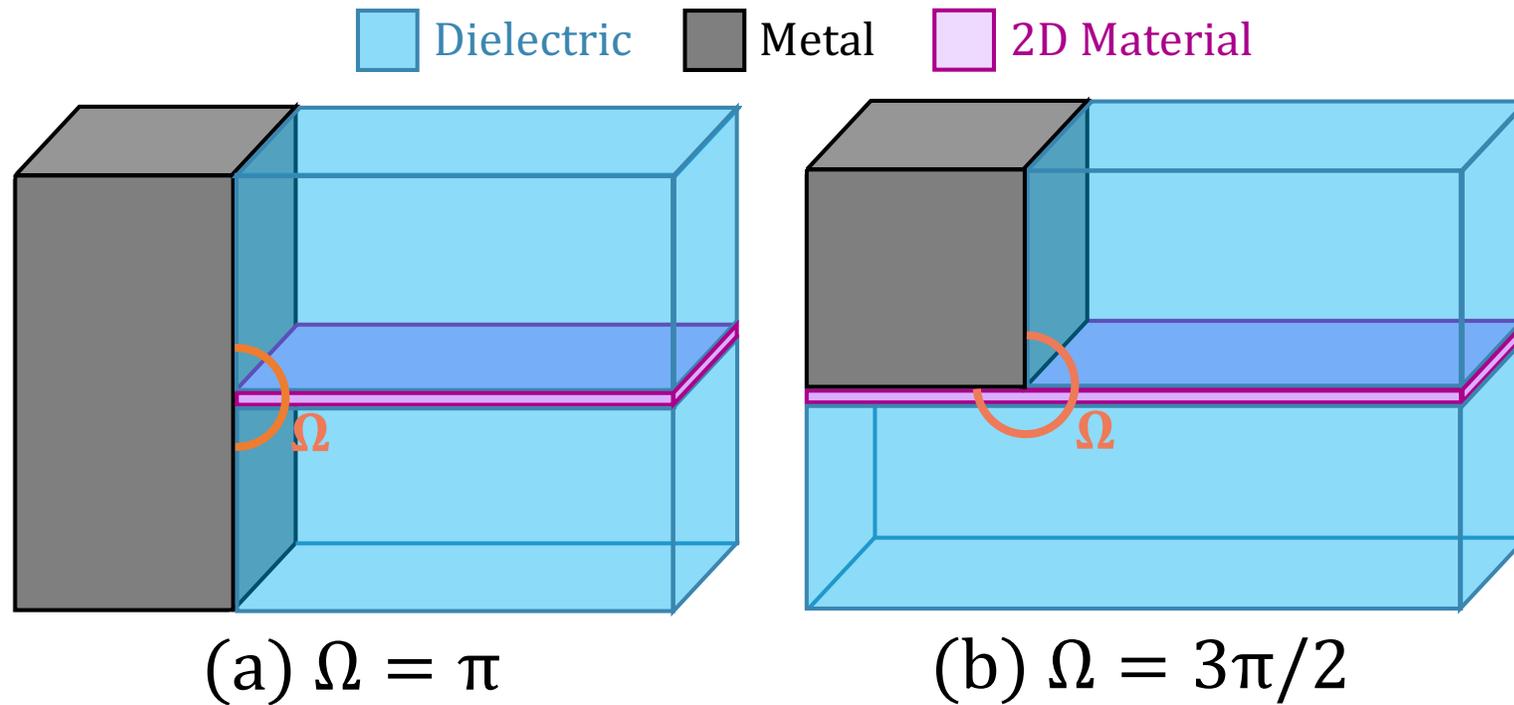


- ~~Method of images~~
- Poisson's equation



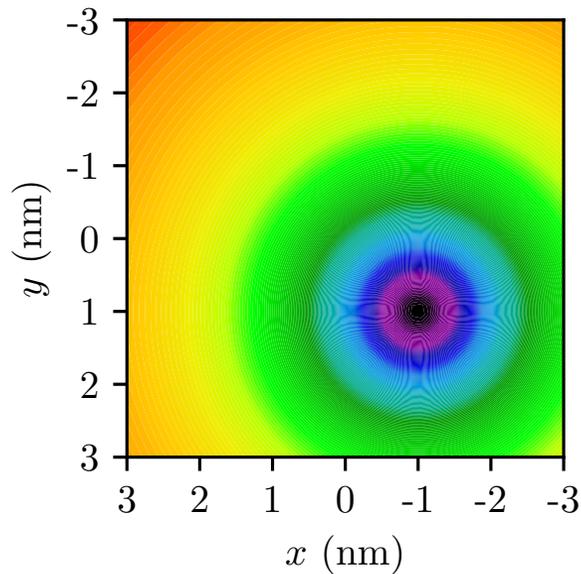
The Problem:

- IFBL necessary to find R_C
- No solution for geometries like (b)

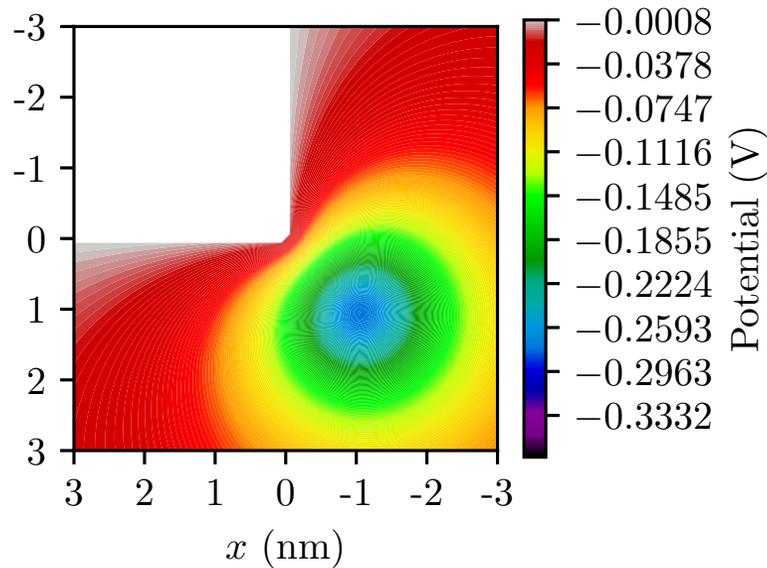


Isolating The Image Potential (V_I)

[3]

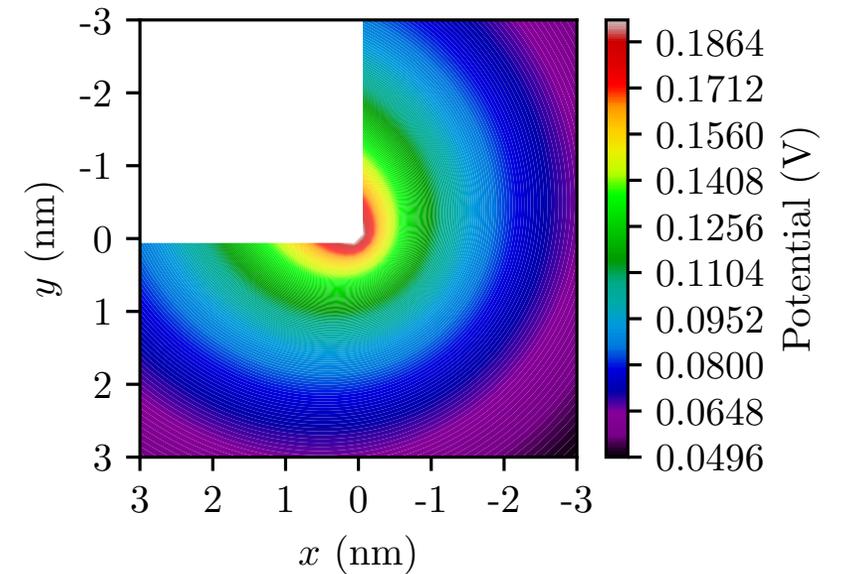


V_C



V

$$V_I = V - V_C$$



V_I

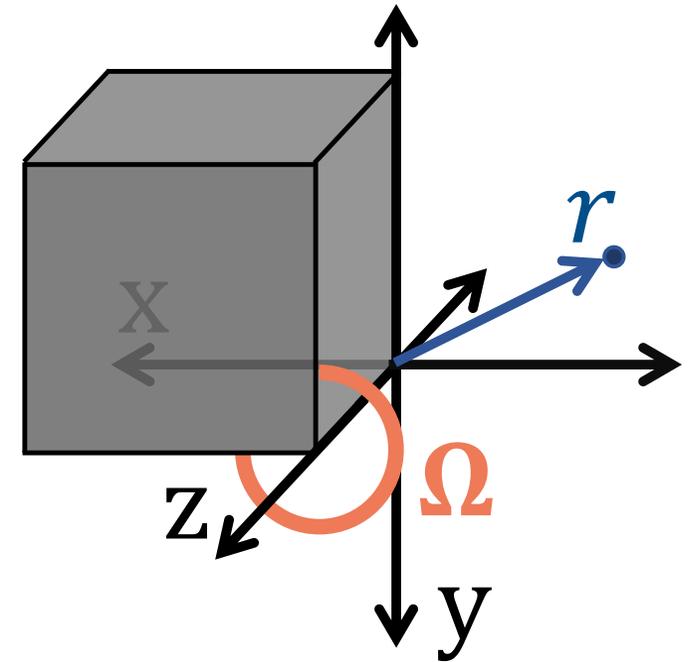
[3] S. Evans, *et al.*, "Image-Force Barrier Lowering for Two-Dimensional Materials: Direct Determination and Method of Images on a Cone Manifold," arXiv, 2023

Determining U_{IFBL} from V_I

To find U_{IFBL} , we solve:

$$U_{\text{IFBL}}(r, \theta) = - \int_0^{-e} \frac{q}{e} V_I(r_0, \theta_0, z_0; r, \theta, z) dq = -\frac{1}{2} e V_I$$

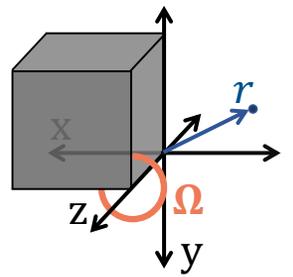
Found from
 $V_I = V - V_C$



q = charge of electron
 e = charge of the electron
inducing U_{IFBL}

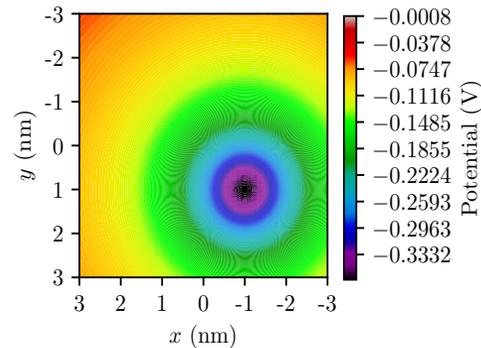
Isolating V_I

- Find V with Poisson's Equation: $V(r, \theta, z) = \frac{e}{\epsilon} \int_0^\infty dk_z Z_{k_z}(z) \int_0^\infty d\alpha C_{\alpha, k_z} \Theta_\alpha(\theta) R_{\alpha k_z}(r)$
 - Solve $\Theta(\theta)$ with 2 sets of boundary conditions:



1. Point charge in free space (V_C)

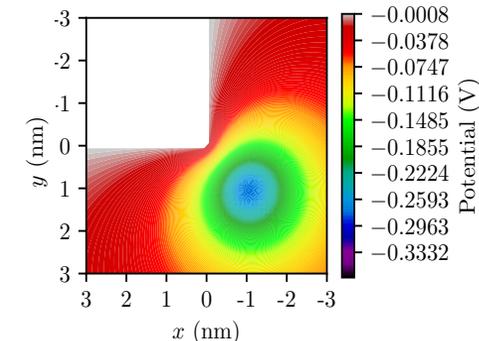
$$\Theta_{C;\alpha}(\theta; \theta_0) = \frac{\cosh(\alpha(\pi - |\theta - \theta_0|))}{\alpha \sinh(\alpha\pi)}$$



2. Point charge near a metal wedge (V)

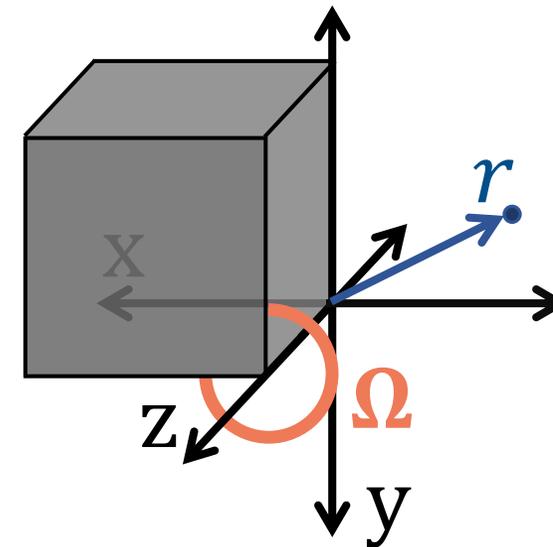
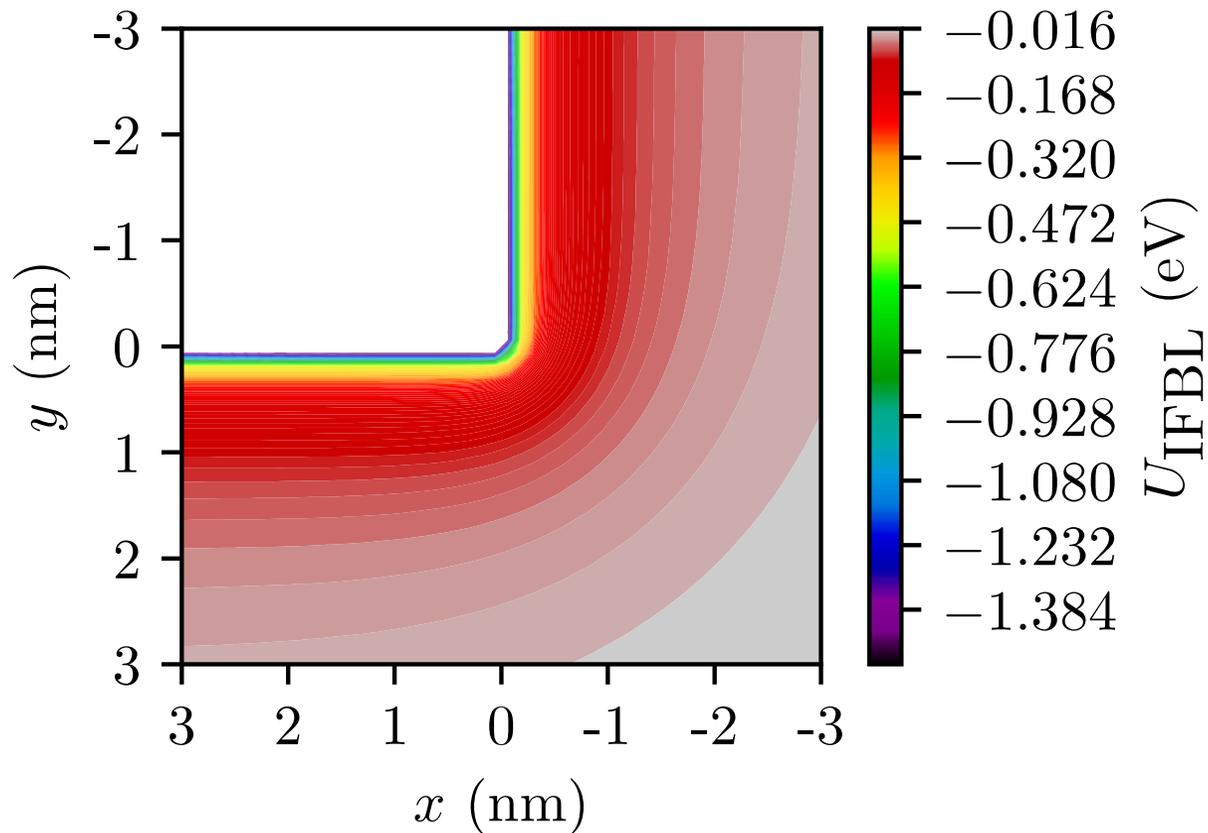
$$\Theta_{;\alpha}(\theta; \theta_0) = \frac{\sinh(\alpha\theta) \cosh(\alpha(\pi - (\Omega - \theta_0)))}{\alpha \sinh(\alpha\pi) \sinh(\alpha\Omega)} + \frac{\sinh(\alpha(\Omega - \theta)) \cosh(\alpha(\pi - \theta_0))}{\alpha \sinh(\alpha\pi) \sinh(\alpha\Omega)}$$

$$\frac{\cosh(\alpha(\pi - |\theta - \theta_0|))}{\alpha \sinh(\alpha\pi)}$$



Results: IFBL Potential Energy

$$U_{\text{IFBL}}(r, \theta) = \frac{-e^2}{8\pi\epsilon r} \int_0^\infty d\alpha \left(\frac{\sinh(\alpha\theta) \cosh(\alpha(\pi - (\Omega - \theta)))}{\sinh(\alpha\Omega) \cosh(\alpha\pi)} + \frac{\sinh(\alpha(\Omega - \theta)) \cosh(\alpha(\pi - \theta))}{\sinh(\alpha\Omega) \cosh(\alpha\pi)} \right)$$

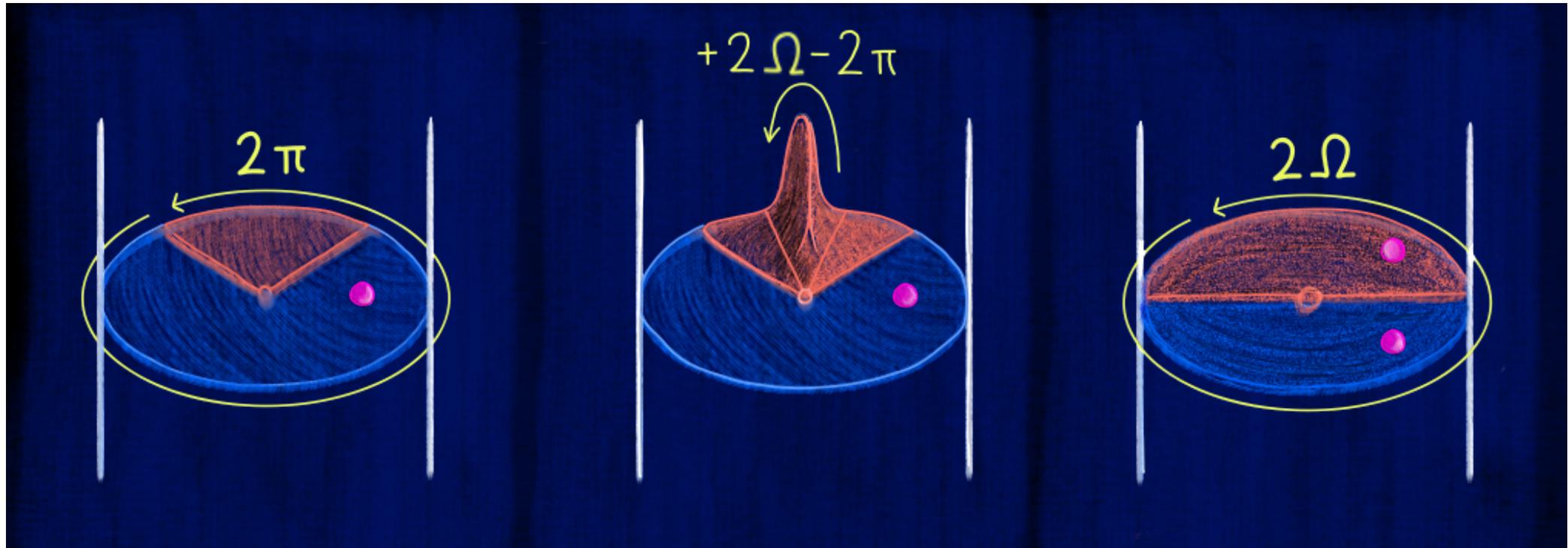


Method of Images on a Cone Manifold

Previously I said that I can not solve asymmetrical problems using method of images...

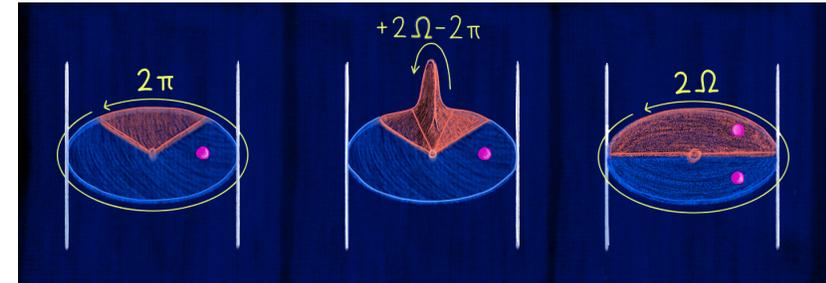
I LIED!

Method of Images on a Cone Manifold



Method of Images on a Cone Manifold (cont.)

- We can now **easily place an image charge**, regardless of asymmetrical geometry
- Results are equal through trig identities

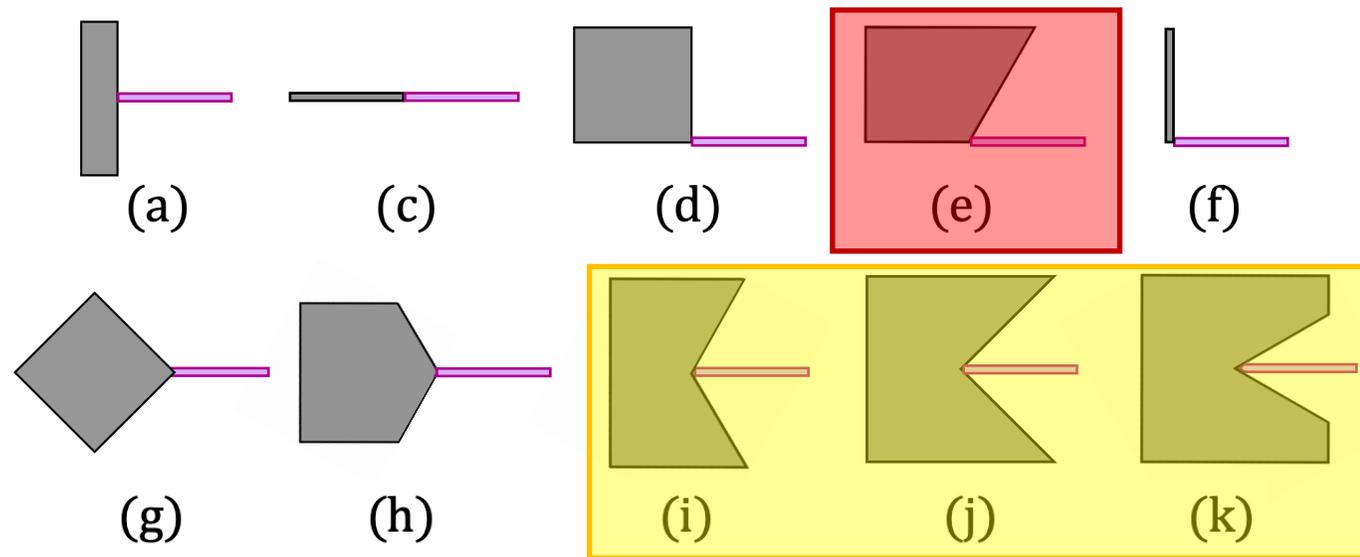


$$U_{\text{IFBL}}(r, \theta) = \frac{-e^2}{8\pi\epsilon r} \int_0^\infty d\alpha \left(\frac{\cosh(\alpha(\Omega - \eta)) \sinh(\alpha\pi)}{\sinh(\alpha\Omega) \cosh(\alpha\pi)} - \frac{\cosh(\alpha(\Omega - 2\theta)) \sinh(\alpha\pi)}{\sinh(\alpha\Omega) \cosh(\alpha\pi)} - \frac{\cosh(\alpha(\pi - \eta))}{\cosh(\alpha\pi)} \right)$$

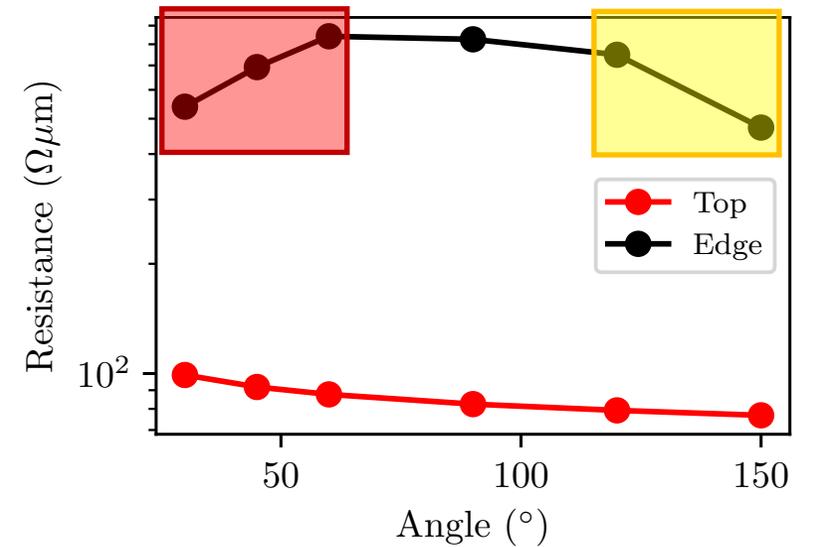
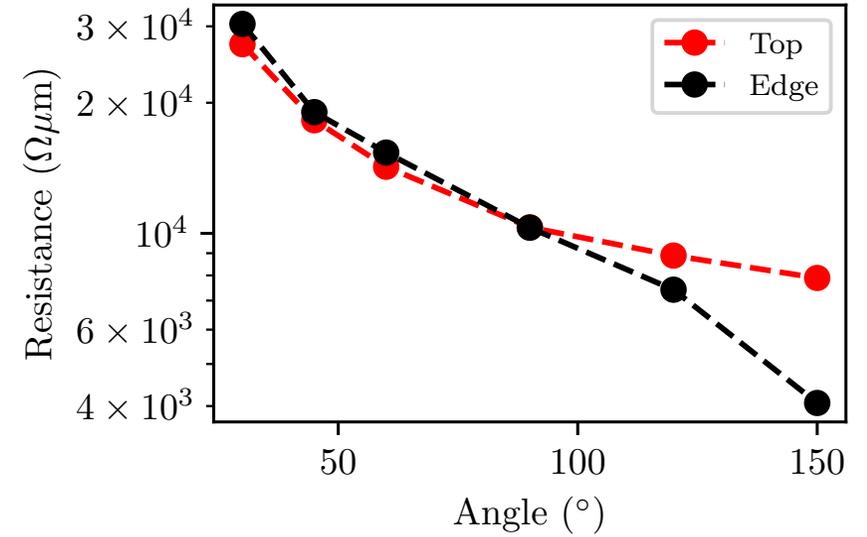
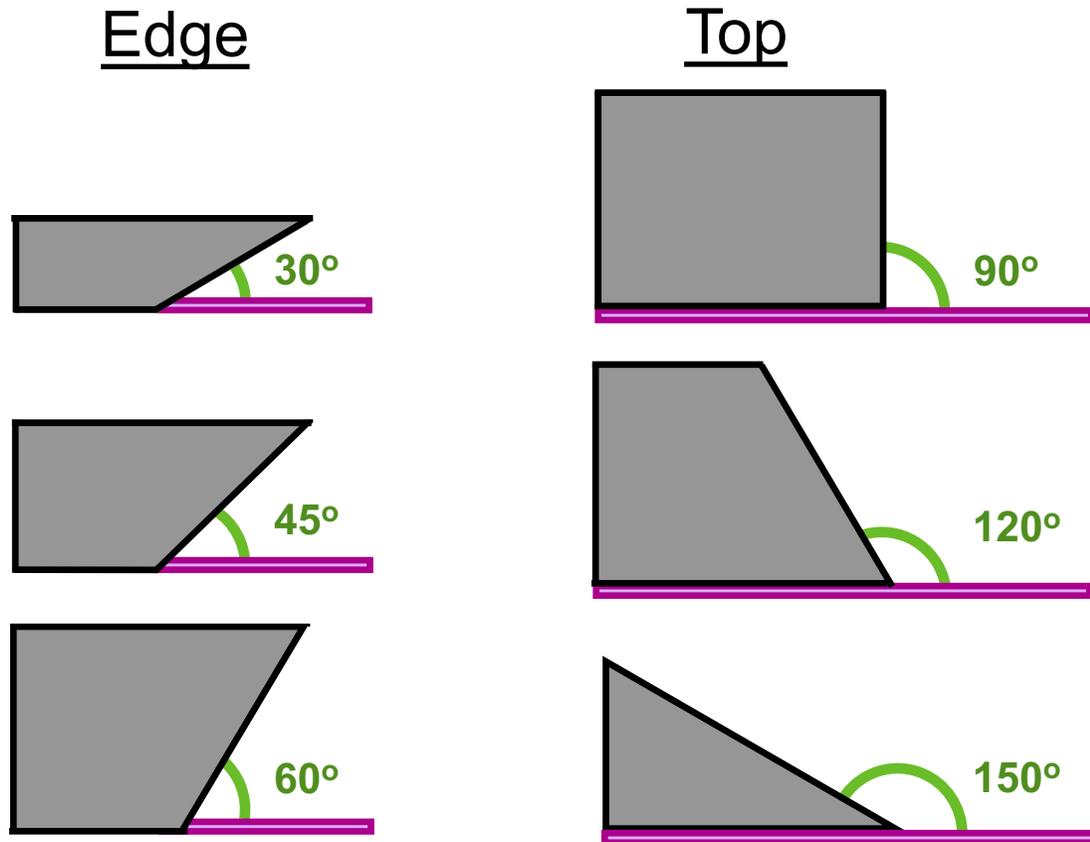
$$U_{\text{IFBL}}(r, \theta) = \frac{-e^2}{8\pi\epsilon r} \left(\int_0^\infty d\alpha \frac{\sinh(\alpha(\Omega - \pi))}{\sinh(\alpha\Omega) \cosh(\alpha\pi)} + \int_0^\infty d\alpha \frac{\cosh(\alpha(\Omega - 2\theta)) \tanh(\alpha\pi)}{\sinh(\alpha\Omega)} \right)$$

Results: U_{IFBL}

	Ω	θ	$U_{\text{IFBL}}/(-e^2/(16\pi\epsilon r))$
a	π	$\pi/2$	1
b	π	θ	$\text{csc}(\theta)$
c	2π	π	0.63
d	$3\pi/2$	π (or $\pi/2$)	0.85
e	$4\pi/3$	π	1.11
f	2π	$3\pi/2$	0.82
g	$3\pi/2$	$\Omega/2 = 3\pi/4$	0.74
h	$4\pi/3$	$\Omega/2 = 2\pi/3$	0.81
i	$2\pi/3$	$\Omega/2 = \pi/3$	1.41
j	$\pi/2$	$\Omega/2 = \pi/4$	1.83
k	$\pi/3$	$\Omega/2 = \pi/6$	2.69



Results: Contact Resistance



Conclusions

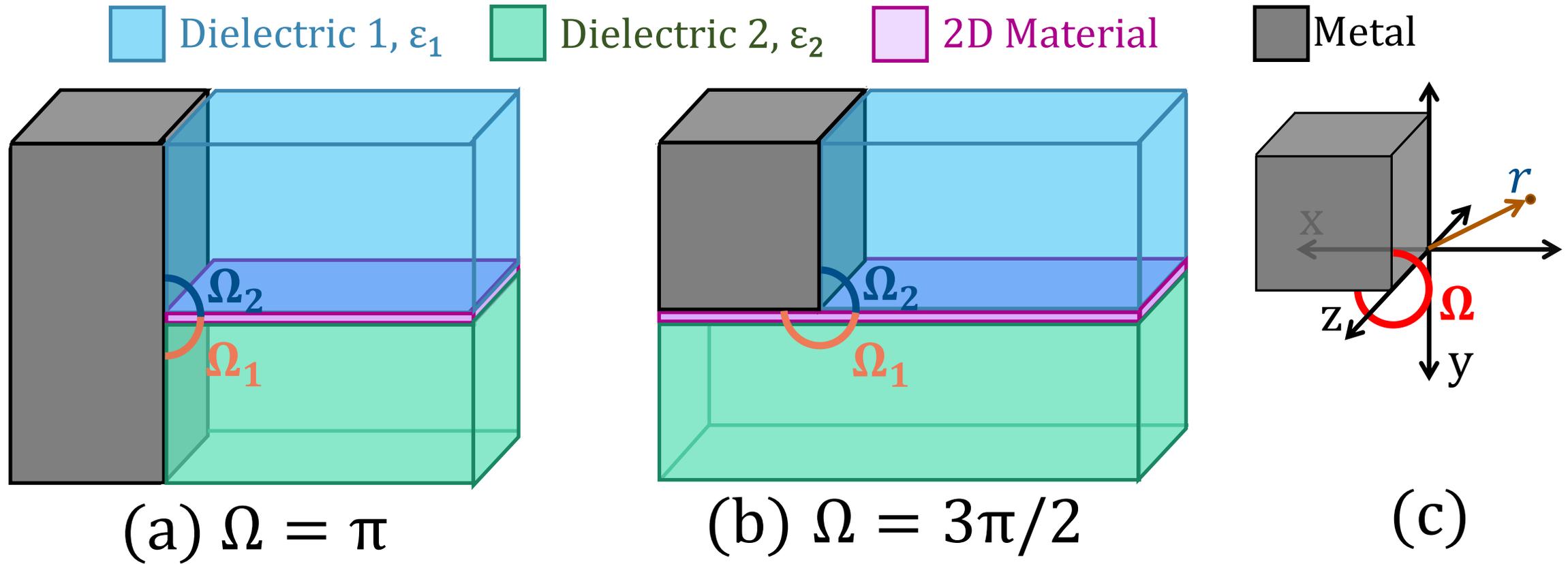
- IFBL necessary for accurate R_C calculation
 - But significantly impacted by contact geometry
- U_{IFBL} determined for asymmetrical geometries:
 - Poisson's equation
 - Cone manifold
- $U_{\text{IFBL}} \propto \frac{1}{\epsilon}$
- Contact angle can act as tunable parameter
 - Impacts Schottky barrier and R_C



Thank you!

Please feel free to ask any questions

Future work



$$U_{\text{IFBL}}(r, \theta) = - \int_0^{-e} \frac{q}{e} V_{\text{I}}(r_0, \theta_0, z_0; r, \theta, z) dq = -\frac{1}{2} e V_{\text{I}} \quad (1)$$

$$\Theta_{\text{I};\alpha}(\theta; \theta_0) = \frac{\sinh(\alpha\theta) \cosh(\alpha(\pi - (\Omega - \theta_0)))}{\alpha \sinh(\alpha\pi) \sinh(\alpha\Omega)} + \frac{\sinh(\alpha(\Omega - \theta)) \cosh(\alpha(\pi - \theta_0))}{\alpha \sinh(\alpha\pi) \sinh(\alpha\Omega)} \quad (2)$$

$$\Theta_{\text{C};\alpha}(\theta; \theta_0) = -\frac{\cosh(\alpha(\pi - |\theta - \theta_0|))}{\alpha \sinh(\alpha\pi)} \quad (3)$$

$$V_{\text{I}} = \frac{e}{4\pi\epsilon\sqrt{rr_0}} \int_0^\infty d\alpha P_{i\alpha-1/2} \left(\frac{(z - z_0)^2 + r_0^2 + r^2}{2rr_0} \right) \quad (4)$$

$$\left(\frac{\sinh(\alpha\theta) \cosh(\alpha(\pi - (\Omega - \theta_0)))}{\sinh(\alpha\Omega) \cosh(\alpha\pi)} + \frac{\sinh(\alpha(\Omega - \theta)) \cosh(\alpha(\pi - \theta_0))}{\sinh(\alpha\Omega) \cosh(\alpha\pi)} \right)$$

$$\Theta_{\text{C};\alpha}^*(\theta; \theta_0) = -\frac{\cosh(\alpha(\Omega - |\theta - \theta_0|))}{\alpha \sinh(\alpha\Omega)} \quad (6)$$

$$U_{\text{IFBL}}(r, \theta) = \frac{-e^2}{8\pi\epsilon r} \int_0^\infty d\alpha \left(\frac{\sinh(\alpha\theta) \cosh(\alpha(\pi - (\Omega - \theta)))}{\sinh(\alpha\Omega) \cosh(\alpha\pi)} + \frac{\sinh(\alpha(\Omega - \theta)) \cosh(\alpha(\pi - \theta))}{\sinh(\alpha\Omega) \cosh(\alpha\pi)} \right) \quad (5)$$

$$U_{\text{IFBL}}(r, \theta) = \frac{-e^2}{8\pi\epsilon r} \int_0^\infty d\alpha \frac{\cosh(\alpha(\Omega - 2\theta)) \sinh(\alpha\pi) + \sinh(\alpha(\Omega - \pi))}{\sinh(\alpha\Omega) \cosh(\alpha\pi)} \quad (7)$$