



Institut Català
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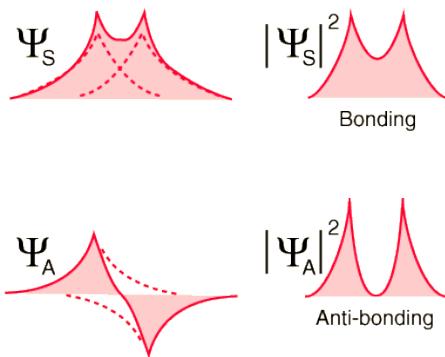
TranSIESTA: Advanced Applications in Electrochemistry and Spintronics

Pablo Ordejón

Catalan Institute of Nanoscience and Nanotechnology – ICN2
Barcelona - Spain

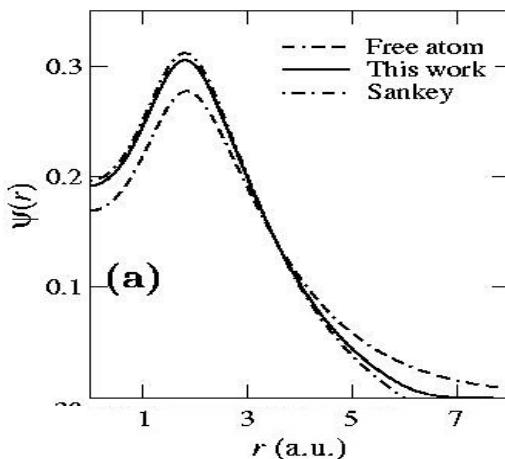
DFT with atomic orbitals

siesta



$$\text{LCAO: } \psi_n(r) = \sum_{\mu} c_{n\mu} \phi_{\mu}(r)$$

(pseudo)atomic orbitals
- Short ranged
- Arbitrarily complete



$$\phi_{\mu}(\vec{r}) = \varphi_{\mu}(|\vec{r}|) Y_{lm}(\theta, \varphi)$$

Spherical harmonics



s



p



d



f

Scaling vs system size

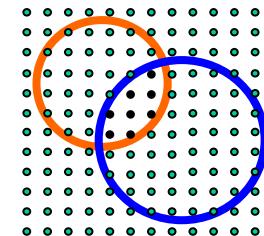


1. Hamiltonian Building:

Always O(N)

Load Balancing for **real-space grid operations**
(inhomogeneous systems)

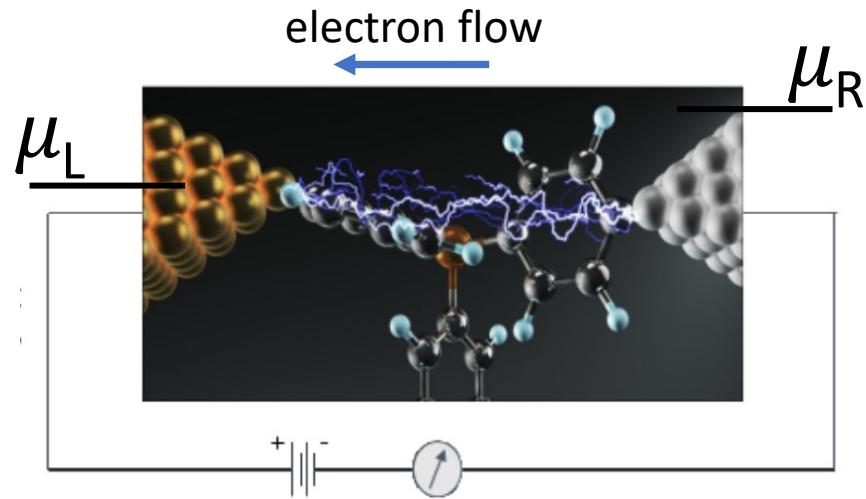
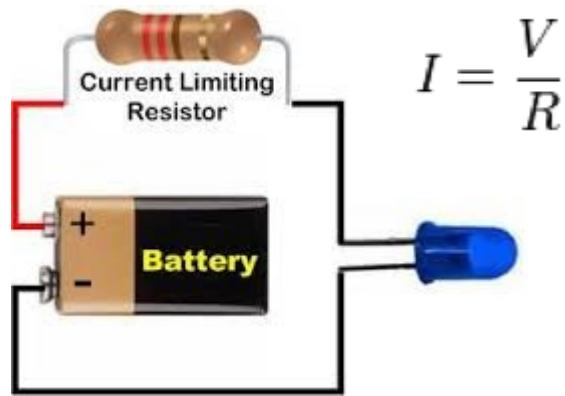
$$\hat{h}_{\mu\nu} = \langle \varphi_\mu | \hat{h} | \varphi_\nu \rangle$$



2. Hamiltonian Solvers: obtaining the charge density from the Hamiltonian

- Standard parallel diagonalization libraries – **SCALAPACK** $O(N^3)$
- New diagonalizers: **ELPA, MRRR** $O(N^3)$
- **PEXSI method** – massive parallelization and reduced scaling: $O(N, N^{3/2}, N^2)$ for 1, 2 and 3D
- $O(N)$ method based on **localized occupied orbitals** (Ordejón et al; Mauri et al)
- $O(N)$ method (**Fermi Operator Expansion**), porting the CheSS library from the BigDFT project to SIESTA.

Challenge: Electronic Transport at the Nanoscale

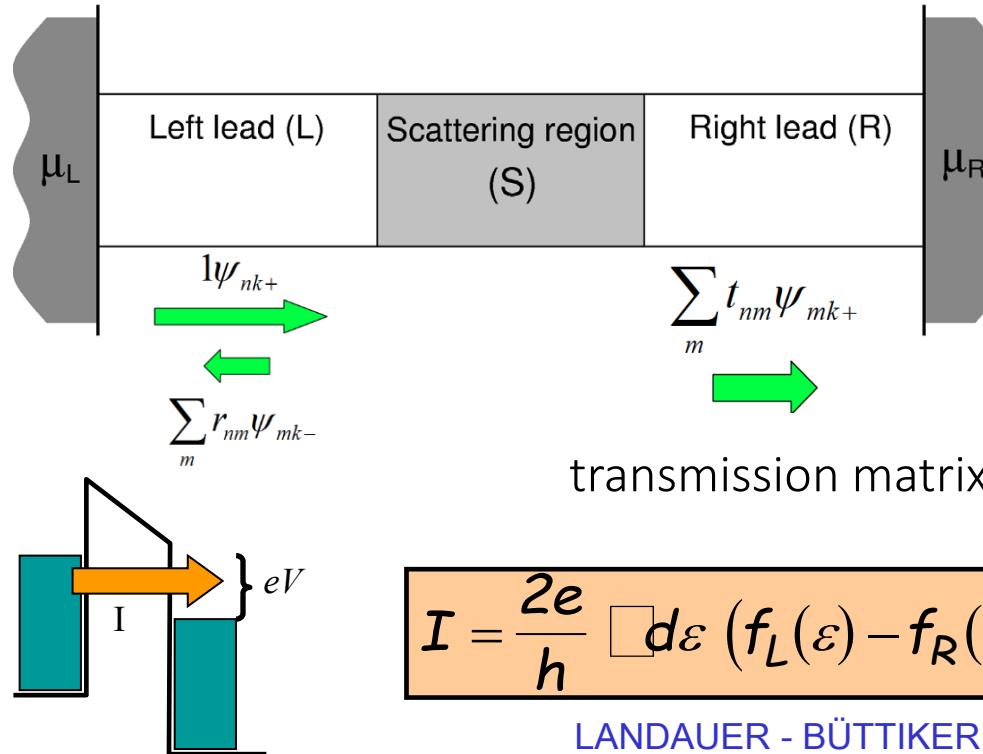


Molecular Electronics

Nanoelectronics
(devices with nanometric features)

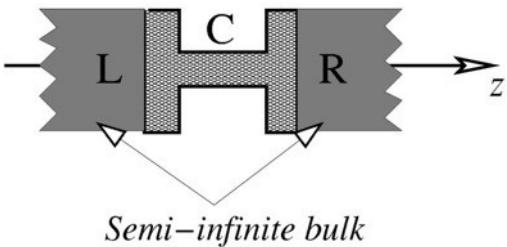
- Open systems (intensive quantities)
- Non-Periodic
- Semi-infinite electrodes
- Out of equilibrium:
 - Chemical potential is not unique
 - Net electronic current flowing

Landauer formulation of electronic transport



Transmission from Green's functions: $t(\varepsilon) = [\Gamma_R(\varepsilon)]^{1/2} \mathbf{G}(\varepsilon) [\Gamma_L(\varepsilon)]^{1/2}$,

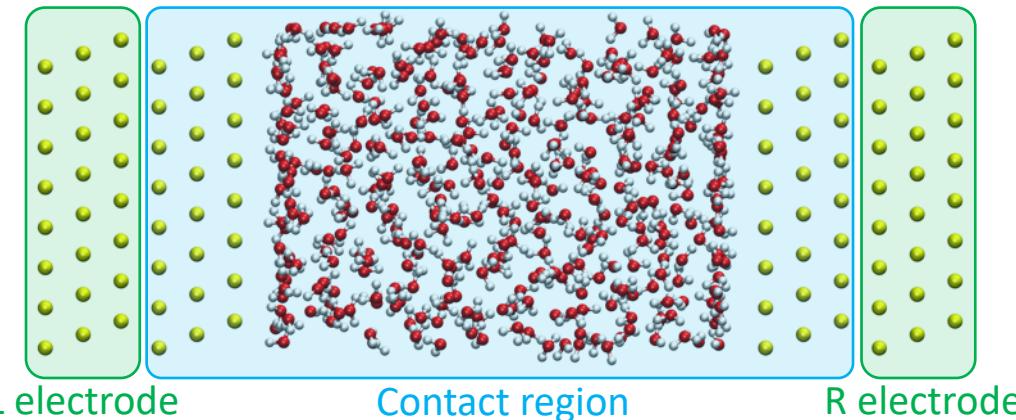
Green's function formalism - TranSIESTA



Green's function

$$G(z) = [z - H]^{-1} \quad (z = \varepsilon + i\eta) \quad \rho(\varepsilon) = -\frac{1}{\pi} \text{Im} [G(z)] \quad (\eta \rightarrow 0^+)$$

$$G(z) = [z - H - \Sigma_L(z) - \Sigma_R(z)]^{-1}$$



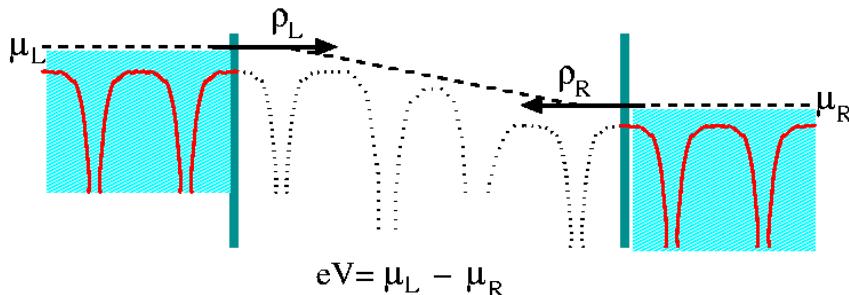
- All screening must happen within the contact region C (L and R are bulk-like)
- Green's Functions are computed in the simulation cell "L+C+R", connected to semi-infinite bulk electrodes (self-energies Σ).

How to handle non-equilibrium distributions



How to compute the density if the chemical potential is not unique?

Non-equilibrium Green's Functions (Keldish formalism):



$$D_{\mu\nu} = \int_{-\infty}^{\infty} d\varepsilon \left[\rho_{\mu\nu}^L(\varepsilon) n_F(\varepsilon - \mu_L) + \rho_{\mu\nu}^R(\varepsilon) n_F(\varepsilon - \mu_R) \right]$$

$$\rho_{\mu\nu}^L(\varepsilon) = \frac{1}{\pi} [\mathbf{G}(\varepsilon) \boldsymbol{\Gamma}_L(\varepsilon) \mathbf{G}^\dagger(\varepsilon)]_{\mu\nu}$$

$$\boldsymbol{\Gamma}_L(z) \equiv i[\boldsymbol{\Sigma}_L(\varepsilon) - \boldsymbol{\Sigma}_L(\varepsilon)^\dagger]/2$$

First principles (DFT): **TranSIESTA**

- Semi-infinite electrodes through Self-Energies
- Non-Equilibrium Green's Functions for finite voltage
- Conductance through Landauer formulation

Brandbyge, Mozos, Ordejon, Taylor, Stokbro PRB 65, 165401 (2002)

Papior, Lorente, Frederiksen, García, Brandbyge, Comp. Phys. Comm. 212, 8 (2017)

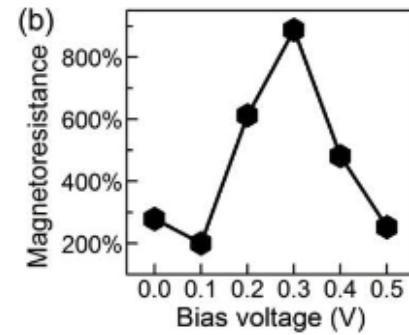
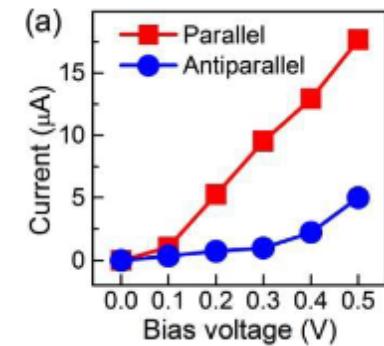
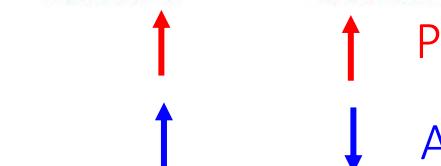
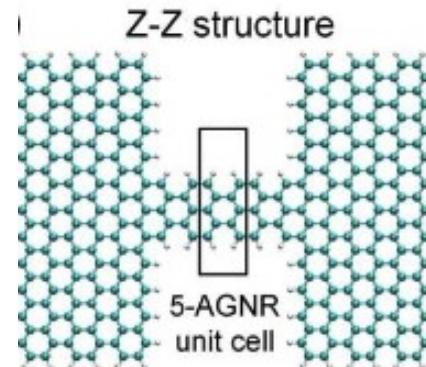
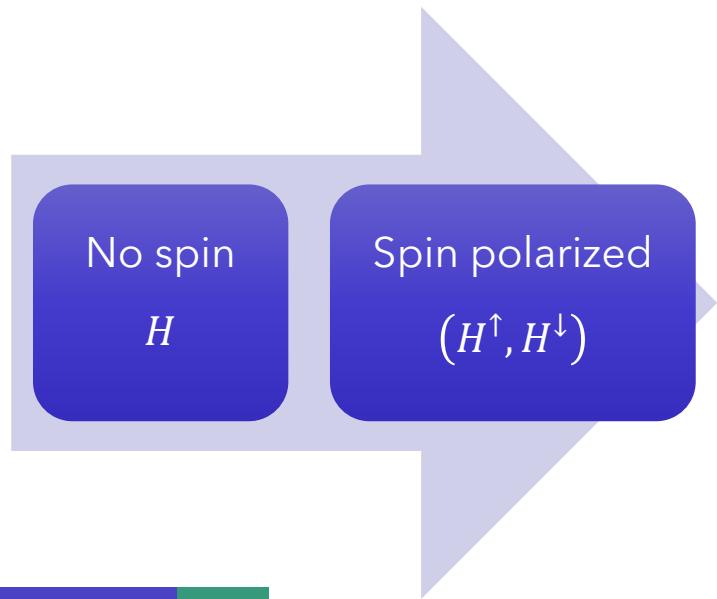
- Accelerated Green's function solvers (linear scaling with contact length)
- Improved parallelization
- Multiterminal devices; Gate voltages;

Spin Transport

Electrochemistry

Spintronics with TranSIESTA

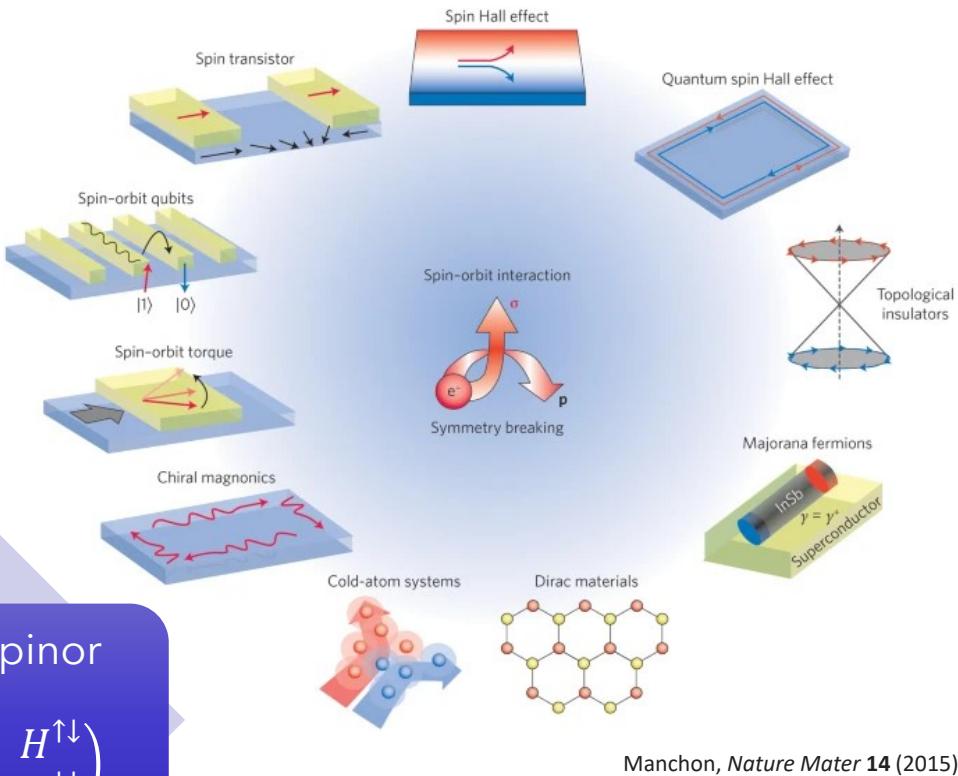
- Spin-polarized DFT
 - Collinear spin directions: common magnetization direction through space
 - No Spin-orbit coupling



S. Li, Y.-W. Son, Su Ying Quek,
Appl. Phys. Lett. 105, 242413 (2014)

- **Non-Collinear magnetism**
- **Spin-Orbit Coupling**

→ require spinor description



Non-eq. Green's Functions with spinors

Non-eq. Green's Function (NEGFF)

$$G_{\mathbf{k}}(z) = \left(z S_{\mathbf{k}} - H_{\mathbf{k}}^{\text{DFT}} - \sum_{\epsilon} \Sigma_{\epsilon, \mathbf{k}}(z) \right)^{-1} \longrightarrow \begin{pmatrix} G_{\mathbf{k}}^{\uparrow\uparrow} & G_{\mathbf{k}}^{\uparrow\downarrow} \\ G_{\mathbf{k}}^{\downarrow\uparrow} & G_{\mathbf{k}}^{\downarrow\downarrow} \end{pmatrix} = \left(z \begin{pmatrix} S_{\mathbf{k}} & \mathbf{0} \\ \mathbf{0} & S_{\mathbf{k}} \end{pmatrix} - \begin{pmatrix} H_{\mathbf{k}}^{\uparrow\uparrow} & H_{\mathbf{k}}^{\uparrow\downarrow} \\ H_{\mathbf{k}}^{\downarrow\uparrow} & H_{\mathbf{k}}^{\downarrow\downarrow} \end{pmatrix} - \sum_{\epsilon} \begin{pmatrix} \Sigma_{\epsilon, \mathbf{k}}^{\uparrow\uparrow} & \Sigma_{\epsilon, \mathbf{k}}^{\uparrow\downarrow} \\ \Sigma_{\epsilon, \mathbf{k}}^{\downarrow\uparrow} & \Sigma_{\epsilon, \mathbf{k}}^{\downarrow\downarrow} \end{pmatrix} \right)^{-1}$$

Spectral function

$$\mathcal{A}_{\epsilon, \mathbf{k}}(z) = G_{\mathbf{k}}(z) \Gamma_{\epsilon, \mathbf{k}}(z) G_{\mathbf{k}}^\dagger(z) \longrightarrow \begin{pmatrix} \mathcal{A}_{\epsilon, \mathbf{k}}^{\uparrow\uparrow} & \mathcal{A}_{\epsilon, \mathbf{k}}^{\uparrow\downarrow} \\ \mathcal{A}_{\epsilon, \mathbf{k}}^{\downarrow\uparrow} & \mathcal{A}_{\epsilon, \mathbf{k}}^{\downarrow\downarrow} \end{pmatrix} = \begin{pmatrix} G_{\mathbf{k}}^{\uparrow\uparrow} & G_{\mathbf{k}}^{\uparrow\downarrow} \\ G_{\mathbf{k}}^{\downarrow\uparrow} & G_{\mathbf{k}}^{\downarrow\downarrow} \end{pmatrix} \begin{pmatrix} \Gamma_{\epsilon, \mathbf{k}}^{\uparrow\uparrow} & \Gamma_{\epsilon, \mathbf{k}}^{\uparrow\downarrow} \\ \Gamma_{\epsilon, \mathbf{k}}^{\downarrow\uparrow} & \Gamma_{\epsilon, \mathbf{k}}^{\downarrow\downarrow} \end{pmatrix} \begin{pmatrix} G_{\mathbf{k}}^{\uparrow\uparrow} & G_{\mathbf{k}}^{\uparrow\downarrow} \\ G_{\mathbf{k}}^{\downarrow\uparrow} & G_{\mathbf{k}}^{\downarrow\downarrow} \end{pmatrix}^\dagger$$

Charge density

$$\rho = \frac{1}{2\pi} \iint_{BZ} d\epsilon dk \sum_i \mathcal{A}_{i, \mathbf{k}}(\epsilon + 0^+) n_F(\epsilon, \mu_i, T_i) e^{-i\mathbf{k}\cdot\mathbf{r}} \longrightarrow \begin{pmatrix} \rho^{\uparrow\uparrow} & \rho^{\uparrow\downarrow} \\ \rho^{\downarrow\uparrow} & \rho^{\downarrow\downarrow} \end{pmatrix} \quad \text{or} \quad \rho \text{ and } \vec{m}$$

Transmission Function for Spinors

Spin channel projected transmission

$$T_k^{\epsilon, \epsilon'}(\epsilon, \vec{\sigma}_n, \vec{\sigma}_m) = \text{Tr} \left\{ \left\langle \vec{\sigma}_m \right| \left(s_k^{\epsilon, \epsilon'} \right)^{\dagger} \left| \vec{\sigma}_n \right\rangle \left\langle \vec{\sigma}_n \right| s_k^{\epsilon, \epsilon'} \left| \vec{\sigma}_m \right\rangle \right\}$$

Conductivity
Current

Transmission can be decomposed into different contributions, e.g.:

$T_k^{\epsilon, \epsilon'}(\epsilon, \uparrow, \uparrow)$ spin-up transmission,

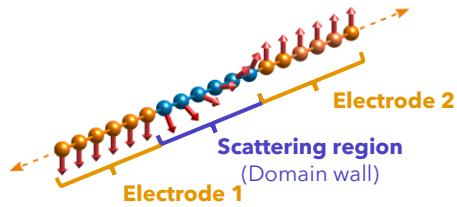
$T_k^{\epsilon, \epsilon'}(\epsilon, \downarrow, \downarrow)$ spin-down transmission,

$T_k^{\epsilon, \epsilon'}(\epsilon, \uparrow, \downarrow), T_k^{\epsilon, \epsilon'}(\epsilon, \downarrow, \uparrow)$ spin-flip

Spin directions can be arbitrary and different for the two electrodes

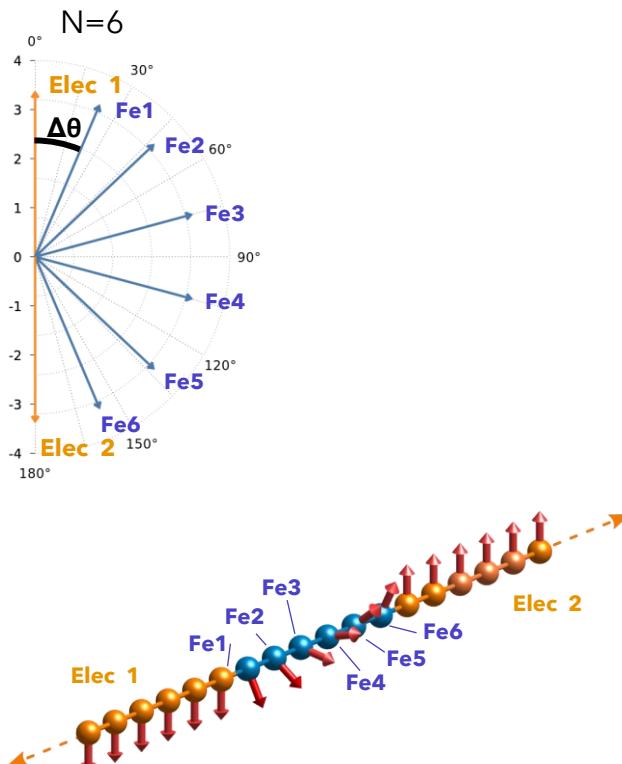
Test case: Domain Walls in Monatomic Fe Chain

- Non-collinear spins
- Prediction of spin orientations
- Anisotropic magnetoresistance
- Spin channel projection



Bloch Domain Wall in Monatomic Chain

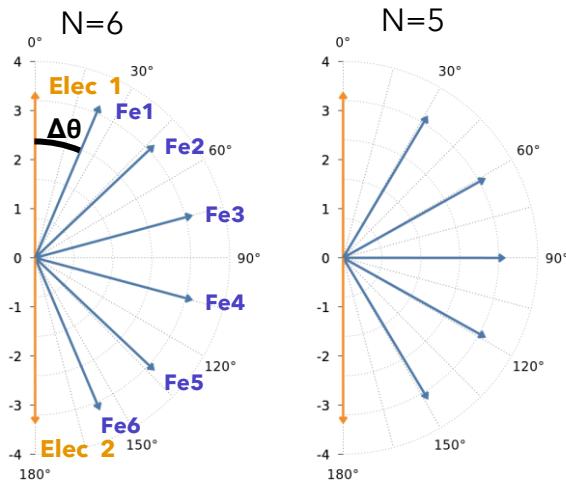
Domain wall length ($\alpha=180^\circ$)



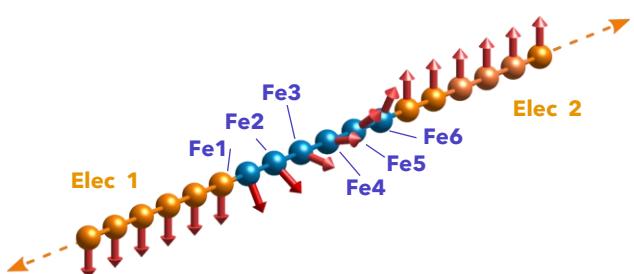
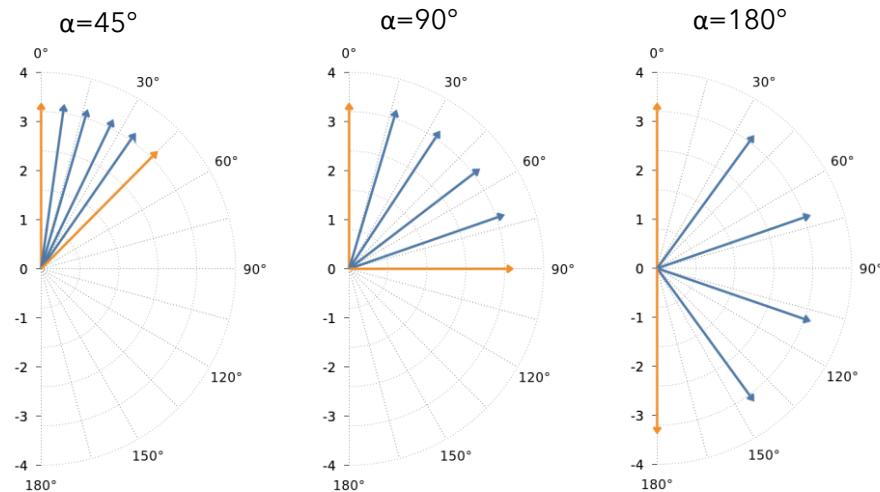
- All spin moments perpendicular to the chain axis
- Uniform change in spin direction along the domain wall

Bloch Domain Wall in Monatomic Chain

Domain wall length ($\alpha=180^\circ$)



Domain wall angle ($N=4$)



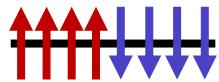
- All spin moments perpendicular to the chain axis
- Uniform change in spin direction along the domain wall

Simple picture: Pristine monatomic chain

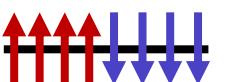
- Uniform magnetic coupling
- **Uniform domain wall ✓**

Domain Wall Conductance - Iron Chain

(a) Abrupt (w/o SOC)



(b) Abrupt (w/ SOC)



(c) Néel (w/ SOC)



(d) Bloch (w/ SOC)



Total (||)

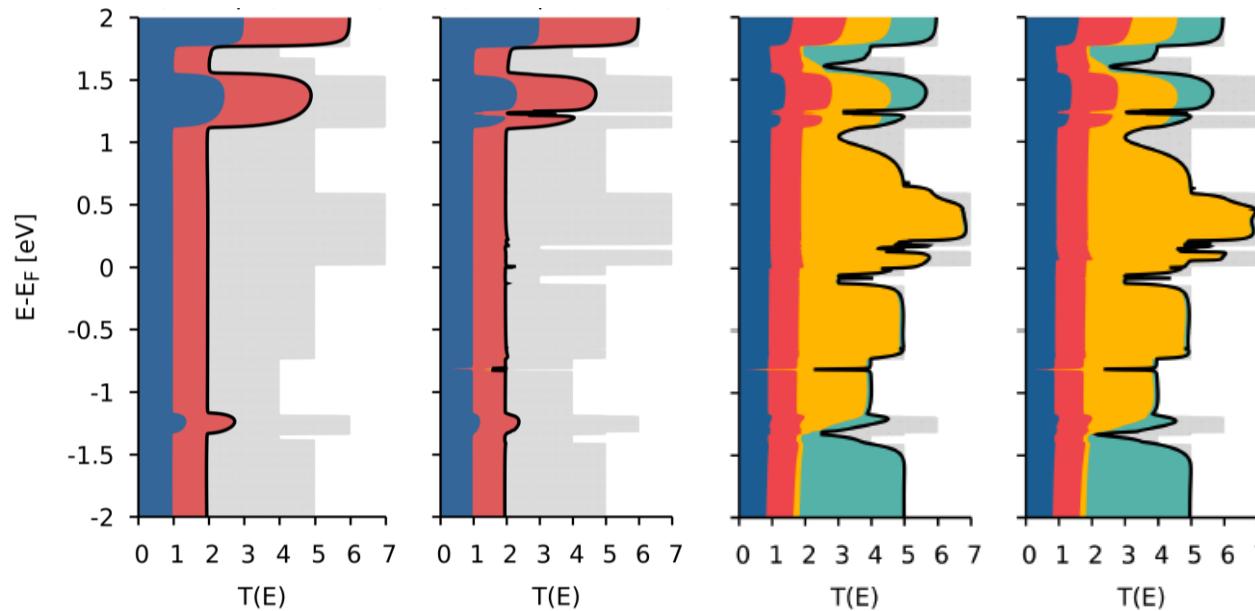
Total

↓↓

↑↑

↑↓

↓↑



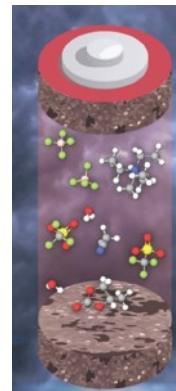
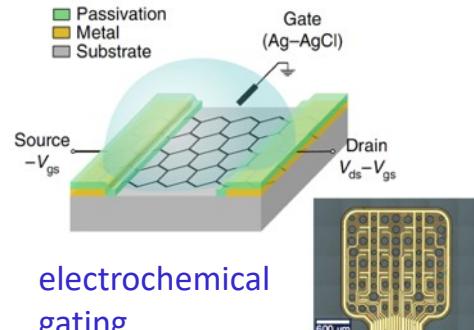
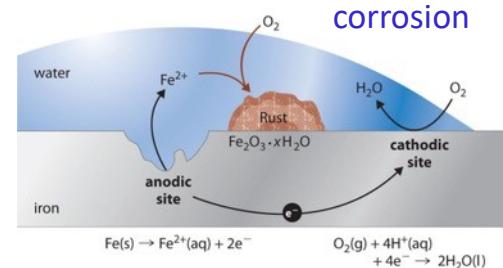
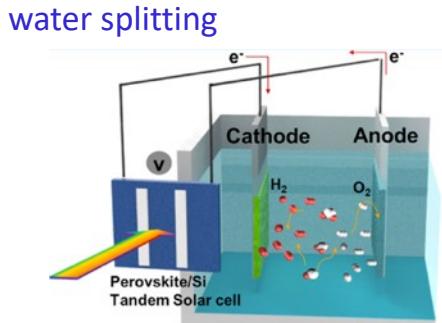
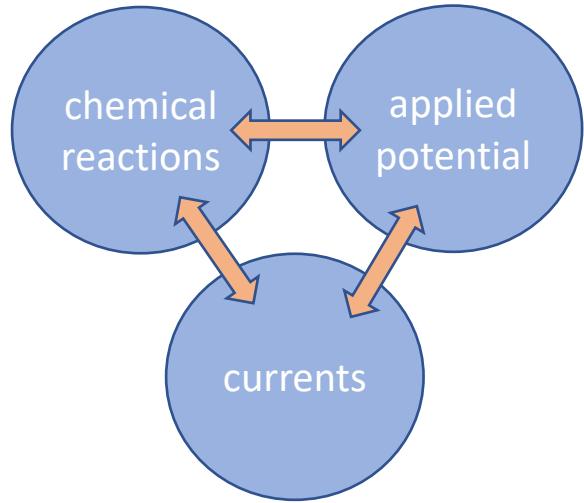
Abrupt domain wall

- No spin flips
- Small SOC effect

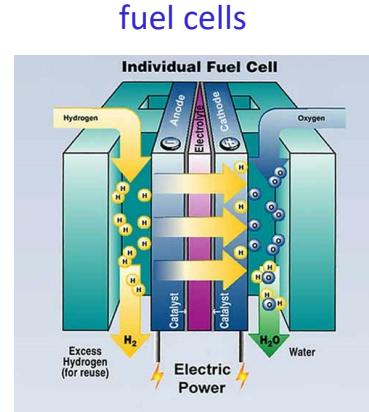
Smooth domain walls

- Significant spin flip
- Very similar for Néel and Bloch

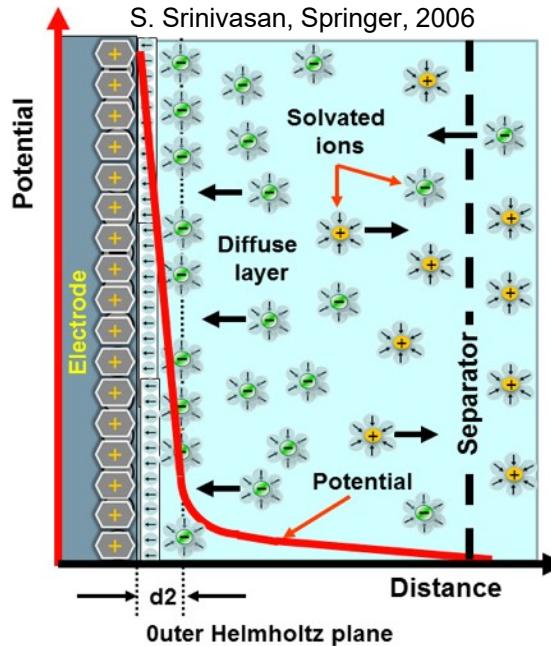
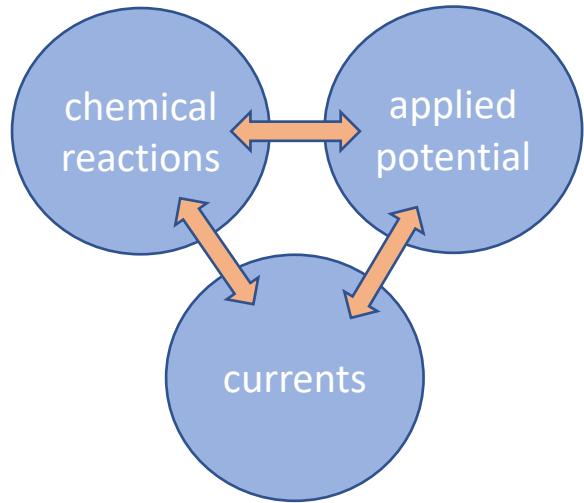
Understanding the electrified interface



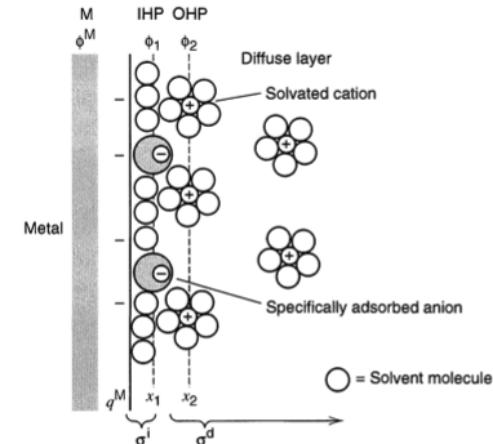
batteries
supercaps



Understanding the electrified interface



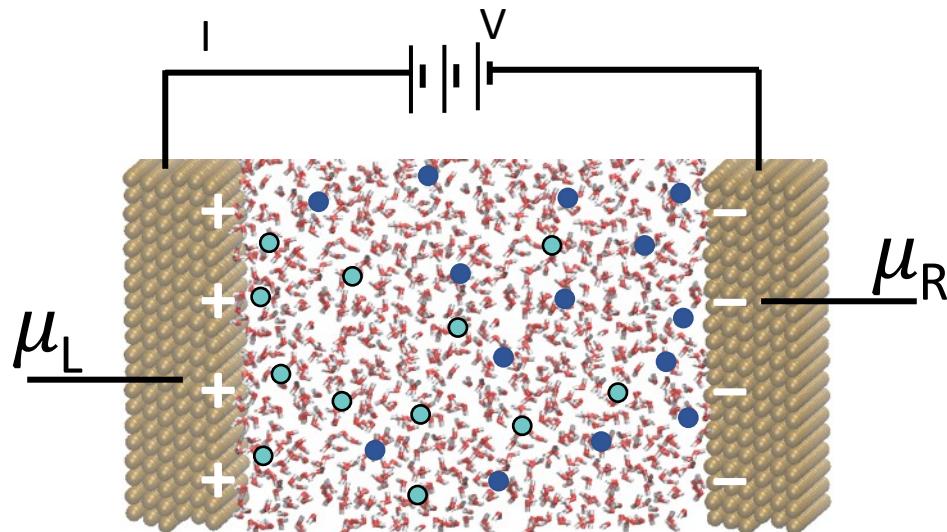
Bard and Faulkner, Wiley, 2001



- Structure of the Double Layer
- Effect of the applied potential
- Currents (faradaic and non-faradaic)
- Chemical reactions

Challenges for first-principles modeling (DFT)

- Open systems (intensive quantities)
- Non-Periodic
- Semi-infinite electrodes
- Out of equilibrium:
 - Chemical potential is not unique
 - Flowing currents (ionic, electronic)
 - Electron transfer reactions (Faradaic)



Proof-of-Concept Simulations

Au(111) electrodes + Water

- pure H₂O
- H₂O with Na, Cl, and Na+Cl (at ~1.25 M)

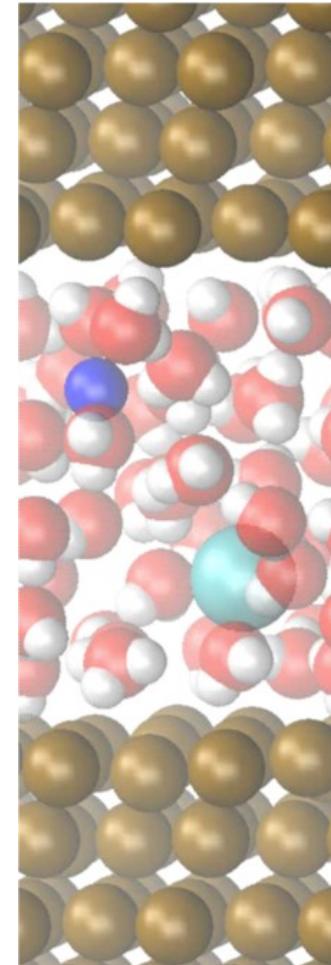
Au(111) surface: (3x4) supercell in-plane

2 layers for each electrode in simulation box
(attached to semi-infinite electrodes)

40 Water Molecules

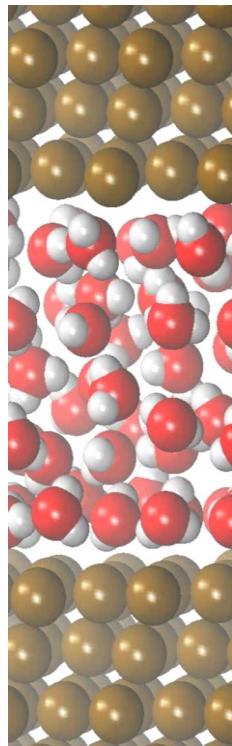
~ 17Å between electrodes (~ 5 water layers)

- DFT: VdW functional of Dion et al., PRL 92, 246401 (2004)
- Molecular Dynamics: 4ps; Berendsen thermostat (300K)
- Γ point in the surface plane
- 200 Ry real space mesh
- DZP basis

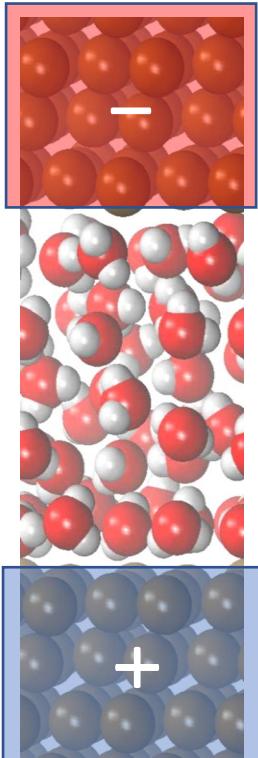


Water at zero and finite bias

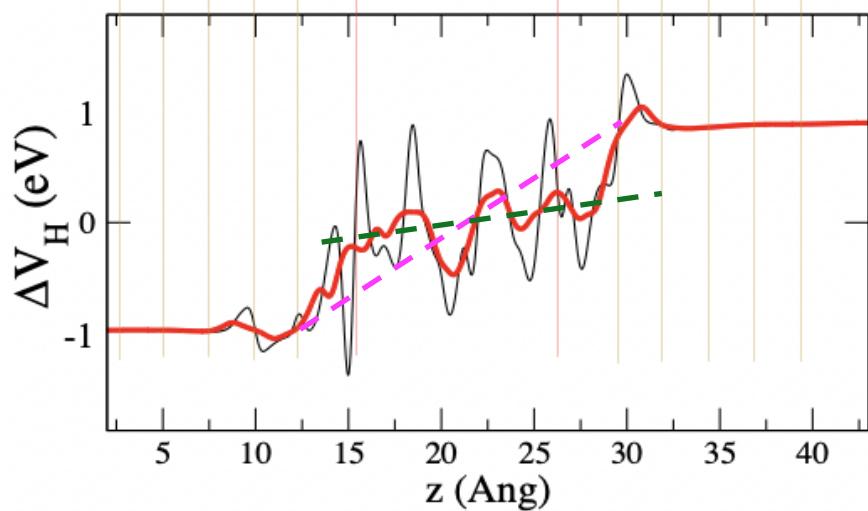
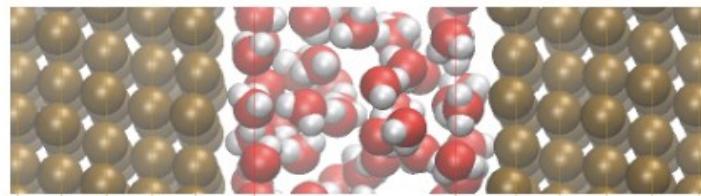
$V = 0 \text{ V}$



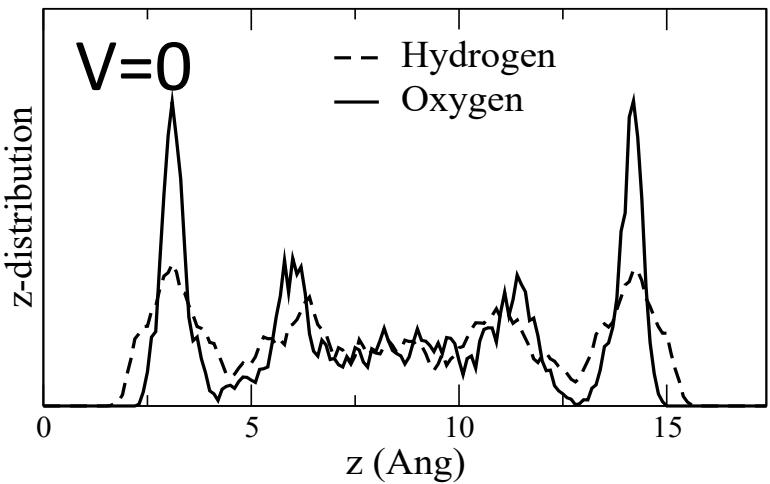
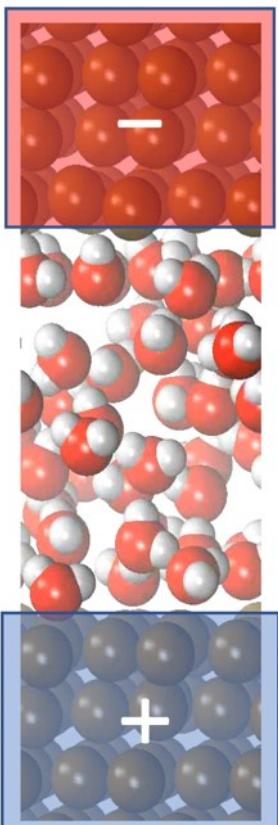
$V = -2 \text{ V}$



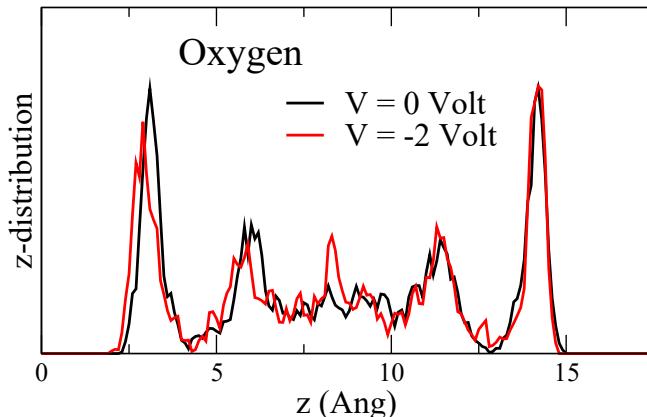
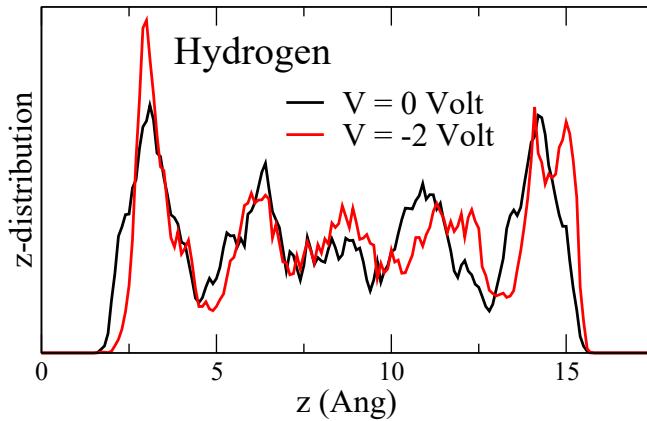
$V = -2 \text{ V}$



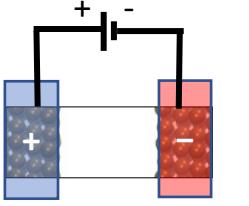
Water Structure

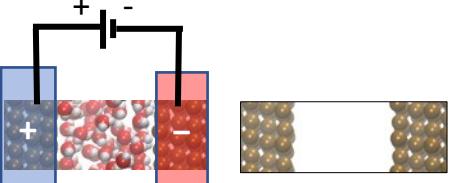


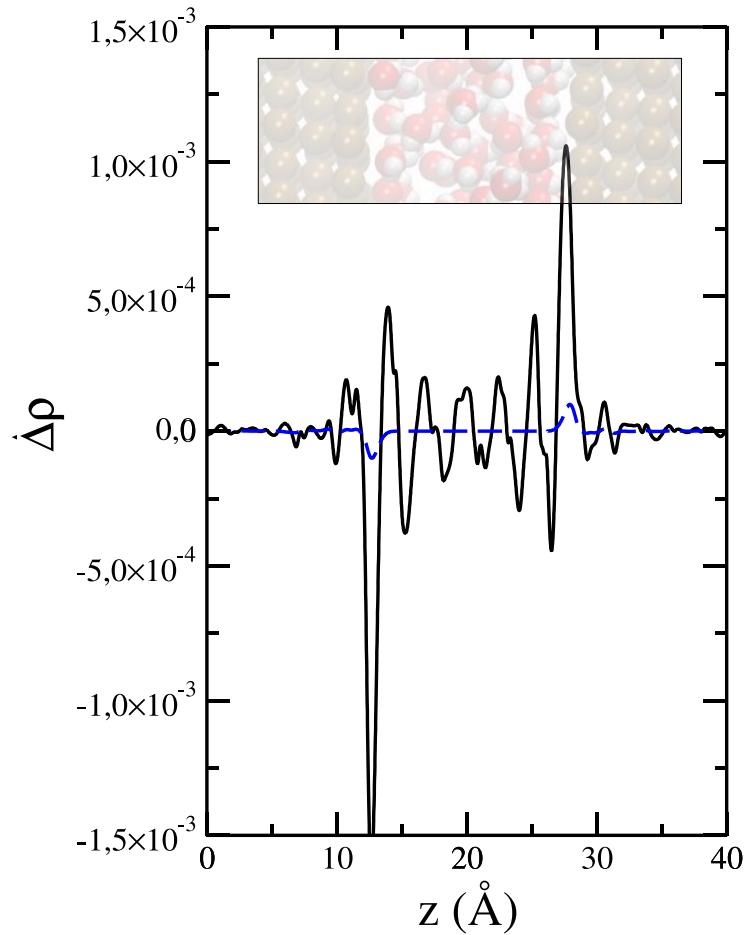
- Stronger layering induced by the bias
- Reorientation of water molecules:
H towards negative; O towards positive electrode



Bias and Surface Charges

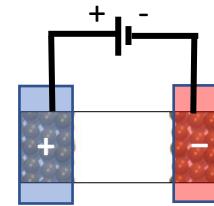
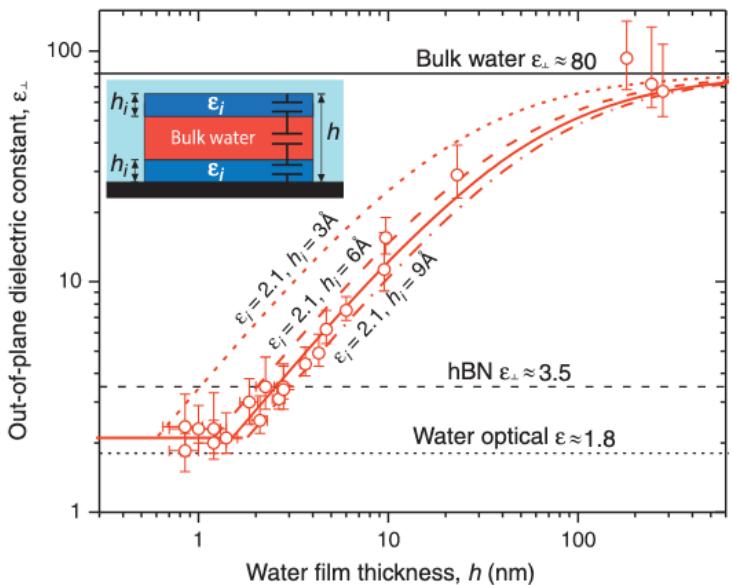

$$\Delta\rho = \rho^V - \rho^{V=0}$$


$$\Delta\rho = \rho^V - \rho_{Au} - \rho_{h2o}$$

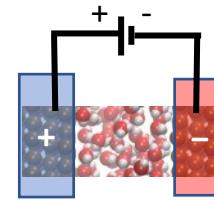


Capacitance and Permittivity

$$V = -0.25 \rightarrow \varepsilon \sim 13$$



$$C = \frac{Q}{V} = \frac{\varepsilon A}{d}$$



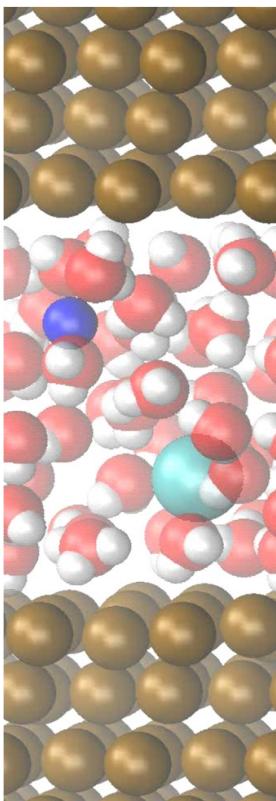
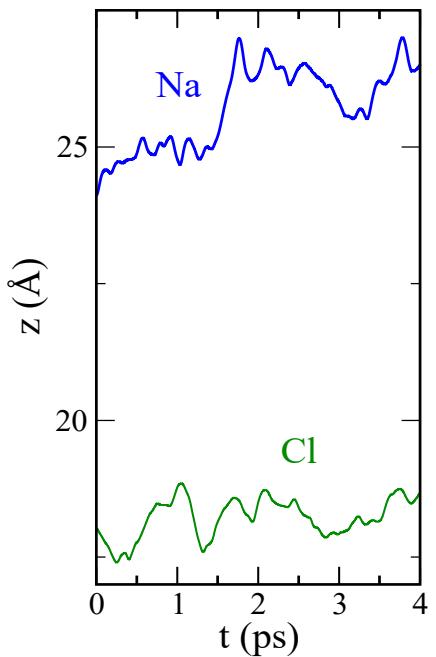
$$\frac{Q}{Q_v} = \varepsilon$$

- Experimentally, for **bulk** water: $\varepsilon \approx 80$
- ε is much reduced for nano-confined water
- Experiment: water **nano-confined** in hBN slits:
 $\varepsilon \sim 3$ for $d \sim 17\text{\AA}$

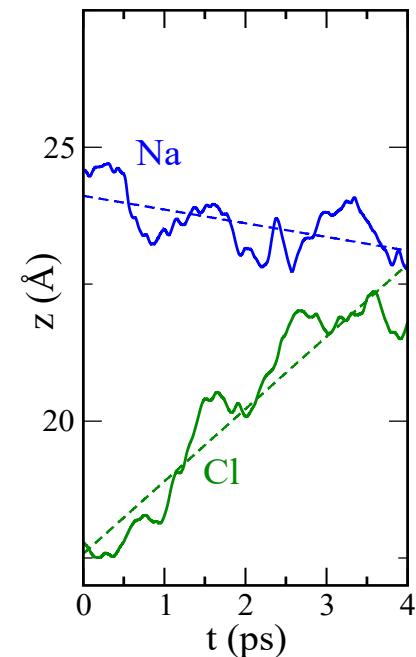
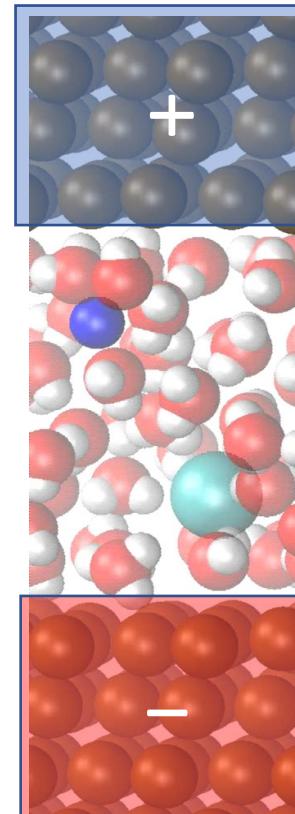
[see e.g. Fumagalli et al, Science **360**, 1339 (2018)]

Ionic species in solution: Cl + Na

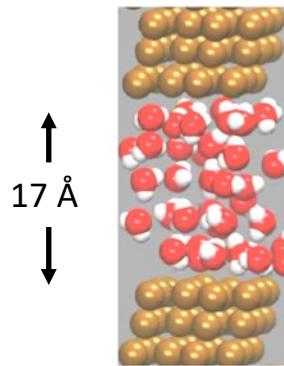
$V = 0 \text{ V}$



$V = 1 \text{ V}$

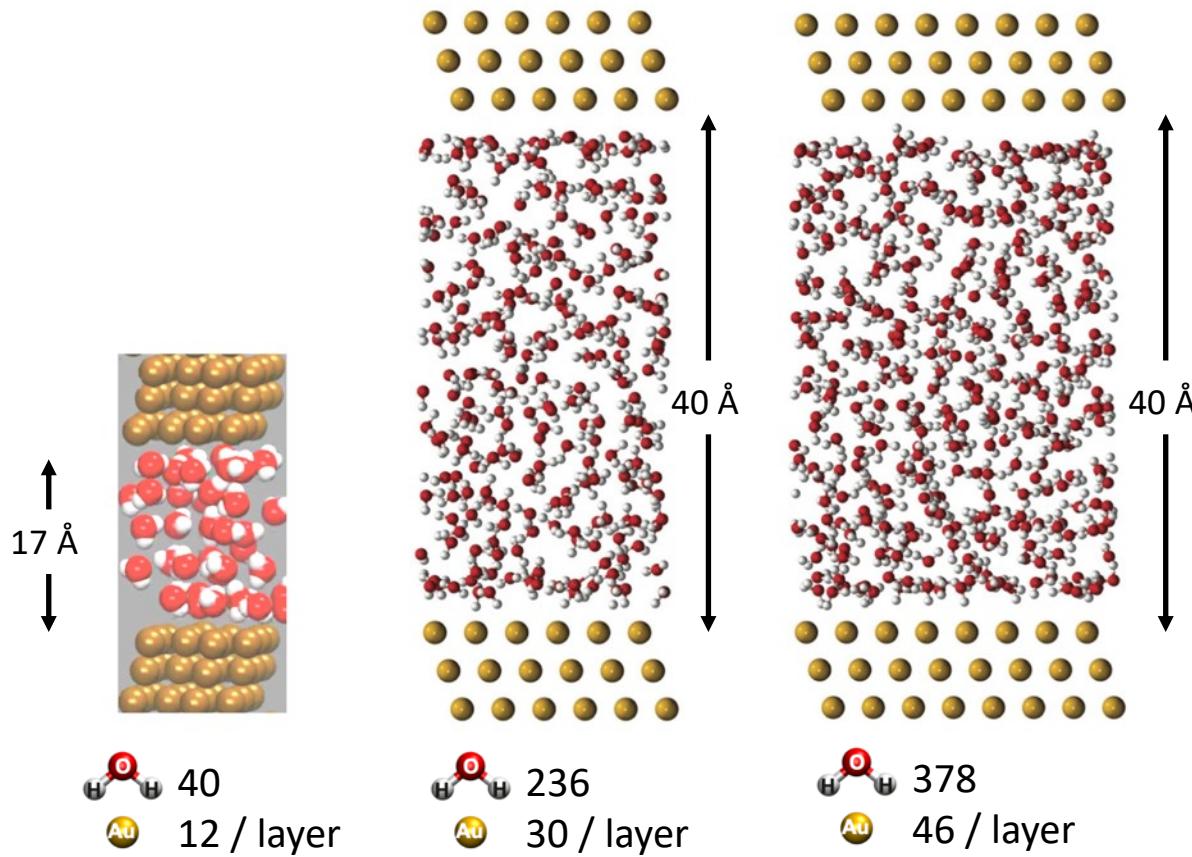


Going to larger systems



40
 12 / layer

Going to larger systems



HPC: Parallelization

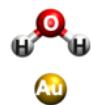
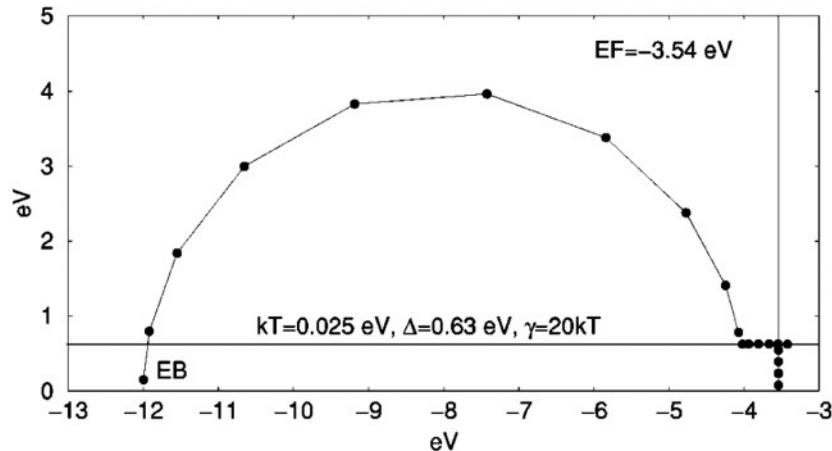
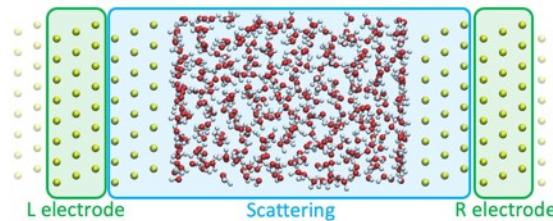


Green's function

$$G(z) = [z - H - \Sigma_L(z) - \Sigma_R(z)]^{-1} \quad (z = \varepsilon + i\eta)$$

Hybrid Parallelization Strategy:

- MPI: energies and k-points over processors (~100)



236
30 / layer

Inversion of
matrices of size
~ 10.000×10.000

HPC: Parallelization

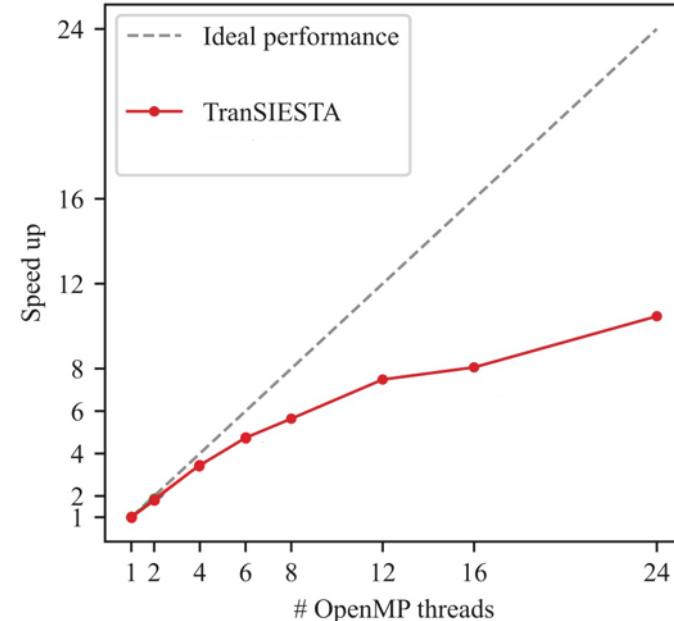


Green's function

$$G(z) = [z - H - \Sigma_L(z) - \Sigma_R(z)]^{-1} \quad (z = \varepsilon + i\eta)$$

Hybrid Parallelization Strategy:

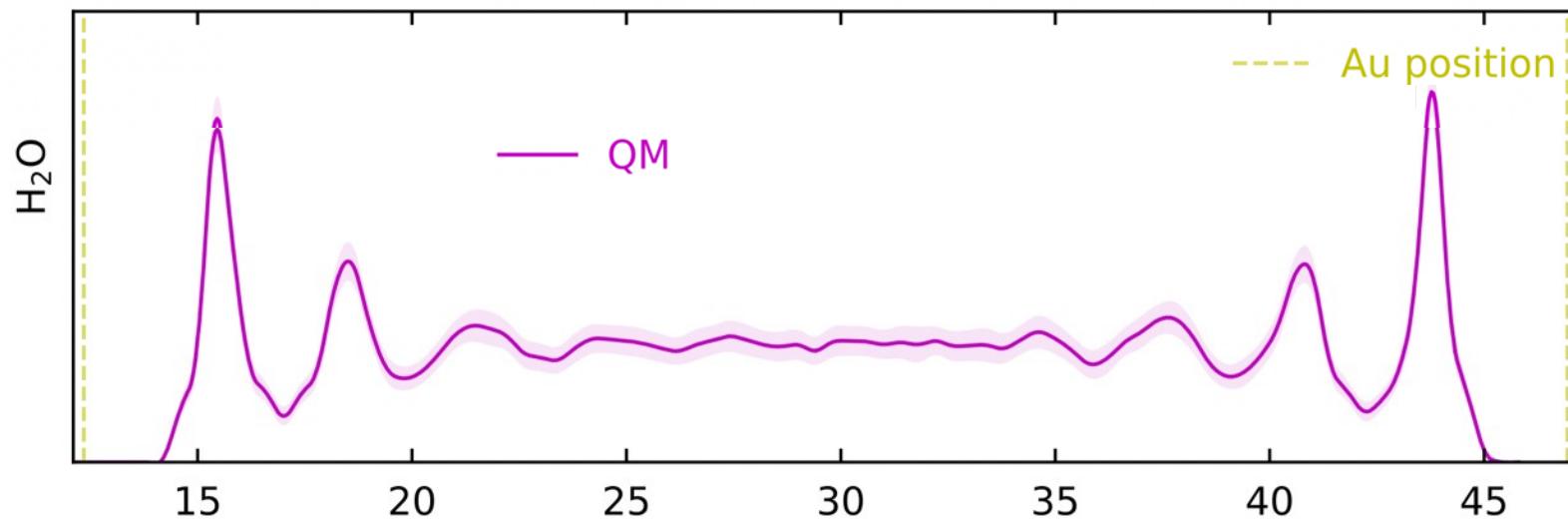
- MPI: energies and k-points over processors (~100)
- OpenMP (shared memory) throughout the code (4-12)
- MPI for distributed inversion (in progress)



Parallel runs in 400 – 1.200 nodes

Massively Parallel runs in 4.000 – 40.000 nodes

Going to larger systems

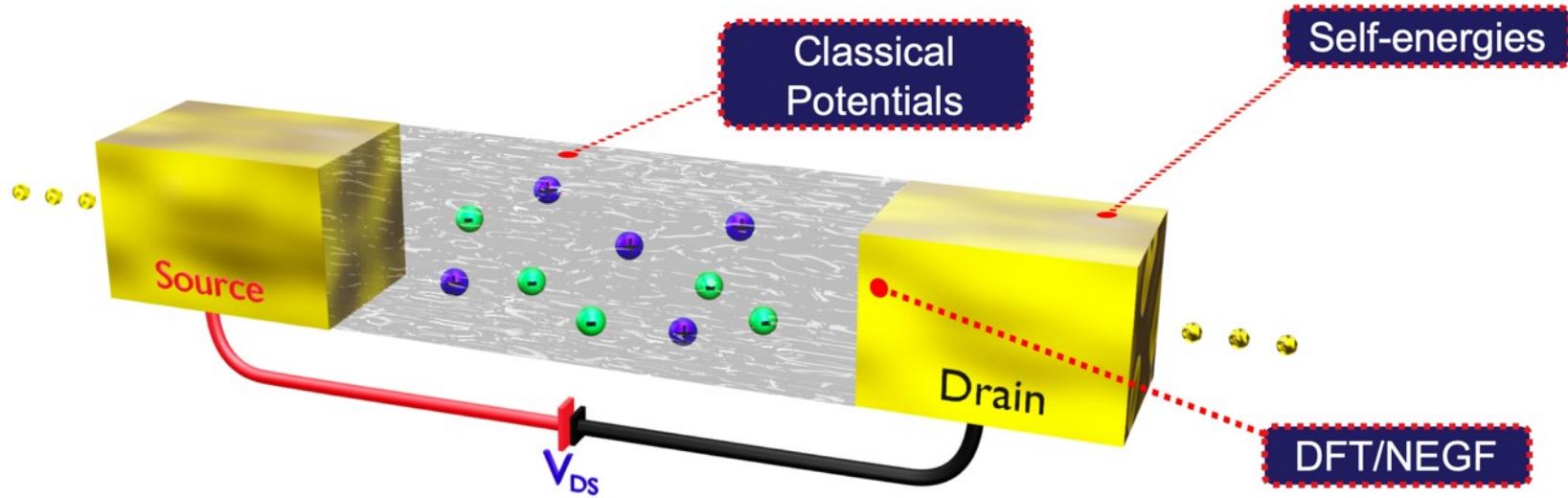


30 ps simulation

Going to larger systems: Hybrid QM/MM approach

Sanz-Navarro et al., Theor. Chem. Acc., 128:825-833 (2011)

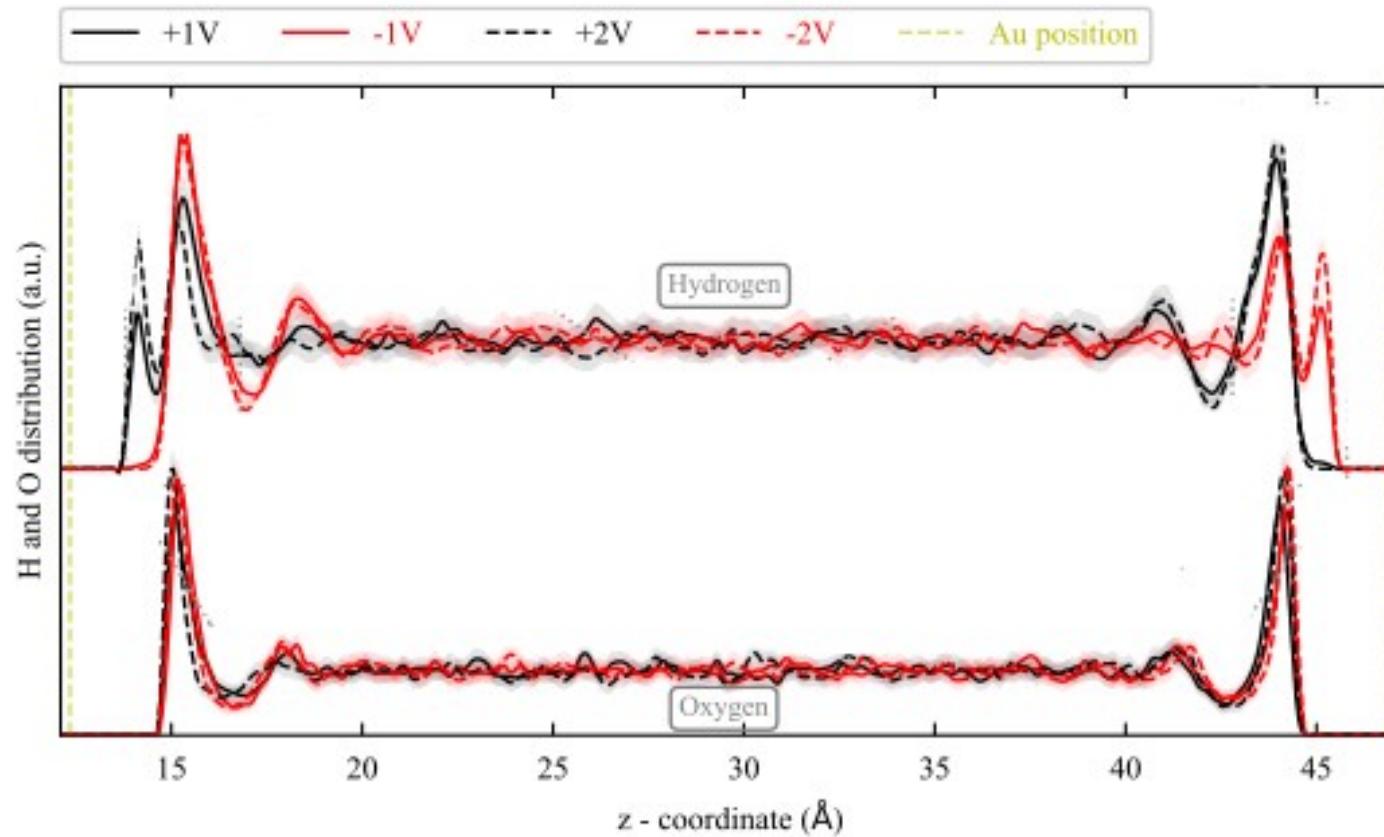
Crespo et al., J. Phys. Chem. B, 107, 13728-13736 (2003)



MM: Solvent, Ions in solution

QM: Metallic electrodes (screening; exchange of electrodes with electrolyte; external bias...)
Chemical species in the electrolyte which undergo chemical reactions

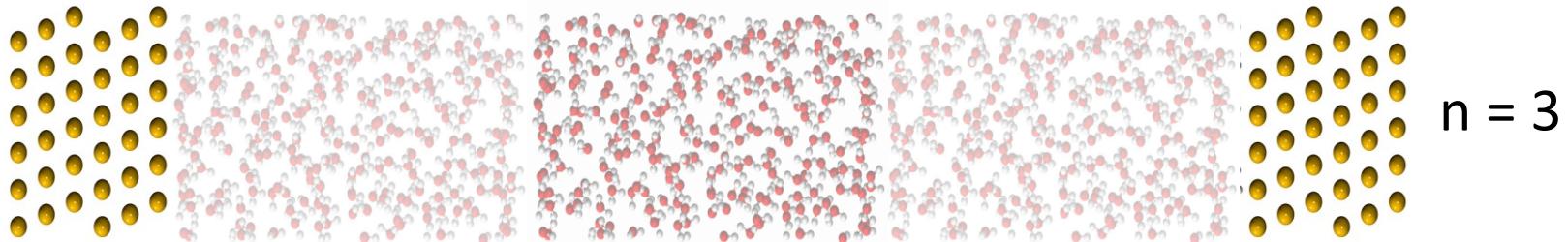
QM/MM: structure of water under bias



Scaling vs system size

Scaling: From 236 to 2360 H₂O molecules

Au – (H₂O)_{236 x n} – Au from n=1 to 10



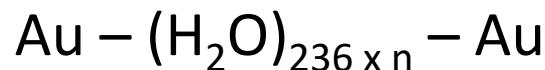
MareNostrum IV @ BSC
Intel Platinum 8160 @ 2.1 GHz

384 cores
96 MPI processes
4 openMP threads / MPI process



Scaling vs system size

Scaling: From 236 to 2360 H₂O molecules

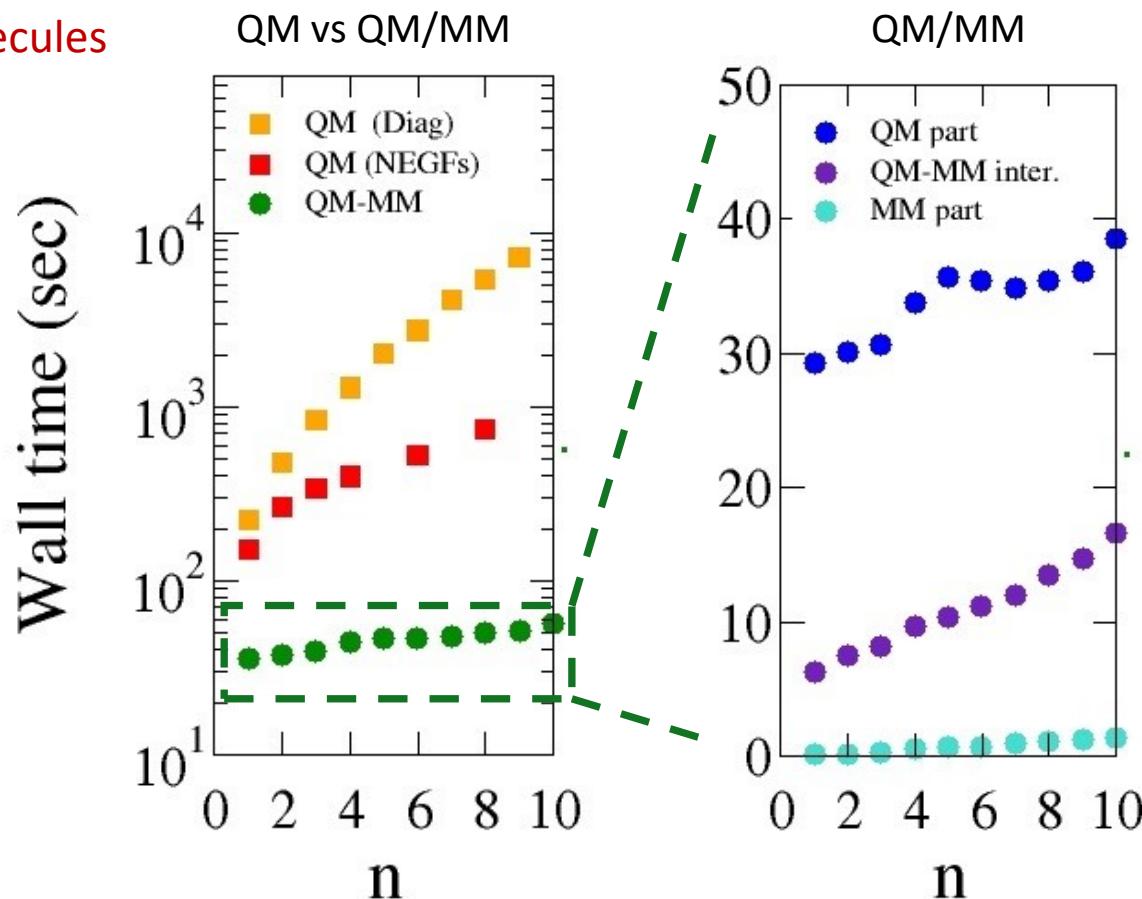


Wall time for 1 MD step

In 384 cores, for n=1:

QM/MM: 1ps / day

QM (NEGF): 0.1ps / day



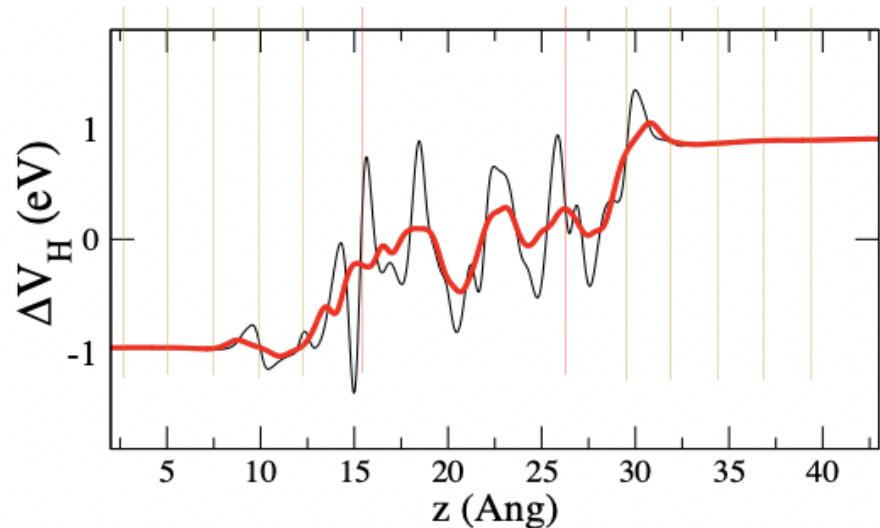
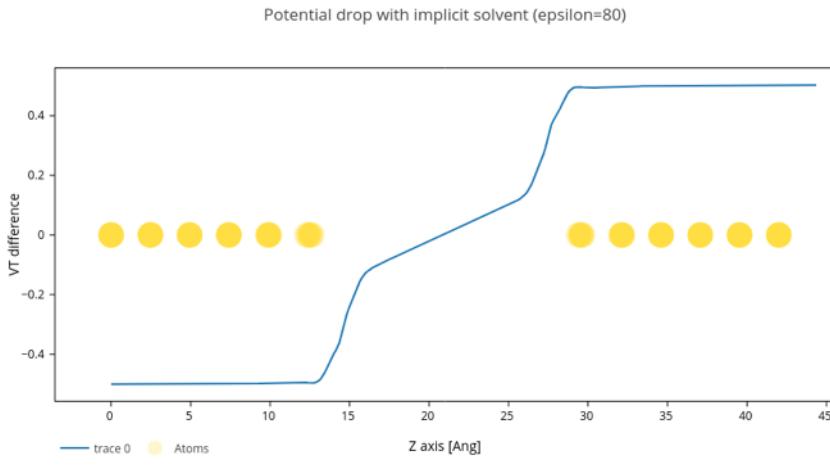
Going beyond - Continuum solvents

A generalized Poisson and Poisson-Boltzmann solver for electrostatic environments 🛒

G. Fisicaro  ; L. Genovese; O. Andreussi; N. Marzari; S. Goedecker

J. Chem. Phys. 144, 014103 (2016)

- Dielectric medium - Generalized Poisson equation (neutral solvent)
- Poisson – Boltzmann equation (ionic solution)



Summary

- **TranSIESTA for spintronics:** Spin-orbit coupling and non-collinear spin states
 - Analysis of the transmission in spin channels; spin flipping processes
-
- **TranSIESTA for electrochemistry:** a tool to study electrified solid / liquid interfaces from first principles
 - Proof of concept: Water on Gold
 - Combining QM with MM: speed-up and scale-up
 - Efficient parallelization for hundreds of cores; near future - tens of thousands
 - Further multiscale approaches (continuous solvents, Poisson-Boltzmann) in progress

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(now at DTU)



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Supercomputing resources at
MareNostrum IV

