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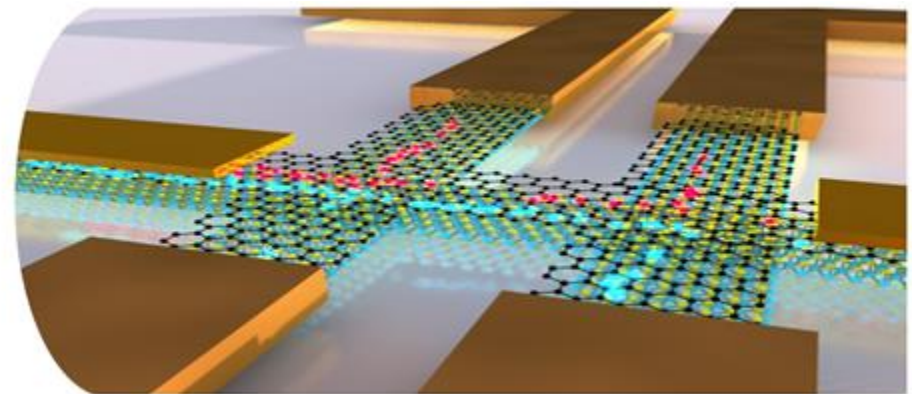
LSQUANT: Linear Scaling Quantum Transport Methodologies

José H. Garcia

Barcelona – Jun 12th, 2022

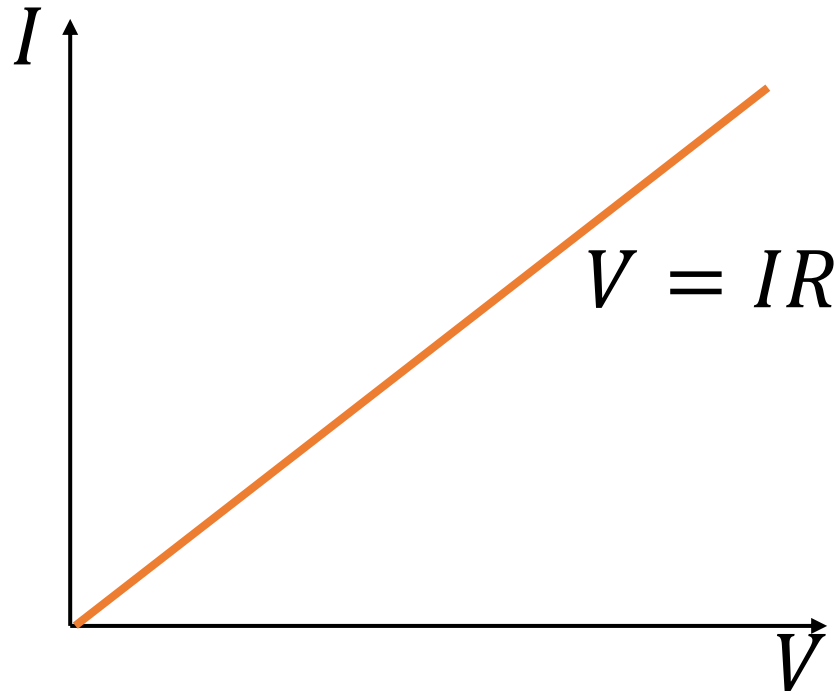
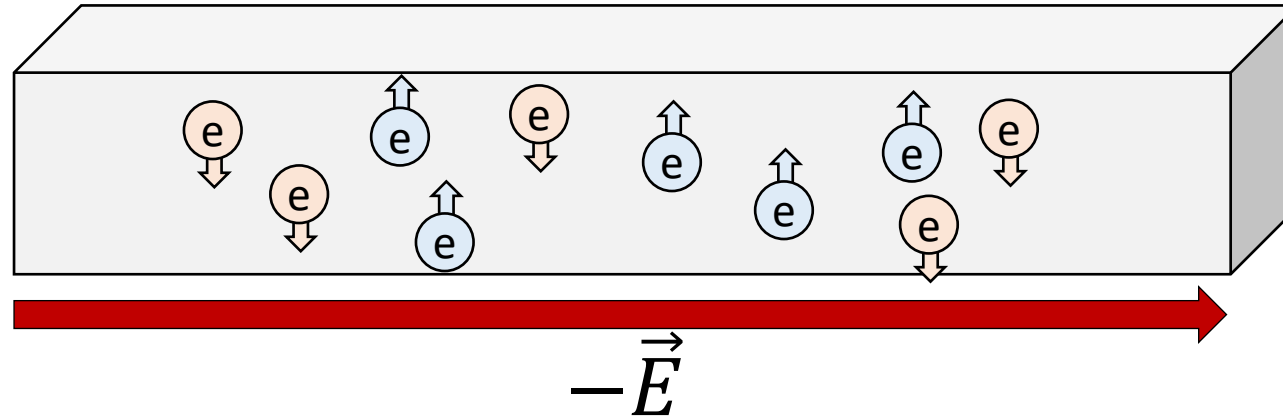
The goals of this presentation

- Gain insights into the importance of electrical response in modern materials science.
- Understand the relationship between the Kubo Formula and electrical response.
- Learn how to numerically solve the Kubo formula in a linear-scaling algorithm.
- Present the up-and-coming LSQuant package

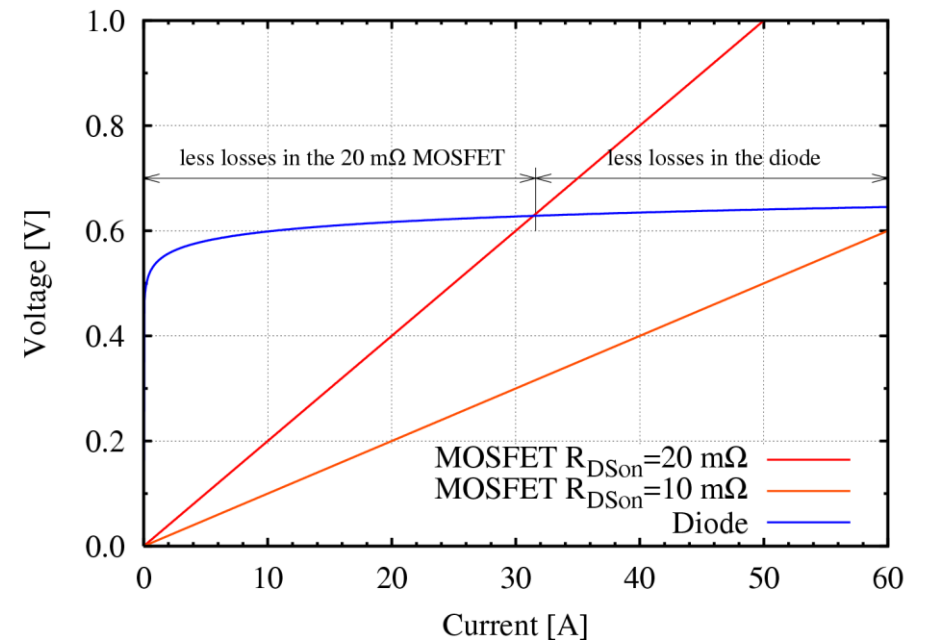


<https://www.lsquant.org/lq/index.php>

Electrical Response

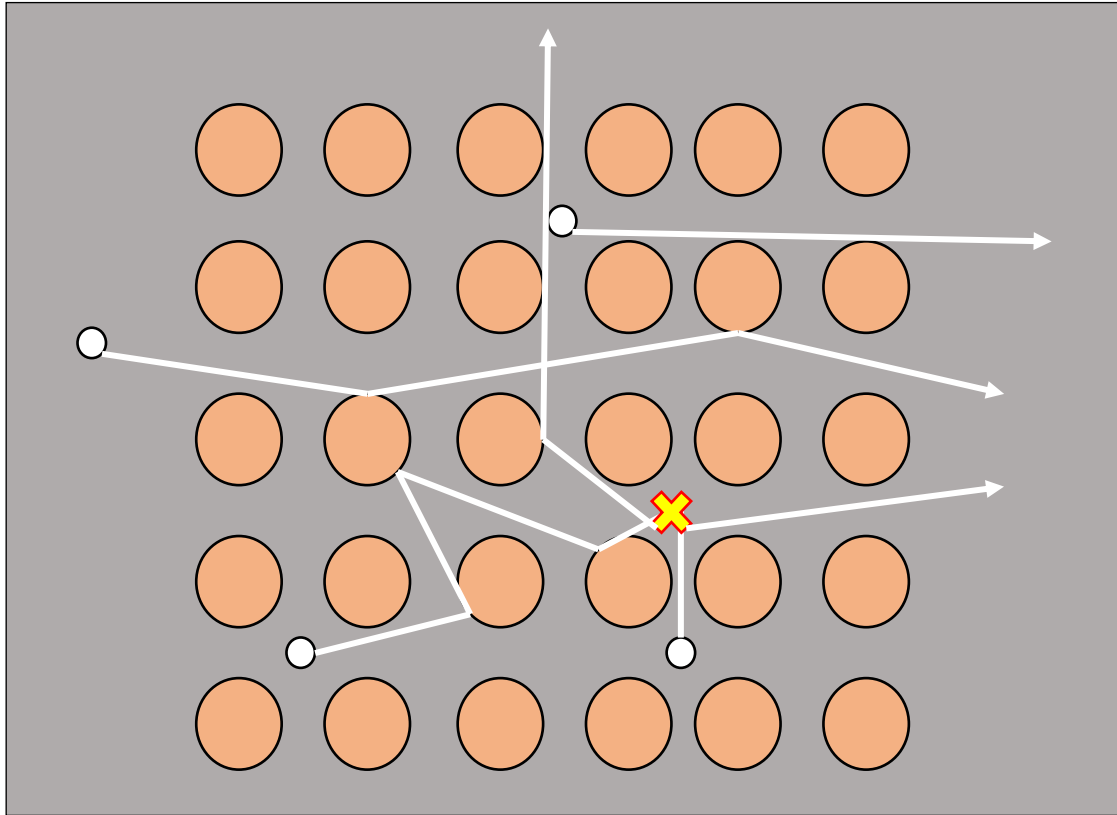


Linear-response



Nonlinear devices

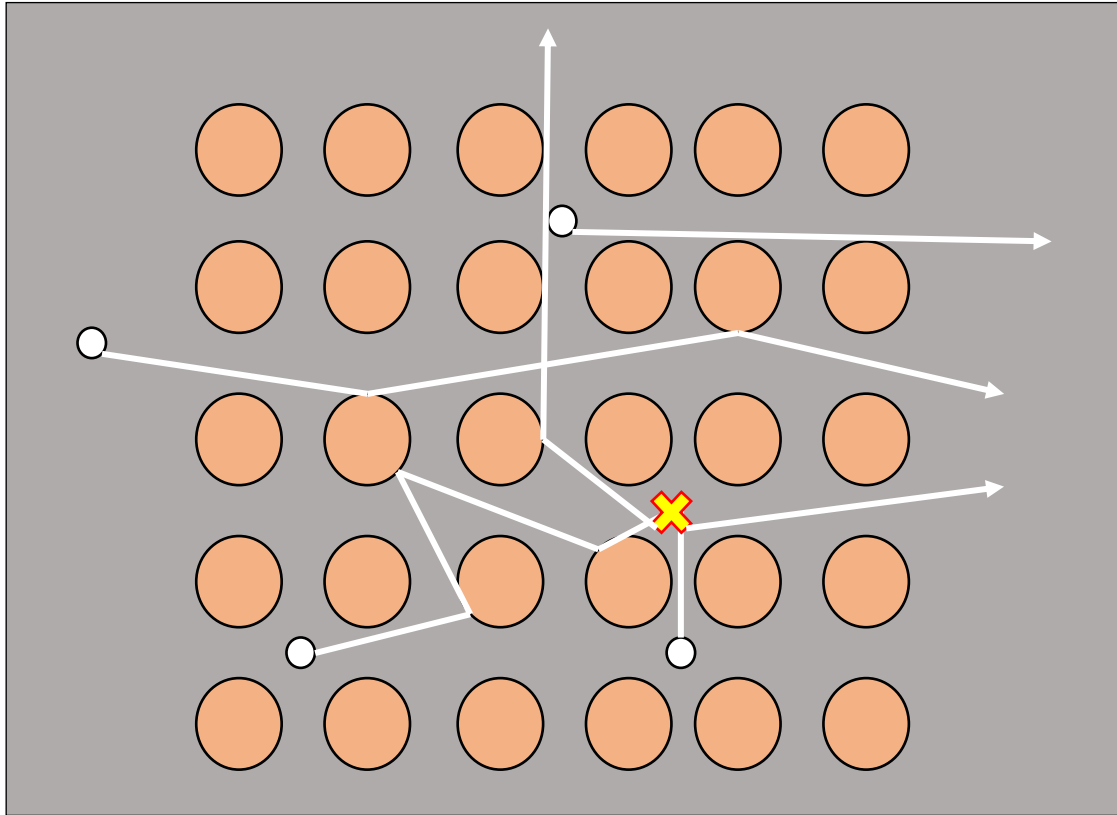
Dissipation



$$\frac{d\vec{p}}{dt} = q\vec{E}$$

Unlimited momentum increment

Dissipation



$$\frac{d\vec{p}}{dt} = q\vec{E} - \frac{\vec{p}}{\tau}$$

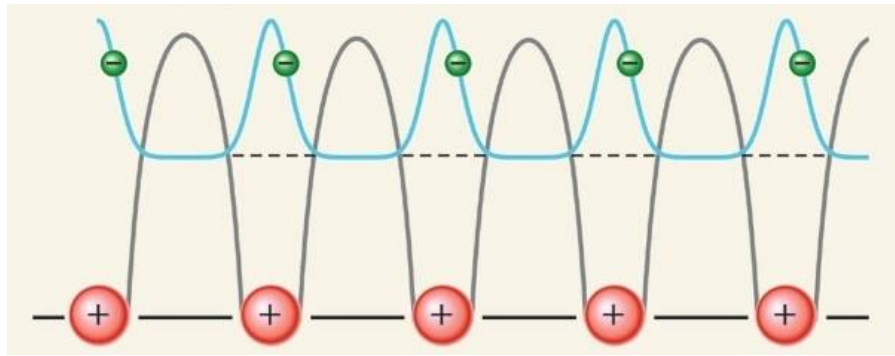
~~Unlimited momentum increment~~

τ

momentum relaxation time

$$\vec{J} = \frac{nq^2\tau}{m} \vec{E} \text{ (Drude model)}$$

Quantum mechanics

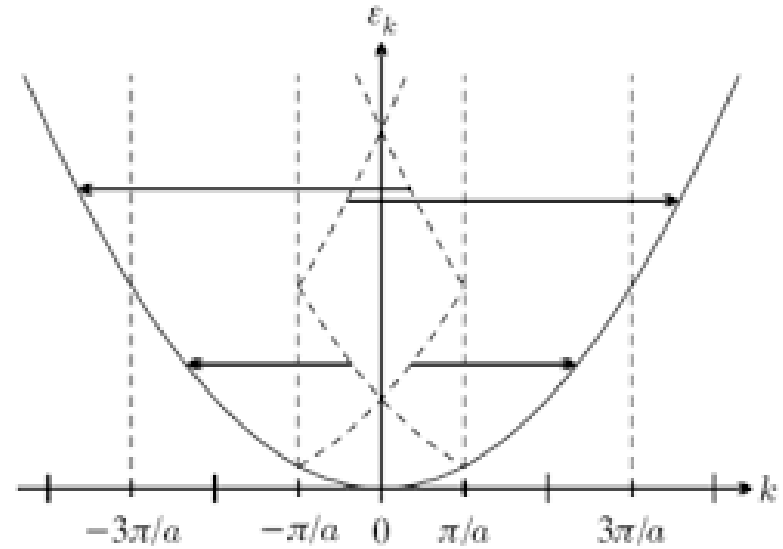


Periodic potential



Bloch's theorem

Nearly free electrons



$$\vec{p} \rightarrow \langle \vec{p} \rangle$$
$$m_e \rightarrow m(\vec{p})$$

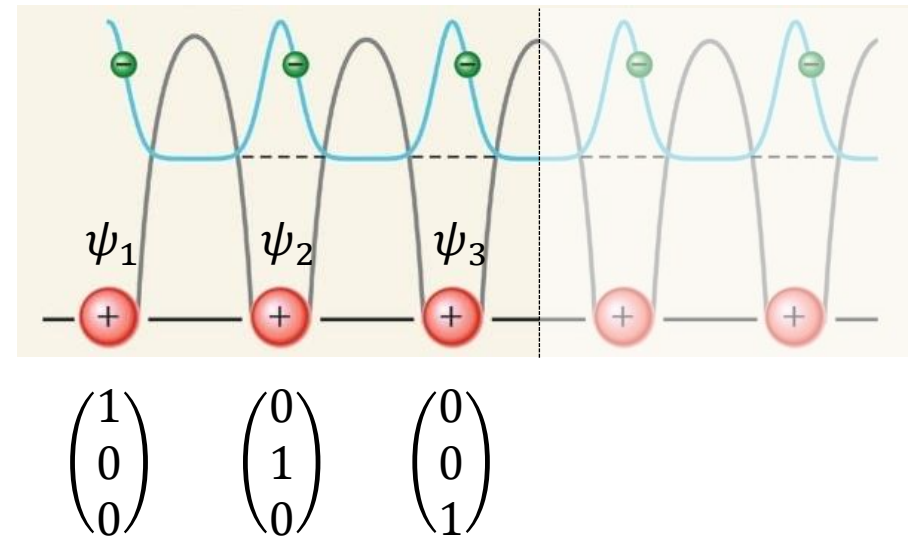
Quantum mechanics



$$i\hbar \frac{d|\Psi(t)\rangle}{dt} = H|\Psi(t)\rangle$$

$$H = \begin{pmatrix} \epsilon_0 & t_{01} & 0 \\ t_{10} & \epsilon_1 & t_{12} \\ 0 & t_{21} & \epsilon_2 \end{pmatrix}$$

$$|\Psi\rangle = \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \end{pmatrix}$$

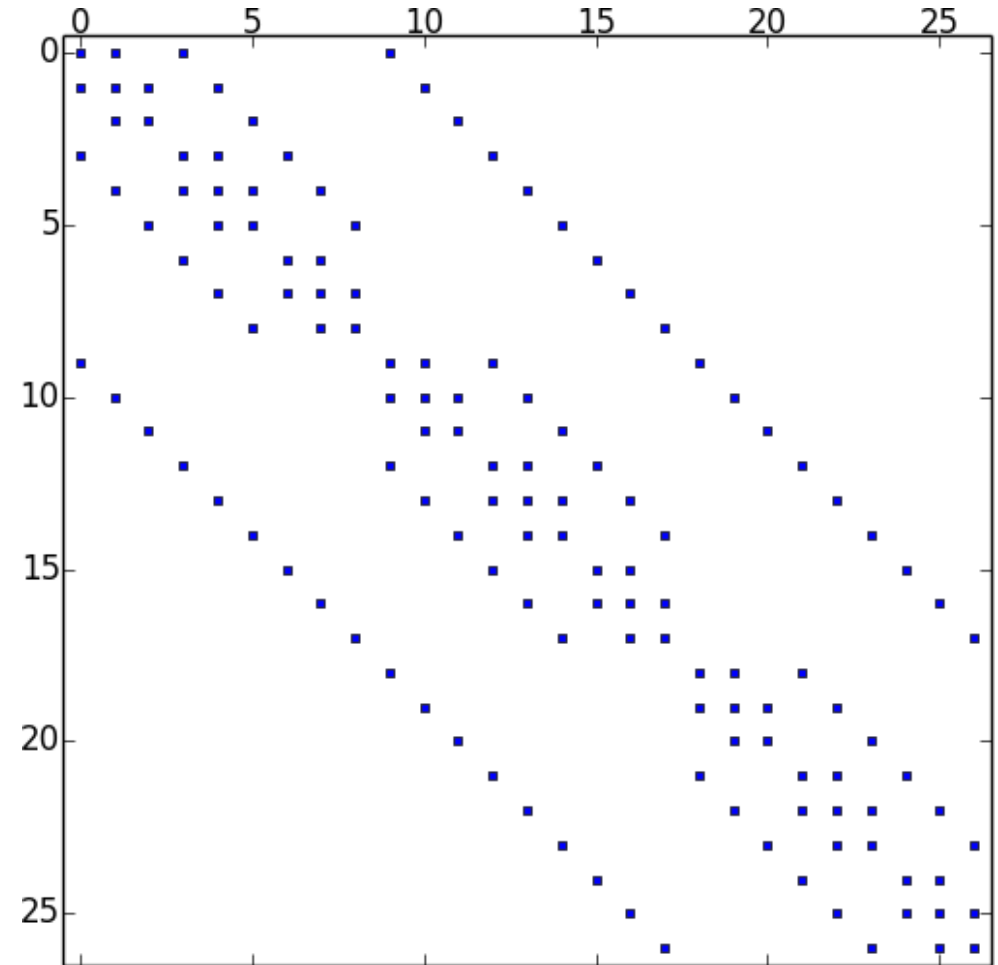


The Hamiltonian as Sparse matrix

$$H = \begin{pmatrix} \epsilon_0 & \cdots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & \epsilon_d \end{pmatrix} \quad d \times d \text{ elements}$$

Column	Row	Matrix element
$\begin{pmatrix} 0 \\ 0 \\ 1 \\ \vdots \\ d \end{pmatrix}$	$\begin{pmatrix} 0 \\ 1 \\ 1 \\ \vdots \\ d \end{pmatrix}$	$\begin{pmatrix} \epsilon_0 \\ t_{01} \\ \epsilon_1 \\ \vdots \\ \epsilon_d \end{pmatrix}$

3 d elements



The macroscopic A

$$i\hbar \frac{d|\Psi(t)\rangle}{dt} = H|\Psi(t)\rangle$$

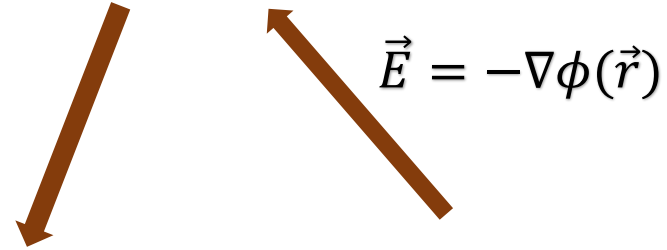
$$\langle \hat{A}(t) \rangle = \langle \Psi(t) | \hat{A} | \Psi(t) \rangle \equiv \text{Tr}[\hat{A}\rho(t)] \quad \text{Density matrix}$$

$$i\hbar \frac{d\rho}{dt} = [H, \rho]$$

Von Neumann equation

The electrical response of A

$$\langle \hat{A}(t, \vec{E}) \rangle \equiv \text{Tr}[\hat{A} \rho(H, \vec{E})]$$


$$\vec{E} = -\nabla\phi(\vec{r})$$

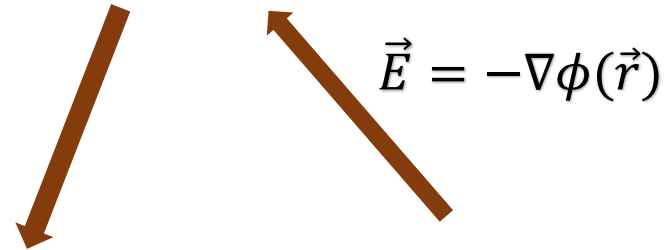
$$H = H_0 + \int d^3r n(\vec{r})\phi(\vec{r})$$

Disordered
Hamiltonian

Minimal
coupling

The electrical response of A

$$\langle \hat{A}(t, \vec{E}) \rangle \equiv \text{Tr}[\hat{A} \rho(H, \vec{E})]$$


$$\vec{E} = -\nabla\phi(\vec{r})$$

$$H = H_0 + \int d^3r n(\vec{r})\phi(\vec{r})$$

Disordered
Hamiltonian

Minimal
coupling

Strong electric fields could
change the ground states!

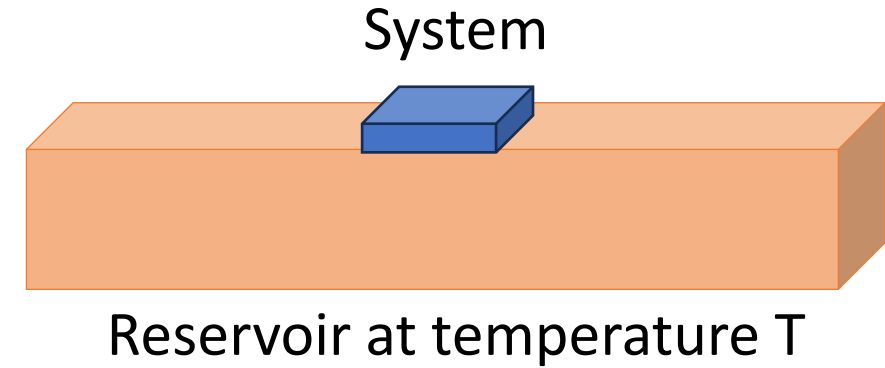
Linear response theory: The Kubo formula

Principal goal:

$$\rho(H, \vec{E}) = \rho_0(H) + \left(\nabla_{\vec{E}} \rho(\vec{E}, H) \right)_{\vec{E}=0} \vec{E}$$

Fundamental Hypotheses

- Thermal equilibrium: $\rho_0(H) \rightarrow \rho_{\text{eq}}(T, \mu)$



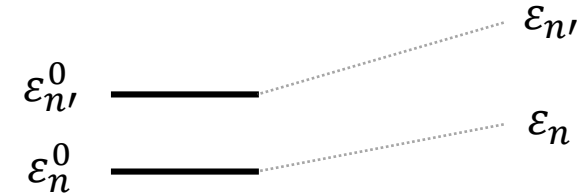
Linear response theory

Principal goal:

$$\rho(H, \vec{E}) = \rho_0(H) + \left(\nabla_{\vec{E}} \rho(\vec{E}, H) \right)_{\vec{E}=0} \vec{E} \chi_{\hat{E}}$$

Fundamental Hypotheses

- Thermal equilibrium: $\rho_0(H) \rightarrow \rho_{\text{eq}}(T, \mu)$
- Adiabatic evolution: $|\varepsilon_n^0\rangle \rightarrow |\varepsilon_n(t)\rangle$



$$\langle \hat{A}(t) \rangle \equiv \text{Tr}[\hat{A}(t) \rho_{\text{eq}}] + \text{Tr}[\hat{A}(t) \chi_{\hat{E}}]$$

Linear response theory: The Kubo formula

$$\langle \hat{A}(t) \rangle = \lim_{t, \tau \rightarrow \infty} \int_0^\beta d\beta' \int_0^t d\tau' e^{-\tau'/\tau} \langle \rho_{eq} \hat{A} [E_\alpha J_\alpha(\tau' - i\hbar\beta')] \rangle$$

- ρ_{eq} : Density matrix in thermal equilibrium
- $E_\alpha J_\alpha(\tau' - i\hbar\beta')$: **History** of the electric field coupled to the current density given a **thermal broadening**

Kubo, Ryogo (1957). "Statistical-Mechanical Theory of Irreversible Processes. I. General Theory and Simple Applications to Magnetic and Conduction Problems". J. Phys. Soc. Jpn. 12 (6): 570–586. doi:10.1143/JPSJ.12.570.⁹

Noninteracting Kubo Formula

$$\langle \hat{A} \rangle = i\hbar \int dE \operatorname{Tr} \left[f(E' - \mu) \delta(E' - H) \hat{A} \frac{dG^+(E')}{dE} \vec{V} \cdot \vec{E}_0 \right]$$

Noninteracting Kubo Formula

$$\langle \hat{A} \rangle = i\hbar \int dE \operatorname{Tr} \left[f(E' - \mu) \delta(E' - H) \hat{A} \frac{dG^+(E')}{dE} \vec{V} \cdot \vec{E}_0 \right]$$

Representable in tight-binding basis

$$H = \begin{pmatrix} \epsilon_0 & t_{01} & 0 \\ t_{10} & \epsilon_1 & t_{12} \\ 0 & t_{21} & \epsilon_2 \end{pmatrix} \quad |\Psi\rangle = \begin{pmatrix} \psi_0 \\ \psi_1 \\ \psi_2 \end{pmatrix}$$

Noninteracting Kubo Formula

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Representable in tight-binding basis

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$$\langle \psi_n | V_x | \psi_m \rangle = i(X_n - X_m) H_{nm}$$

Noninteracting Kubo Formula

$$\langle \hat{A} \rangle = i\hbar \int dE \operatorname{Tr} \left[f(E' - \mu) \delta(E' - H) \hat{A} \frac{dG^+(E')}{dE} \vec{V} \cdot \vec{E}_0 \right]$$

Singular Functions

$$G^+(E') = \lim_{\eta \rightarrow 0} \frac{1}{H - E + i\eta}$$
$$\int dE' \delta(E' - H) f(E') = f(H)$$

Numerical implementation

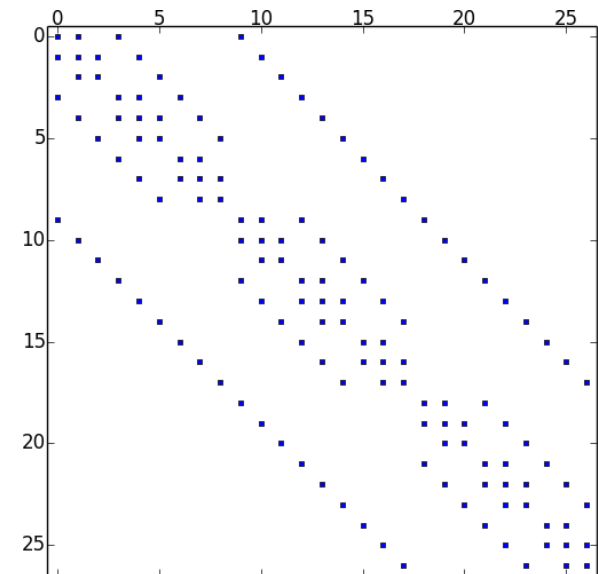
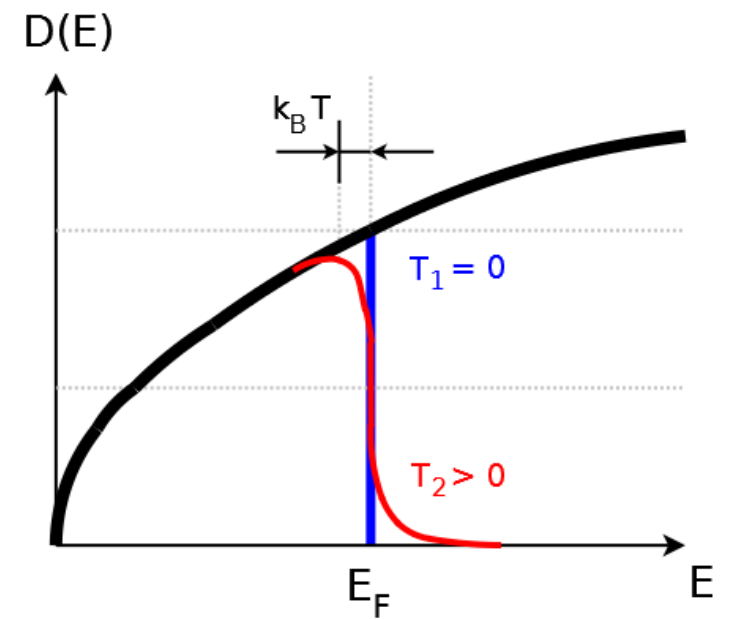
The kernel polynomial method

$$DOS \equiv \text{Tr}[\delta(E' - H)]$$

$$\text{Tr}[\delta(H - E)] = \sum_n \delta(E_n - E)$$

Polynomial expansion

$$\delta(H - E) = \sum_{n=0}^{\infty} c_n p_n(H)$$



The kernel polynomial method

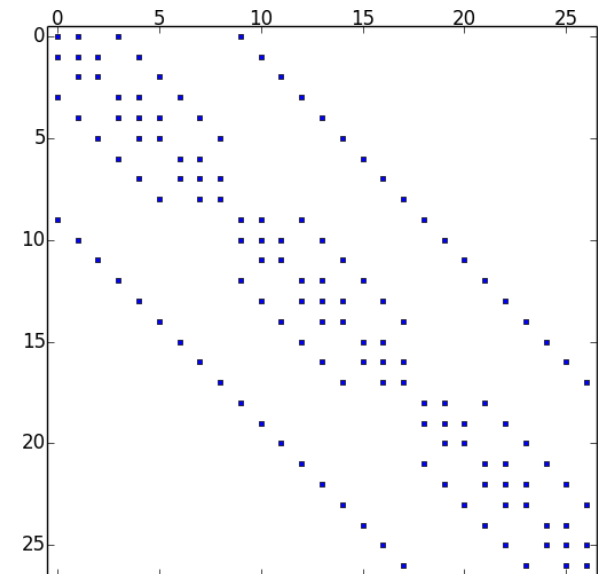
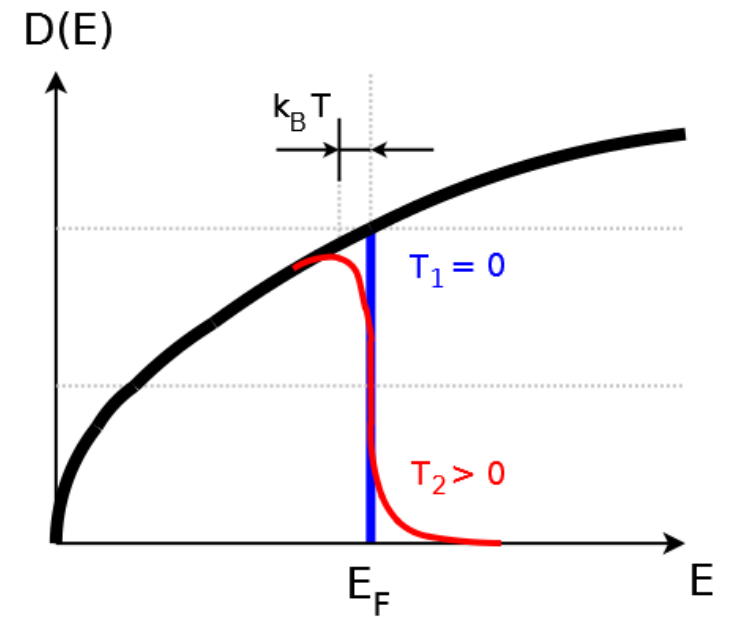
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Polynomial expansion

$$\delta(H - E) = \sum_{n=0}^{\infty} c_n p_n(H)$$

$$O(N^3) \rightarrow O(N^2)$$



The random phase approximation

$$DOS \equiv \text{Tr}[\delta(E' - H)]$$

$$DOS = \langle r | \delta(H - E) | r \rangle$$

$$= \sum_{n,m=0}^d \xi_n \xi_m \langle n | \delta(H - E) | m \rangle$$

$$|r\rangle = \begin{pmatrix} \xi_0 \\ \vdots \\ \xi_d \end{pmatrix}$$

$$O(N^2) \rightarrow O(N)$$

Choosing the right distribution

$$\begin{aligned} \xi_n \xi_m &= 0 \\ \xi_n^2 &= 1 \end{aligned}$$

The Chebyshev polynomials

$$\text{Tr}[\delta(H - E)] = \sum_{n=0}^M c_n \langle r | p_n(H) | r \rangle$$

$$T_m(x) \equiv \text{Cos} (m \text{ArcCos} (x))$$

$$T_{n+1}(x) \equiv 2 T_n(x) - T_{n-1}(x)$$

$$T_0(x) \equiv 1,$$

$$T_1(x) \equiv x ,$$

$$T_2(x) \equiv 2x^2 - 1,$$

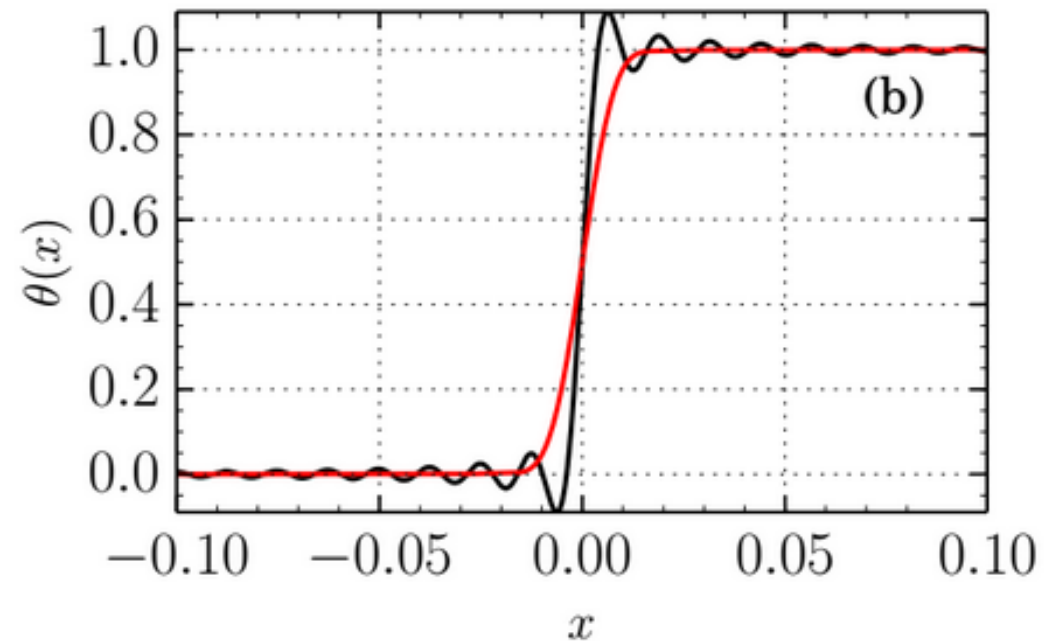
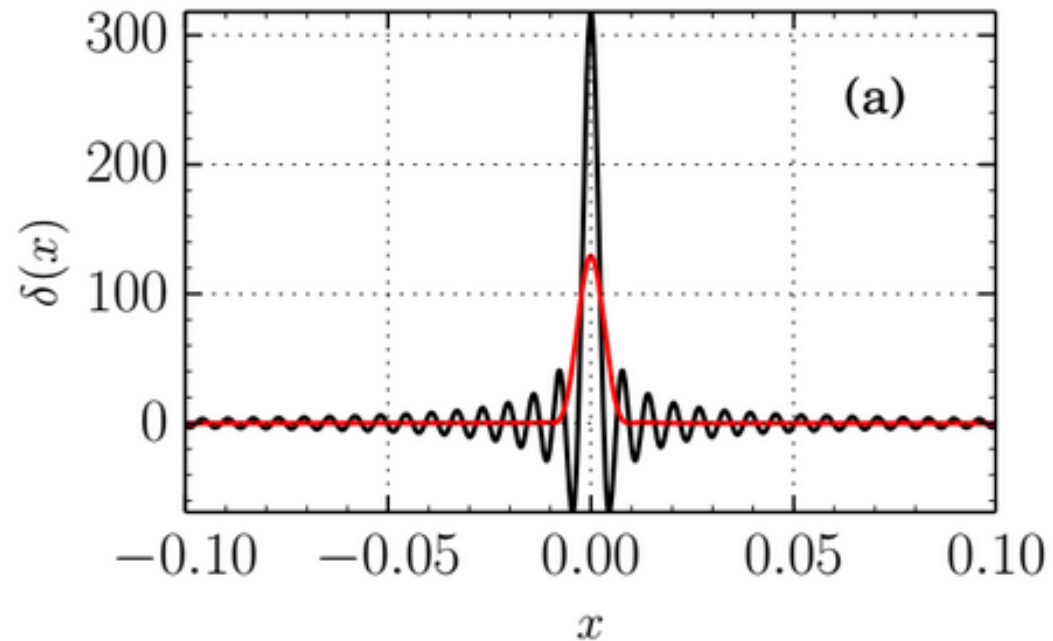
$$T_3(x) \equiv 4x^3 - 2x$$

$$O(N^2) \rightarrow O(M N)$$

The Gibbs oscillations

$$\delta(x - x_0) = \frac{2}{\pi\sqrt{1-x^2}} \sum_{m=0}^M \frac{g_m T_m(x_0) T_m(x)}{1 + \delta_{m0}}$$

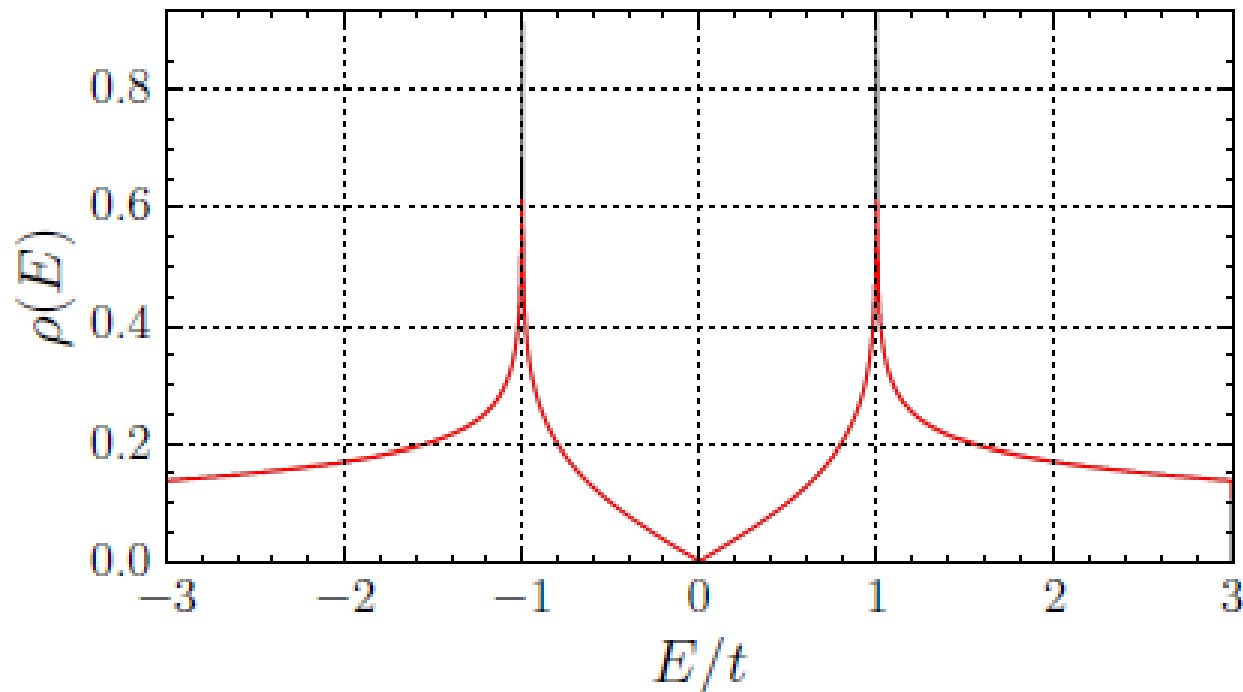
$$g_m = \frac{\sinh\left(\left(1 - \frac{m}{M}\right)\lambda\right)}{\sinh(\lambda)}$$



The final density of states

$$\text{DOS} = \frac{2\alpha}{W_H} \frac{2}{\pi\sqrt{1 - \tilde{\varepsilon}_F}} \sum_{m=0}^M g_m T_m(\tilde{\varepsilon}_F) \mu_m$$

$$\mu_m = \text{Tr}[T_m(\tilde{H})] = \langle\langle \langle r | T_m(\tilde{H}) | r \rangle \rangle\rangle$$



$$\tilde{H}, \tilde{\varepsilon}_F \in (-\alpha, \alpha) \quad \tilde{H} = \frac{2\alpha}{W_H} (H - E_c), \quad \tilde{\varepsilon}_F = \frac{2\alpha}{W_H} (\varepsilon_F - E_c)$$

The final density of states

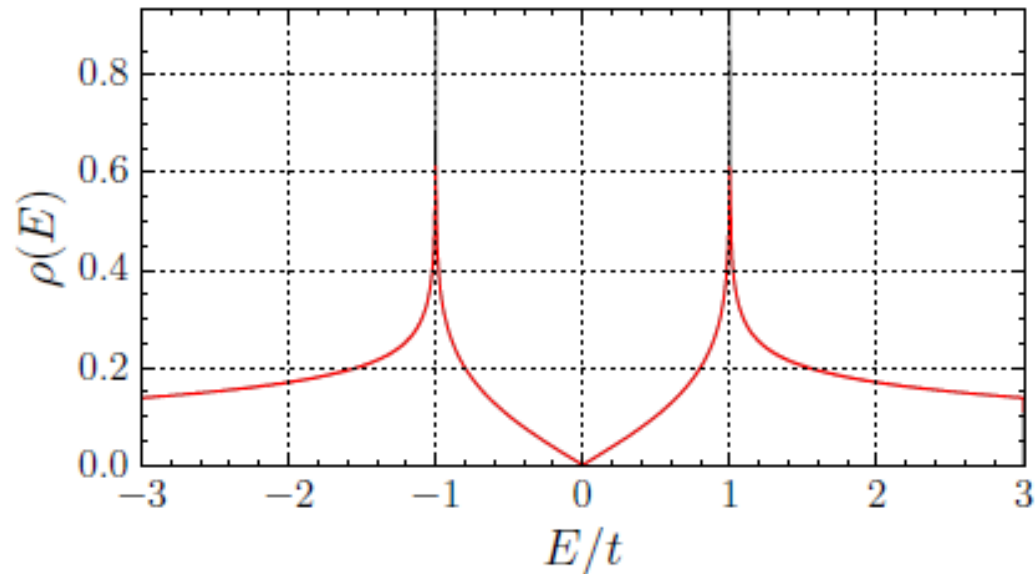
The density of states using KQuant

```
import matplotlib.pyplot as plt
from k_quant.system import System

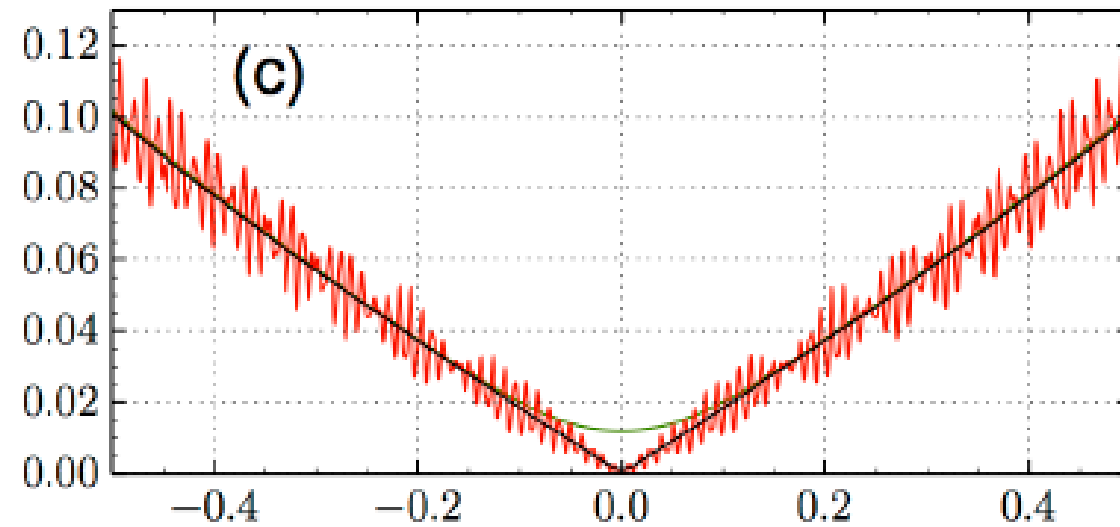
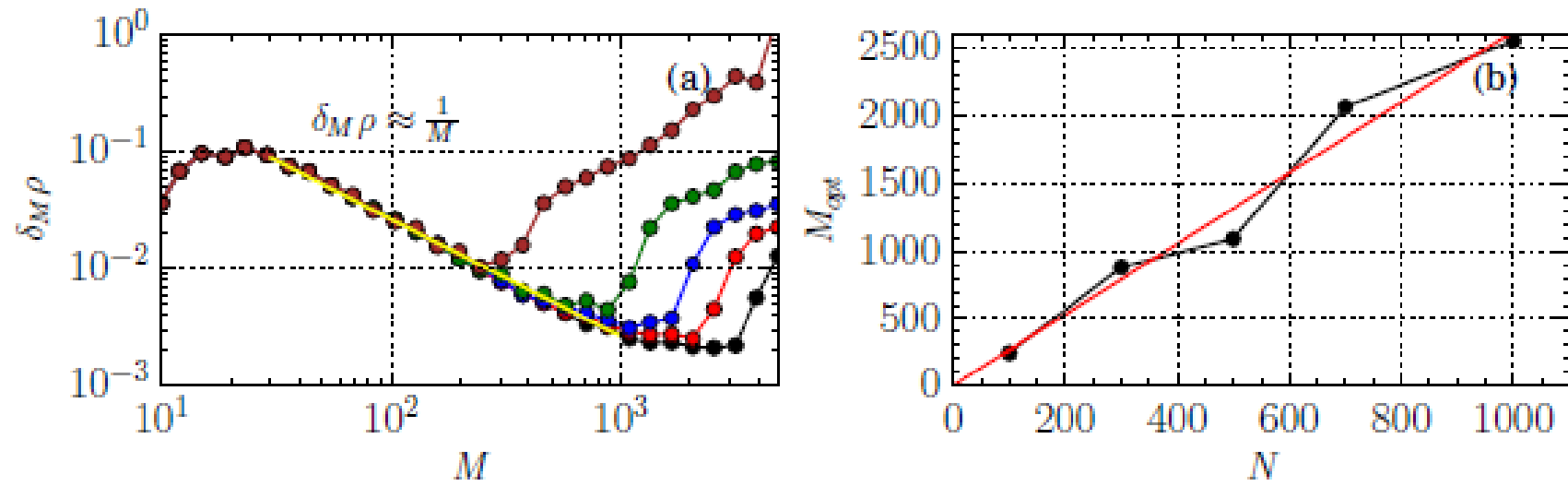
import kpm
kpm.safe_CUTOFF= 0.99;

wann_syst = System( dimensions = (10000,1,1), w90_inp="linear_chain")
dens = kpm.Density(wann_syst, bounds=(-1,1))
dens.ComputeMoments( broadening =0.05)

plt.plot(*dens.spectral_average() )
```



Convergence of the Chebyshev polynomials



The LSQT

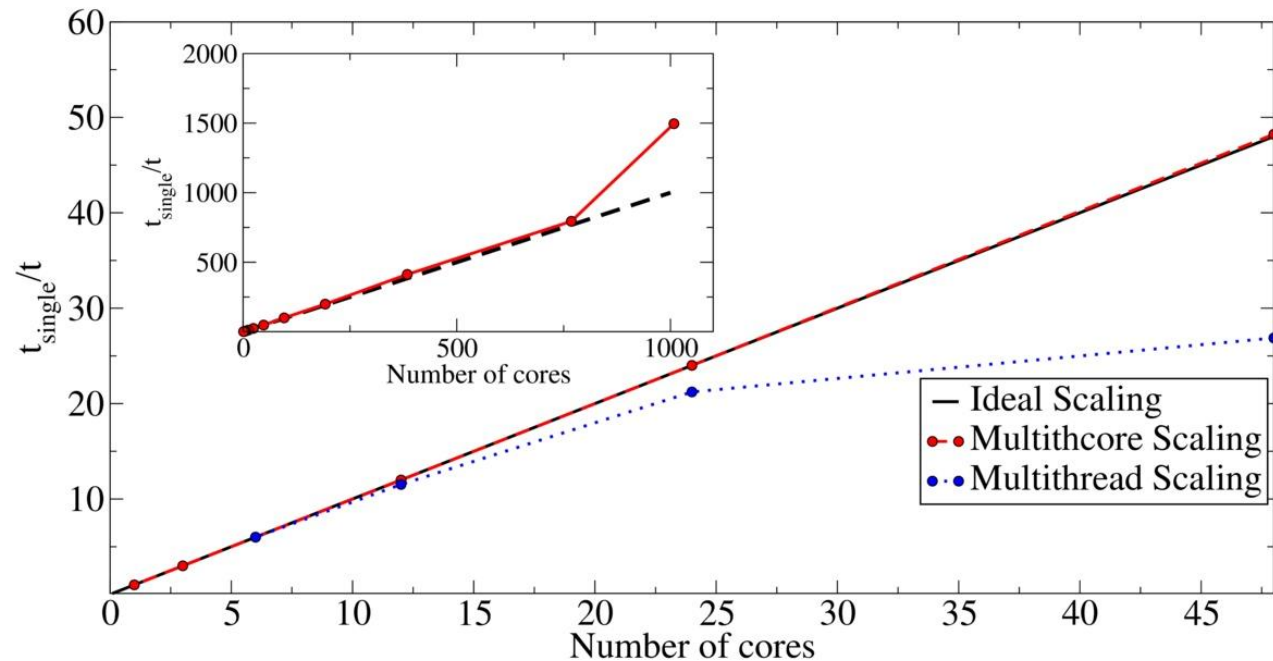
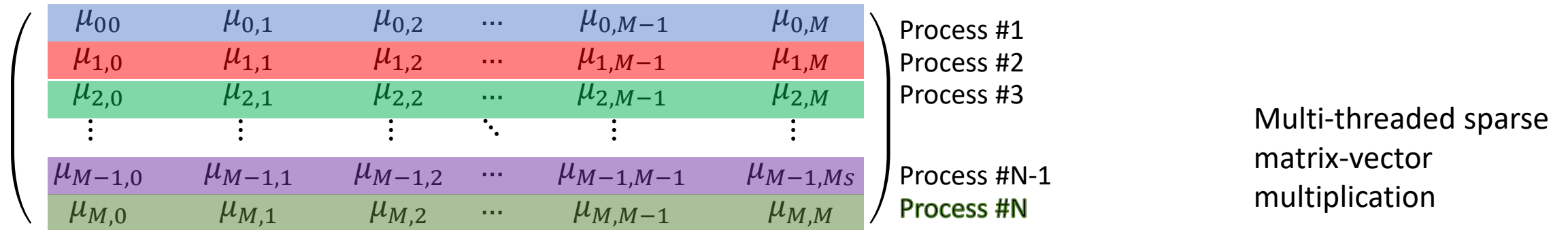
$$\langle \hat{A} \rangle = i\hbar \int dE \operatorname{Tr} \left[f(E' - \mu) \delta(E' - H) \hat{A} \frac{dG^+(E')}{dE} \vec{V} \cdot \vec{E}_0 \right]$$

$$\langle \hat{A}(\mu, T) \rangle = \sum_{m=0}^M \sum_{n=0}^M \Gamma_{m,n}(\mu, T) \operatorname{Tr} \left[T_m(H) \hat{A} T_n(H) \hat{V}_\beta \right] \mu_{m,n}$$

Chebyshev expansion moment matrix

Scaling of the LSQT

$$\mu_{m,n} = \text{Tr} [T_m(H)j_x T_n(H)j_x]$$

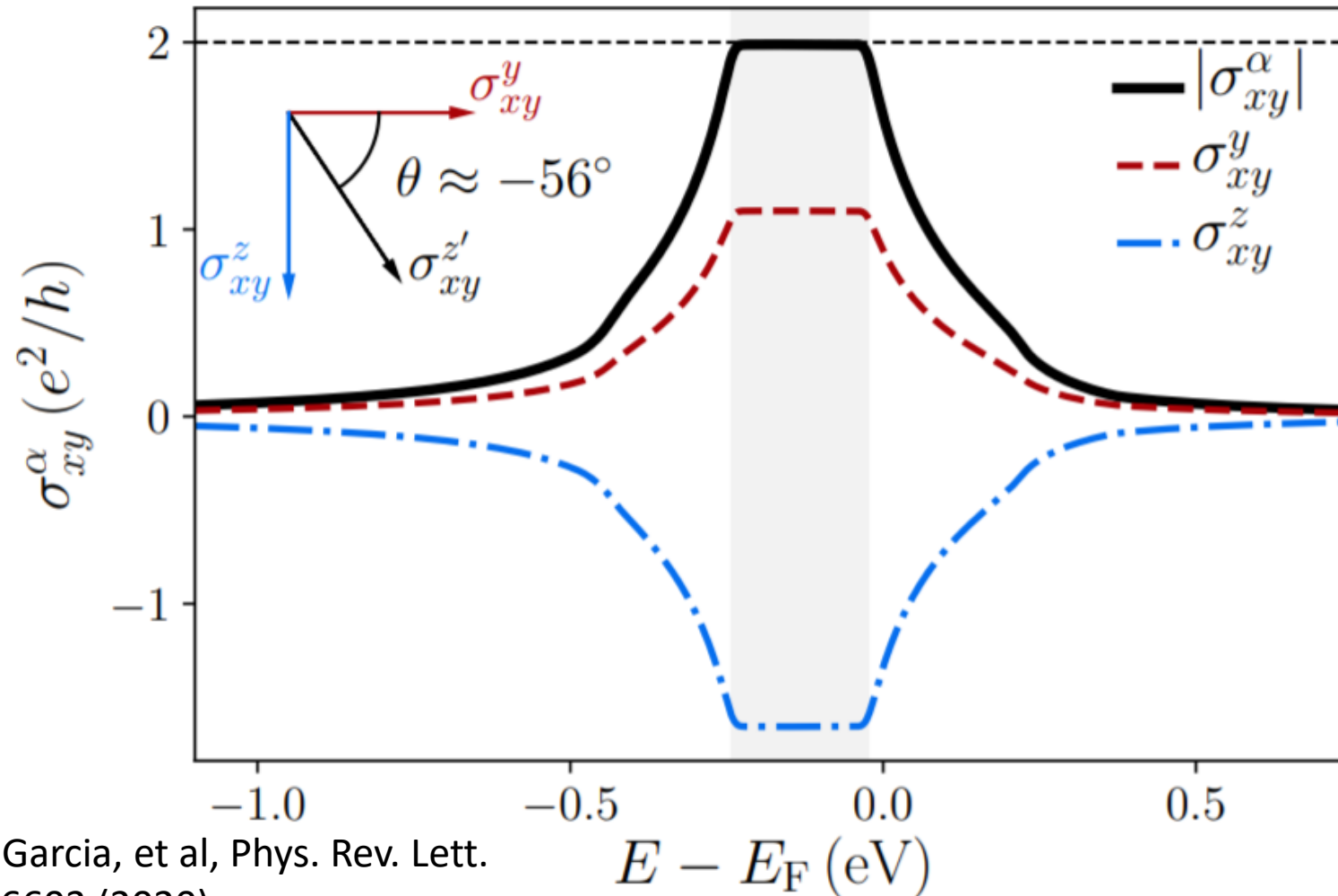


Applications

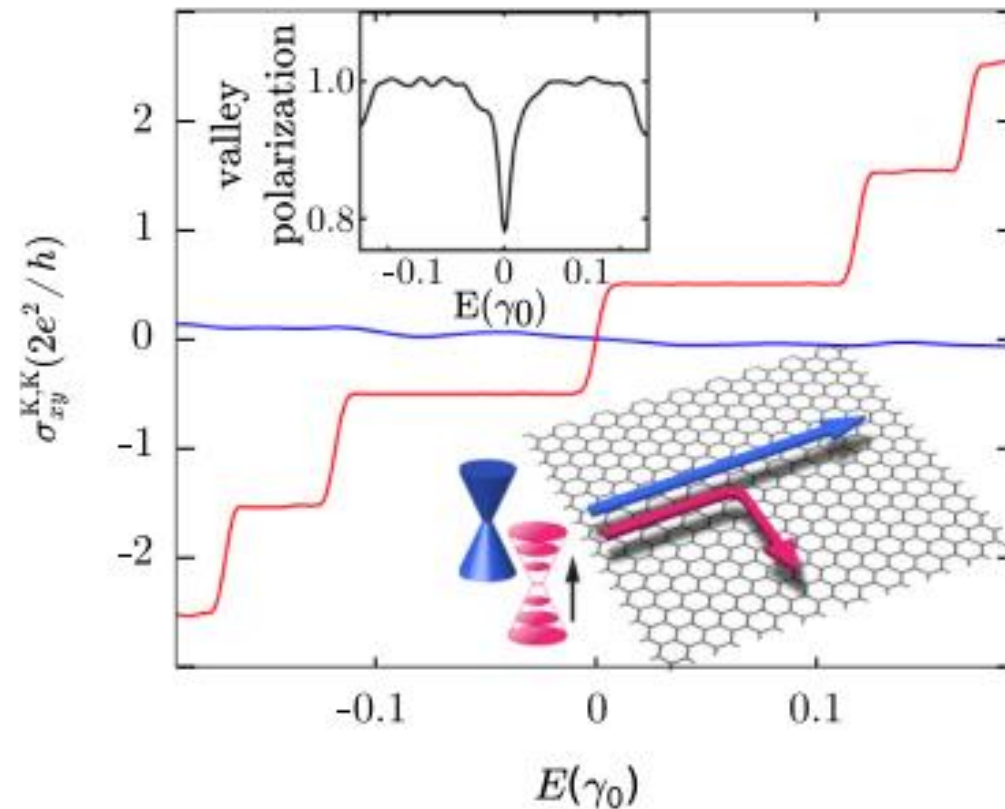
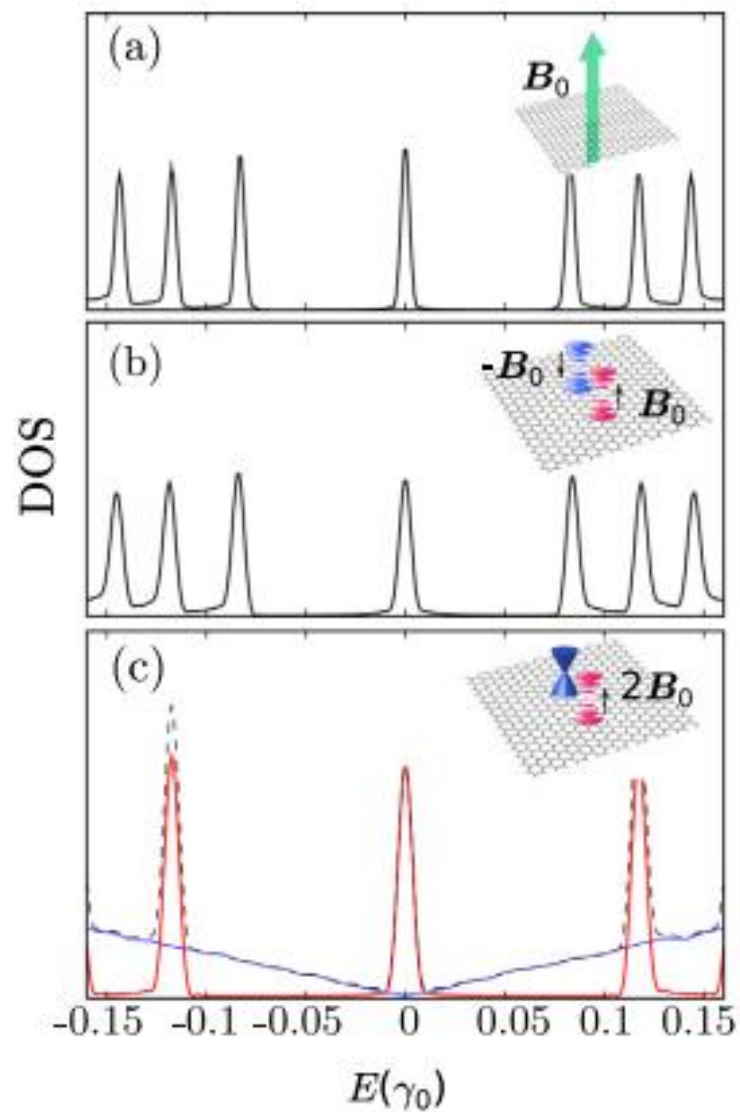
Applications: Canted quantum Spin Hall effect

$$\langle \vec{S} \rangle_{E_F} \approx \langle S_0 \rangle_{E_F} \hat{u}(\theta)$$

$$\theta \equiv \arctan(\Lambda_z/\Lambda_y) \approx -56^\circ$$



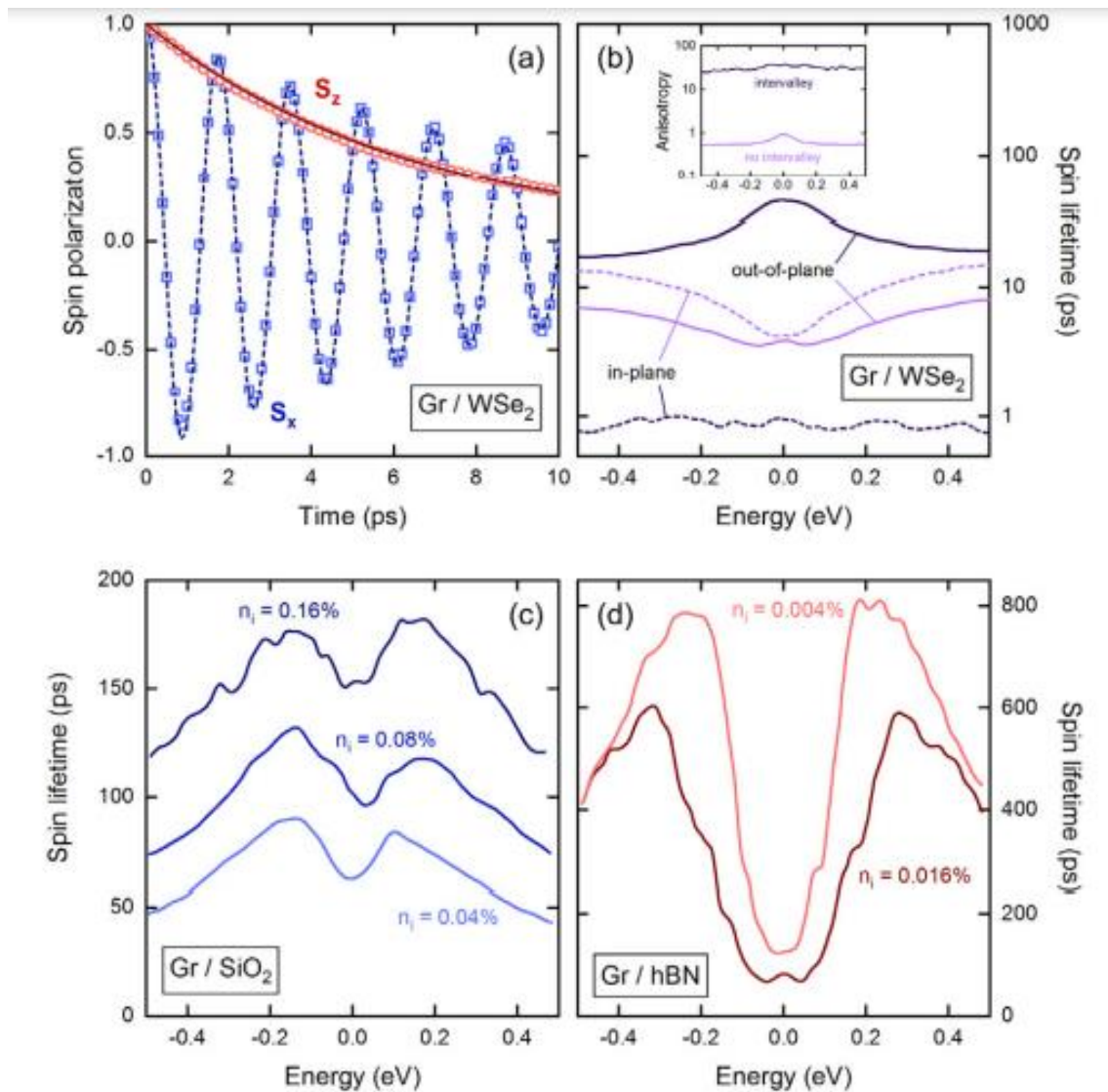
Applications: Valley Hall effect



$$\hat{J}_\alpha^v = \frac{1}{2} (P_v^+ \hat{J}_\alpha P_v^+ - P_v^- \hat{J}_\alpha P_v^-).$$

$$P_v^\pm = \sum_{\mathbf{k}} \theta(|\mathbf{k} - \mathbf{K}^\pm| - R) |\mathbf{k}\rangle \langle \mathbf{k}|$$

Applications: Spin relaxation



$$\mathbf{S}(E, t) = \frac{1}{2} \frac{\langle \phi(t) | \hat{\mathbf{S}} \delta(E - \hat{H}) | \phi(t) \rangle + \text{h.c.}}{\langle \phi(t) | \delta(E - \hat{H}) | \phi(t) \rangle},$$



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