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Recombination Time in Drift-Diffusion Models of Graphene Field-Effect Transistors

Pedro C. Feijoo, Ferney A. Chaves, D. Jiménez



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Outline

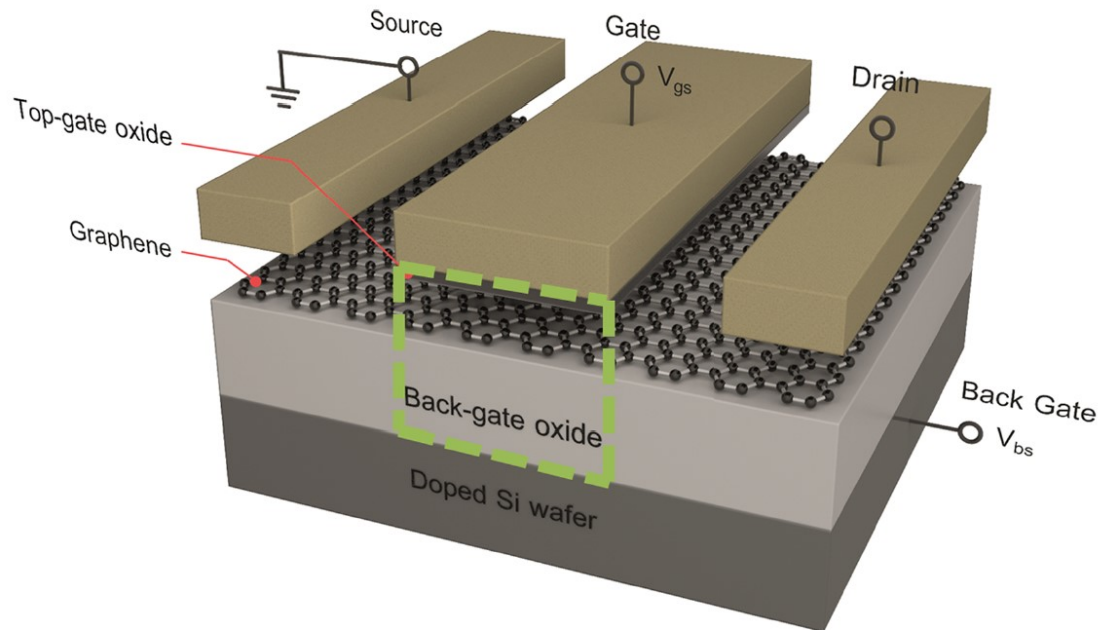
- Introduction
- Mathematical model
- Results
- Conclusions

Recombination Time in Drift-Diffusion Models of
Graphene Field-Effect Transistors

INTRODUCTION

Introduction

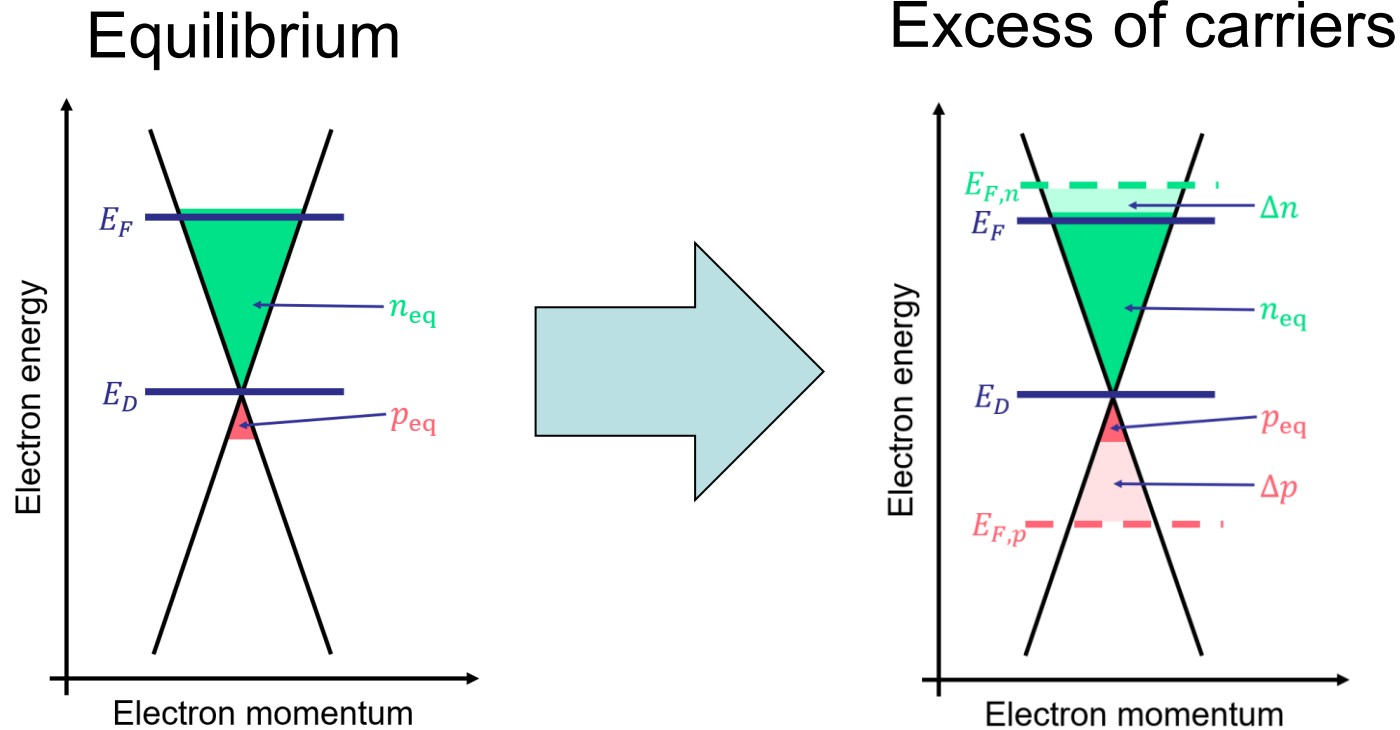
- Graphene Field-Effect Transistor (GFET)
 - High mobility and saturation velocity
 - Applications in radiofrequency and photonics



Introduction

Recombination time in graphene

$\tau \approx 1 - 100$ ps have been reported



$$U = \frac{\Delta n}{\tau}$$

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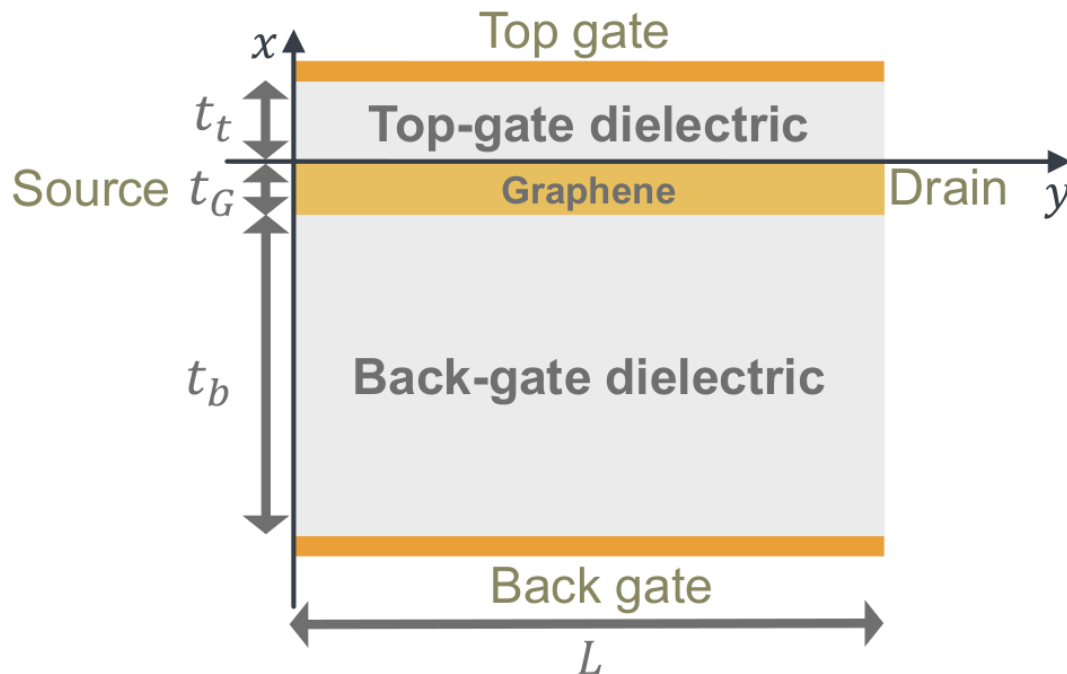
MATHEMATICAL MODEL

Mathematical model

- Poisson's equation

$$\nabla \cdot [\epsilon(x, y) \nabla \phi(x, y)] = \rho_{\text{free}}(x, y)$$

$$\rho_{\text{free}}(x, y) = \frac{q[p(y) - n(y)]}{t_G}$$



Boundary conditions

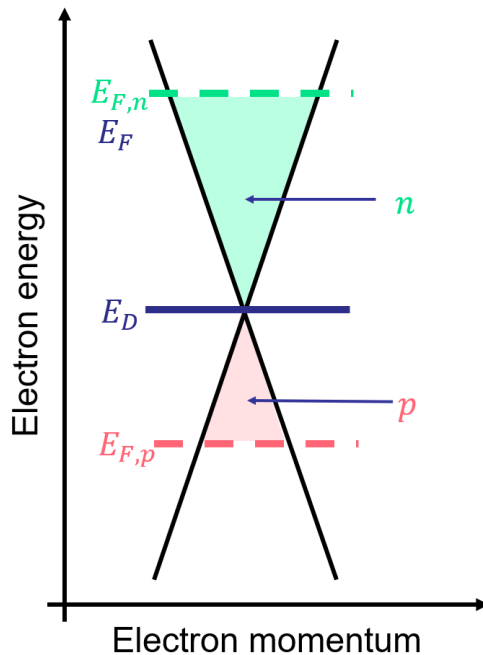
$$\phi(t_t, y) = V_G - V_{gs0}$$

$$\left. \frac{\partial \phi}{\partial y} \right|_{y=0} = \left. \frac{\partial \phi}{\partial y} \right|_{y=L} = 0$$

$$\phi(-t_G - t_b, y) = V_B - V_{bs0}$$

Mathematical model

- Carrier concentration



$$n = \sigma_0 + \frac{2}{\pi} \left(\frac{kT}{\hbar v_F} \right)^2 \mathcal{F}_1 \left[\frac{E_{F,n} - E_D}{kT} \right]$$

$$p = \sigma_0 + \frac{2}{\pi} \left(\frac{kT}{\hbar v_F} \right)^2 \mathcal{F}_1 \left[\frac{E_D - E_{F,p}}{kT} \right]$$

$$E_D(y) = -q\phi \left(x = -\frac{t_G}{2}, y \right).$$

Electron and hole concentrations are deduced from the linear dispersion relation and the Fermi-Dirac statistics.

Mathematical model

- Current continuity equation

$$J_n(y) = n(y)\mu_n(y) \frac{dE_{F,n}(y)}{dy}$$

$$J_p(y) = p(y)\mu_p(y) \frac{dE_{F,p}(y)}{dy}$$

$$\frac{dJ_n(y)}{dy} = qU(y)$$

$$\frac{dJ_p(y)}{dy} = -qU(y)$$

$$J_{DS} = J_n(y) + J_p(y)$$

$$\mu_n(y) = \mu_p(y) = \mu(y) = \frac{\mu_{LF}}{\sqrt[1 + \left(\frac{\mu_{LF}}{v_{sat}} \left| -\frac{\partial \phi}{\partial y} \right|_{x=-\frac{t_G}{2}} \right)^2]^\beta}}$$

Boundary conditions

$$E_{F,n}(L) = E_{F,p}(L) = -qV_D$$

$$E_{F,n}(0) = E_{F,p}(0) = -qV_S$$

Our model solves self-consistently Poisson's equation and current continuity.

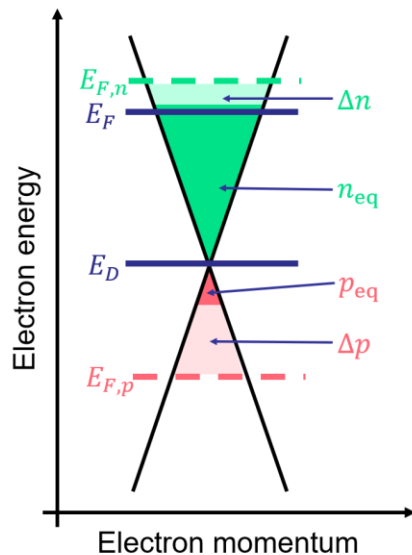
Mathematical model

- Recombination model

$$\frac{dJ_n}{dy} = -\frac{dJ_p}{dy} = qU$$

$$U = \frac{\partial \Delta n}{\partial t} = \frac{\partial \Delta p}{\partial t} = -\frac{\Delta n}{\tau} = -\frac{n - n_{eq}}{\tau}$$

We consider band-to-band recombinations and charge neutrality.



$$U \propto -\Delta n$$

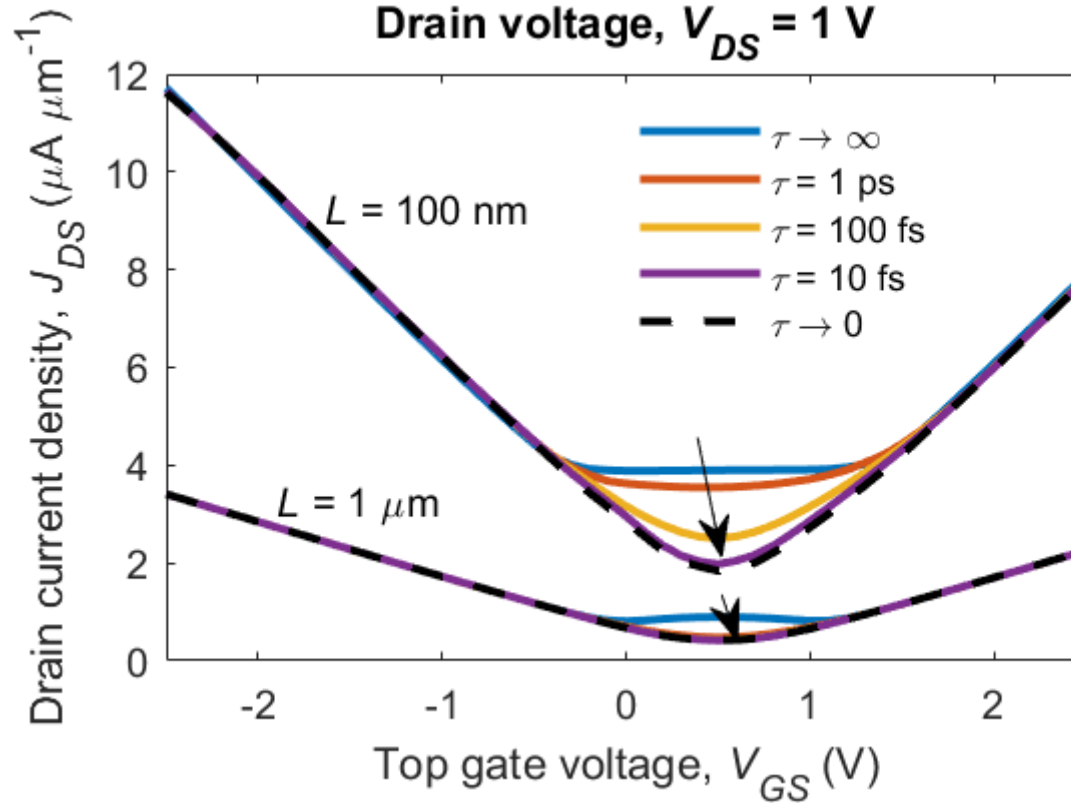
$$\Delta p \approx \Delta n$$

$$U = -\frac{n - n_{eq}(n, p, E_D)}{\tau}$$

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RESULTS

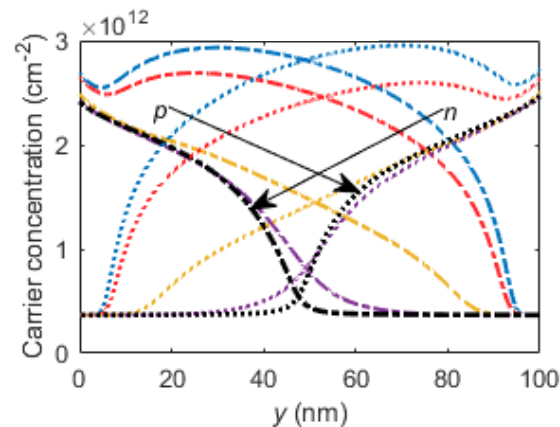
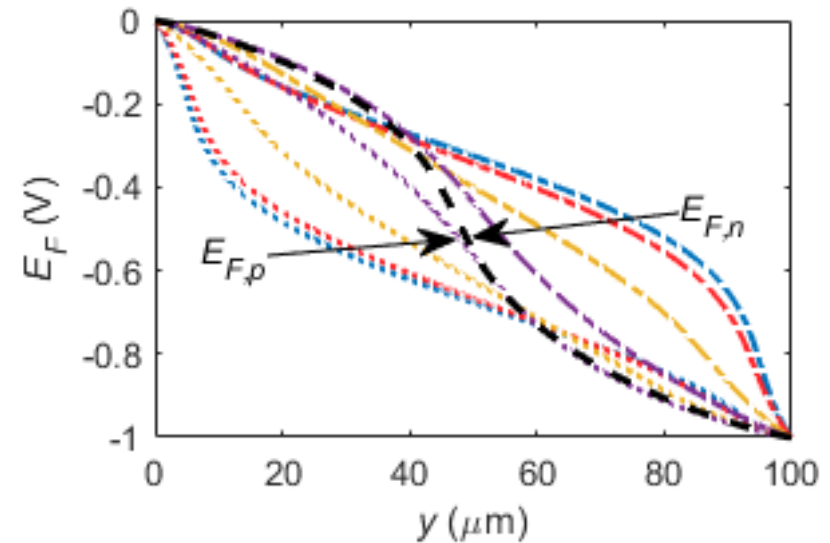
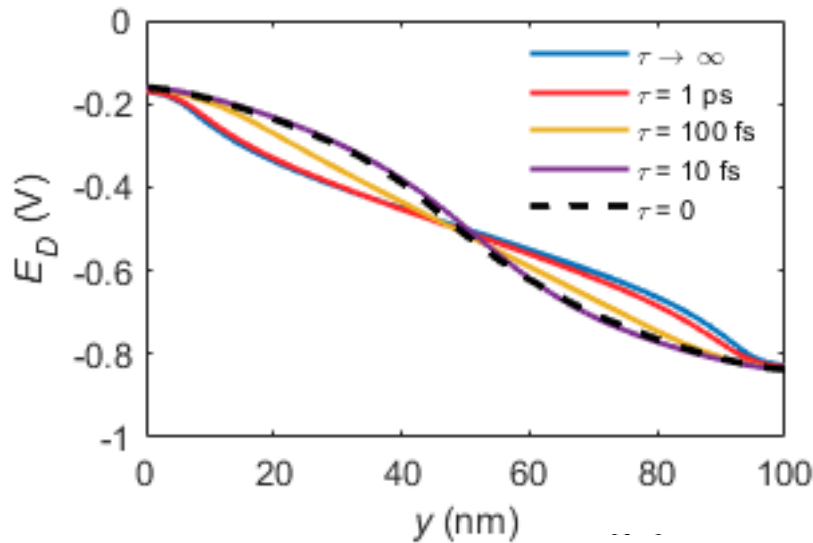
Transfer curves



Recombination mainly affects at biases close to Dirac point.

Dirac point:

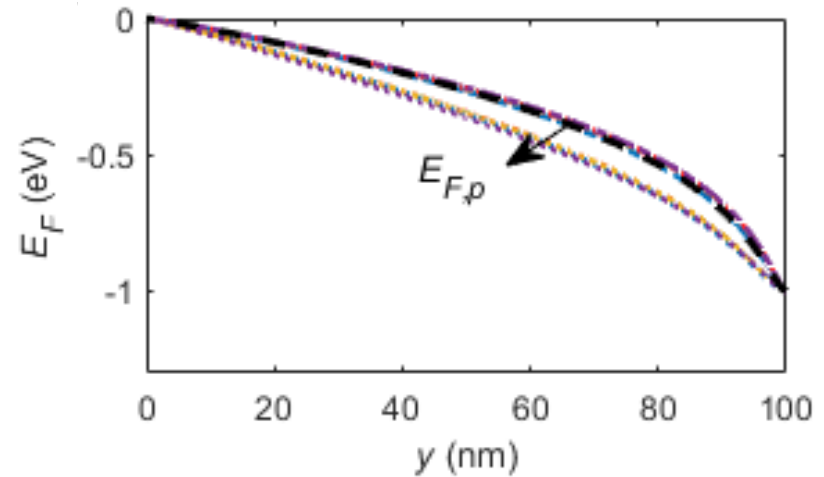
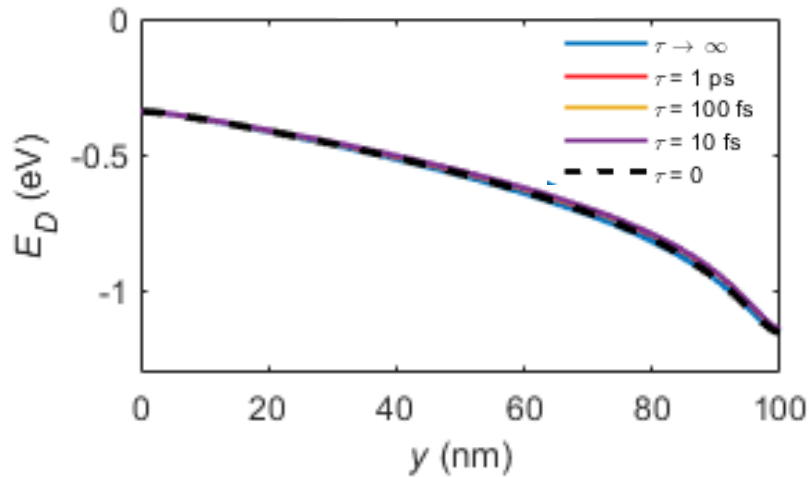
$V_{GS} = 0.5 \text{ V}; V_{DS} = 1 \text{ V}$



$L = 100 \text{ nm}$

Far from Dirac point:

$V_{GS} = 2 \text{ V}; V_{DS} = 1 \text{ V}$



$L = 100 \text{ nm}$

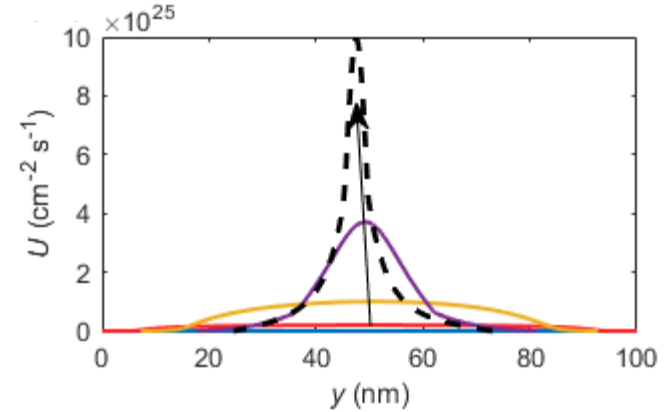
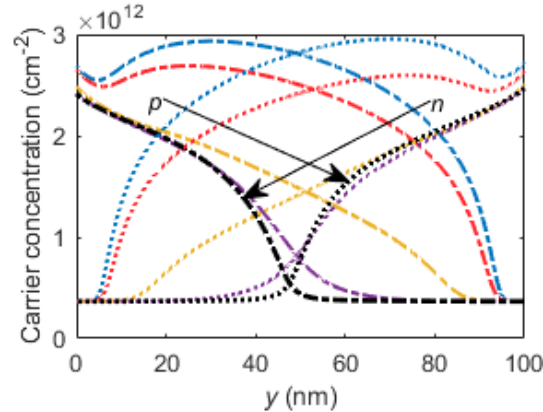
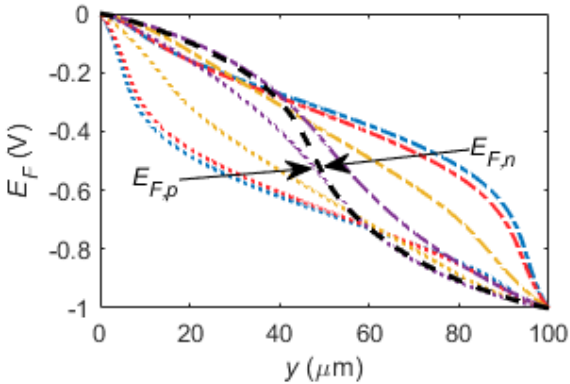
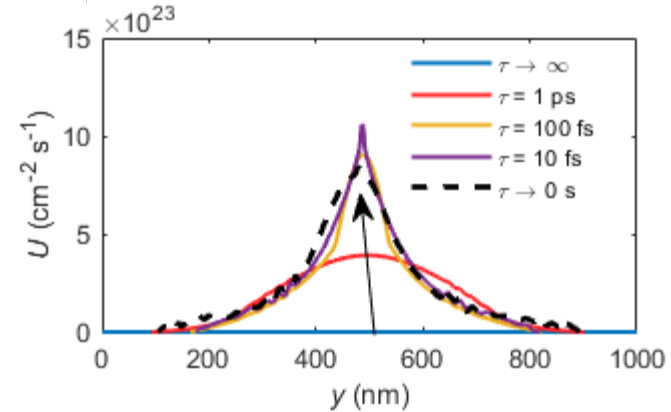
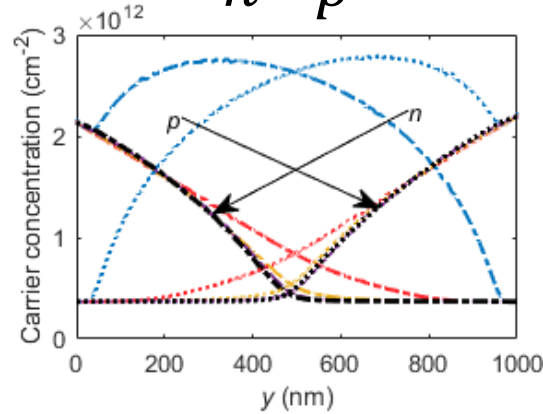
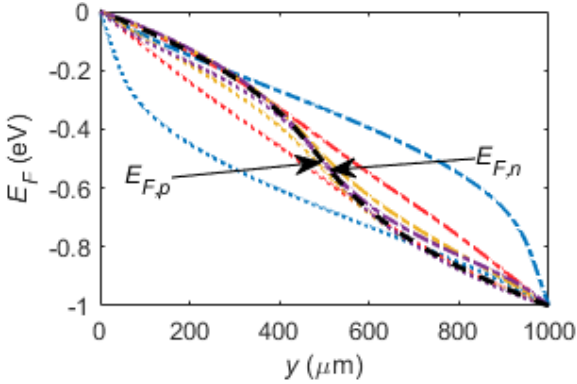
Different channel lengths

$L = 1 \mu\text{m}$

n p

U

$E_{F,n}$ $E_{F,p}$



$L = 100 \text{ nm}$

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CONCLUSIONS

Conclusions

- Effect of recombination in GFETs. We considered band-to-band recombination with charge neutrality
- Recombination affects mainly around the Dirac bias, where charge transport is ambipolar
- For a 100 nm channel,
 - $\tau < 10$ fs \rightarrow single quasi-Fermi level approximation ($\tau \rightarrow 0$ s)
 - $\tau > 1$ ps \rightarrow null recombination rate ($U = 0$)
- These limits are larger for longer channels

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Thanks for your attention!