

General Maxwell-Bloch modelling of self-induced transparency in N-level atom

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INTRODUCTION

We investigate self-induced transparency (SIT) effect caused by Rabi oscillations in N -level atom systems implemented in multiple quantum well superlattice. The common models typically implement a two-level Maxwell-Bloch model [1] and consider SIT effect in quantum systems that ideally match with theoretical predictions. The realistic quantum system, as a heterostructure superlattice, however have design-dependent parameters and a ladder of quantum states which cannot be justified by only two levels. We develop a model in which we couple full Maxwell equations to Bloch equations comprised of N levels. We consider scattering of electrons with longitudinal optical phonons, acoustic phonons, interface roughness, impurities and other electrons that provide a realistic disipater within Bloch equations which is similar to the modelling of transport in terahertz quantum cascade lasers [2]. Our model allows propagation of arbitrary input signal and offers detail information of dynamics within N -level atom. This model opens new insights in developing strategies for passive mode locking of THz QCLs via SIT effect.

We analyse ideal two level atom (Fig. 1 and 2) and an exemplary superlattice (Fig. 3) with calculated dephasing time of 0.6 ps, and driven by 3 ps optical pulse at 2 THz frequency (Fig. 4). The system generates partial Rabi flops between the levels at the pump's frequency and complex multiphoton absorption (Fig. 5) that can only be modelled by N -level model.

MODEL

The general formulation of Maxwell-Bloch model for N -level atom consists of two Maxwell equations:

$$\begin{aligned} \frac{\partial H}{\partial t} &= \frac{1}{\mu} \frac{\partial E}{\partial z} \\ \frac{\partial E}{\partial t} &= \frac{1}{\epsilon} \frac{\partial H}{\partial z} - \frac{\sigma}{\epsilon} E - \frac{1}{\epsilon} \frac{\partial P}{\partial t} \end{aligned} \quad (1)$$

where the polarisation term is calculated by coupling this equation to the minimal set of $N(N + 1)/2$ Bloch equations:

$$\begin{aligned} \frac{\partial P}{\partial t} &= 2en_{3D} \sum_j \sum_{i>j} \left(\frac{z_{ji}}{\tau_{||ij}} X_{ij} - z_{ji}\omega_{ij} Y_{ij} \right) \\ \frac{\partial \rho_{ii}}{\partial t} &= -\frac{\rho_{ii}}{\tau_i} + \sum_{j \neq i} \left(\frac{2eE}{\hbar} z_{ij} Y_{ji} \text{sgn}(j-i) + \frac{\rho_{jj}}{\tau_{ji}} \right) \end{aligned}$$

Write only for $i > j$ on the left hand side :

$$\begin{aligned} \frac{\partial X_{ij}}{\partial t} &= \omega_{ij} Y_{ij} - \frac{X_{ij}}{\tau_{||ij}} \\ &+ \frac{eE}{\hbar} \sum_{k \neq i,j} (z_{ik} Y_{kj} \text{sgn}(k-j) - z_{jk} Y_{ik} \text{sgn}(i-k)) \\ \frac{\partial Y_{ij}}{\partial t} &= -\omega_{ij} X_{ij} - \frac{Y_{ij}}{\tau_{||ij}} - \frac{eE}{\hbar} \sum_{k \neq i,j} (z_{ik} X_{kj} - z_{jk} X_{ik}) \\ &- \frac{ez_{ij}E}{\hbar} (\rho_{jj} - \rho_{ii}) \end{aligned} \quad (2)$$

where ρ_{ii} (populations) are diagonal elements of the density matrix, while the off-diagonal elements are written as $\rho_{ij} = X_{ij} + jY_{ij}$ allowing a system of equations in real plane, τ_i represent the state lifetimes, τ_{ij} are scattering times while $\tau_{||ij}$ are dephasing times modelled as in [2], z_{ij} are dipole elements and n_{3D} is volume doping density. The Maxwell equations are solved on a Yee grid in order to ensure numerical stability of the system. The normalisation condition $\sum_i \rho_{ii} = 1$ is implemented by replacing Bloch equation for the last state as $\rho_{NN} = 1 - \sum_{i \neq N} \rho_{ii}$ and in order to increase the numerical accuracy of the system, Bloch equations are solved within predictor-corrector algorithm. For $N = 2$, this system directly folds into two level atom model [1]. Quantum system parameters ($z_{ij}, \tau_i, \tau_{ij}, \tau_{||ij}$) are calculated by self-self consistent Schrödinger-Poisson algorithm [2] which offers realistic values as this model provides high degree of accuracy for modelling transport in THz QCLs. The input field $E(z = 0, t)$ can take an arbitrary function, and we analysed secant signal for ideal two level atom and experimental pulse for superlattice quantum system.

REFERENCES

- [1] R.W. Ziolkowski *et.al*, PRA 52, 4, 3082 (1995).
- [2] A. Demić *et.al*, AIP Advances 9 no. 9, 095019 (2019).

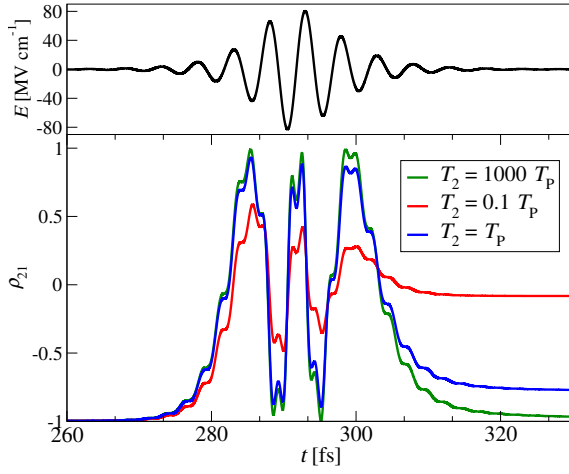


Fig. 1: Top: Secant signal $E_0^{6\pi} = 63 \text{ MV cm}^{-1}$ that analytically causes three Rabi SIT flops of an ideal two level atom from [1]. Bottom: Population inversion when dephasing time (T_2) is larger, equal and smaller than signal width T_P .

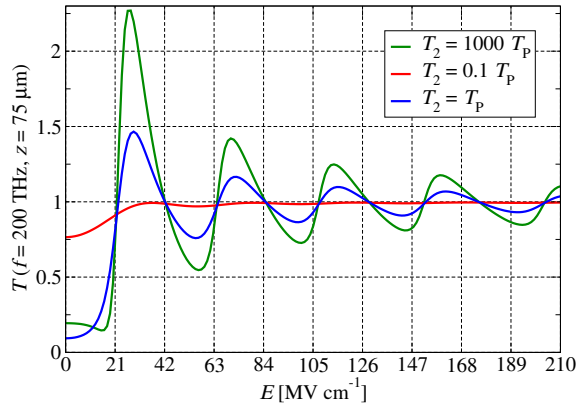


Fig. 2: Transmittance at resonant frequency (200 THz) of ideal two level atom from [1] for different electrical field amplitudes when dephasing time (T_2) is larger, equal and smaller than signal width T_P . Oscillations at $E_0^{2\pi} = 21 \text{ MV cm}^{-1}$ can be clearly observed.

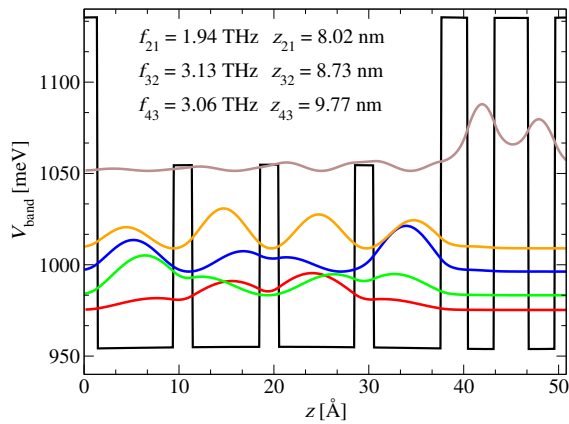


Fig. 3: Bandstructure of superlattice $\text{GaAs}/\text{Al}_x\text{Ga}_{1-x}\text{As}$. designed for absorption at 1.94 THz. Conduction band profile and wavefunction moduli squared are presented. The layer structure is **28/80/20/71/20/80/20/71/28/28/36/28**, barriers are shown in bold text ($x=0.11$ and 0.2) and underlined wells are doped by 2 cm^{-1} .

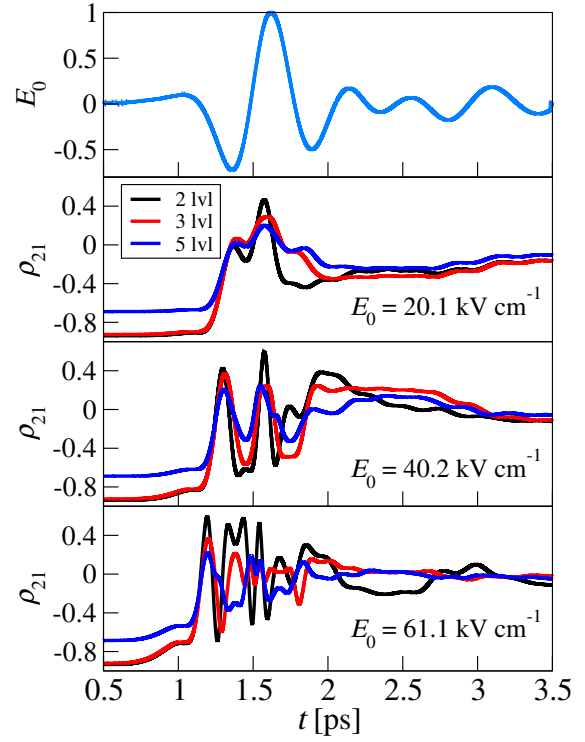


Fig. 4: Top: Experimental input signal generated at resonance of $\sim 2 \text{ THz}$ presented with normalised amplitude that can be varied experimentally and within the model. Bottom: Population inversion at different input field amplitudes when 2,3 and 5 levels are considered in model in Eqs. (1) and (2). Quantum system parameters are calculated for the structure in Fig. 3, displaying dephasing time $\tau_{||21} = 0.6 \text{ ps}$.

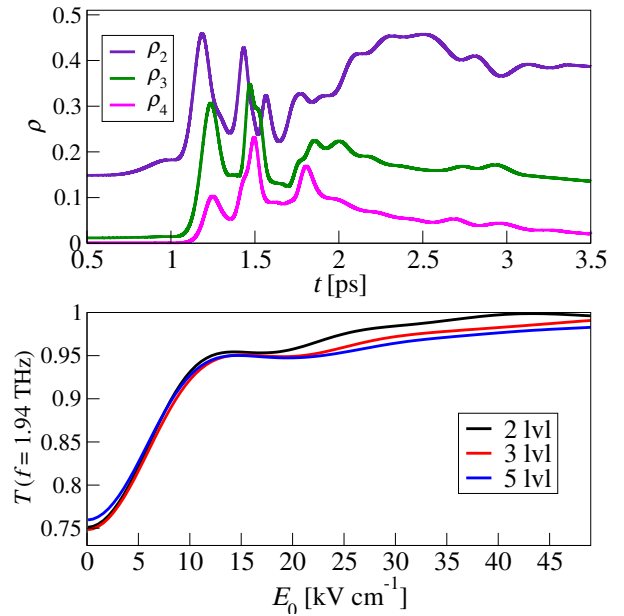


Fig. 5: Top: Populations (in %) at input electric field $E_0 = 40.2 \text{ kV cm}^{-1}$ when five levels are considered in the model. As the signal has spectral contributions $\sim 3 \text{ THz}$ and structure's dipole elements are significantly large, complex multiphoton absorption is observed. This, combined with short dephasing time is responsible for the incomplete Rabi oscillations in Fig. 4. Bottom: Transmittance at resonant frequency (1.94 THz) for different electrical field amplitudes when 2,3 and 5 levels are considered in the model.